

Stability and composition of CH₄-rich clathrate hydrates in the present martian subsurface

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Abstract

In the past years, multiple detections of methane in Mars' atmosphere have raised numerous questions about its potential sources. Independently of methane formation mechanism(s), CH₄ produced in the past or at the present day could be stored in subsurface reservoirs such as clathrate hydrates. In this work, global maps of stability depth of CH₄-rich clathrate hydrates in the martian soil are obtained by including in the subsurface model a top layer with thermal properties that fit with the thermal inertia derived from MGS TES observations (Putzig and Mellon, 2007) and by taking into account the spatial variations of the surface heat flow (Parro et al., 2017). In addition, the spatial distribution of stable methane clathrates is investigated in the presence of eutectic NaCl and CaCl₂ brines. The influence of gas phase composition on clathrate stability and on guests abundance is also examined by considering the CO₂-CH₄-N₂ and CO₂-CH₄-H₂ mixtures.

On present-day Mars, the stability conditions of CH₄-rich clathrates are

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met at a depth of a few meters in high latitude regions and at a few tens of meters deep at the equator. The top of their stability zone is the deepest (~ 68 m) in regions where methane has been locally reported, especially in the area observed by Mumma et al. (2009). It is also in that particular region that the methane clathrate stability zone is the least extended in the martian subsurface, its base reaching ~ 8 km deep in the presence of pure water. This depth can even be 4 times shallower if methane clathrates form in the presence of eutectic CaCl_2 brine. Moreover, the incorporation of nitrogen and hydrogen in mixed $\text{CO}_2\text{-CH}_4$ clathrates allows a better trapping of methane but their presence increases the formation pressure and therefore their stability depth.

Keywords:

Mars, Methane, Clathrate hydrates

1. Introduction

Several detections of methane (CH_4) in the martian atmosphere have been reported from Earth-based and Mars orbit instruments (Formisano et al., 2004; Krasnopolsky et al., 2004; Geminale et al., 2008, 2011; Mumma et al., 2009; Fonti and Marzo, 2010) with abundances ranging up to tens of parts per billion by volume (ppbv). Moreover, in situ measurements performed at Gale crater by the Curiosity rover showed some peaks up to ~ 7 ppbv and an apparent seasonal cycle of the CH_4 background levels (Webster et al., 2015, 2018). However, more time series of methane abundance estimates would be necessary to confirm this periodic variation (Gillen et al., 2020). Recently, first measurements from the ExoMars Trace Gas Orbiter (TGO) mission, especially designed to look for trace levels of methane, did not show any CH_4 detection with an upper limit of about 0.05 ppbv (Korablev et al., 2019).

The only methane sink currently accepted on Mars is its photochemical destruction through its photolysis by the Lyman α line above ~ 80 km altitude, and its oxidation by OH or $\text{O}(^1\text{D})$ below (Haberle et al., 2017). This leads to a CH_4 chemical lifetime of about 300 terrestrial years (Lefèvre and Forget, 2009), much larger than both vertical and horizontal mixing times (10 days and 0.5 year respectively) (Krasnopolsky et al., 2004). Accordingly, methane was expected to be well-mixed in the martian atmosphere. Both remote-sensing and in situ observations have nonetheless evidenced substantial variations of the CH_4 mixing ratio implying an atmospheric lifetime of less than 200 days (Lefèvre and Forget, 2009) and suggesting the presence of localized release and fast removal processes. Methane variability on short time scales has raised numerous questions and the validity of CH_4 detections

26 have been debated (Zahnle et al., 2011; Villanueva et al., 2013; Zahnle, 2015;
27 Webster et al., 2018) on the basis that methane retrievals were distorted
28 due to improper corrections of telluric lines or that some residual terrestrial
29 methane was present in the rover chamber. However, erroneous detections
30 seem surprising given that unsuccessful searches for CH₄ on Mars have also
31 been reported from ground-based (Mumma et al., 2009; Villanueva et al.,
32 2013) and Curiosity (Webster et al., 2013) measurements.

33 Although an in situ spike detection has been recently confirmed from orbit
34 by the Mars Express team who reanalyzed data recorded in 2013 (Giuranna
35 et al., 2019), the existence of unknown and rapid destruction mechanisms is
36 required to explain simultaneously the presence of intermittent CH₄ peaks
37 at the surface and the TGO global non-detection. Several methane loss
38 processes have been proposed such as heterogeneous chemistry (Lefèvre and
39 Forget, 2009), triboelectricity (Farrell et al., 2006; Atreya et al., 2007) or
40 physical and chemical sequestration in the soil (Gough et al., 2010; Jensen
41 et al., 2014) but none of them has been identified to date.

42 The observed methane variability results from a competition between
43 sinks and sources. It has been proposed that methane on Mars could origi-
44 nate from biological (Atreya et al., 2007; Weiss et al., 2000), geological (Oze
45 and Sharma, 2005; Etiope and Sherwood Lollar, 2013) or exogenous (Kress
46 and McKay, 2004; Keppler et al., 2012) generation mechanisms. On the other
47 hand, CH₄ produced in the past or at the present day could be stored in sub-
48 surface reservoirs such as clathrate hydrates (Chastain and Chevrier, 2007)
49 or zeolites (Holmes et al., 2015; Mousis et al., 2016). Clathrate hydrates, also
50 called clathrates, are crystalline compounds constituted by cages formed by

51 hydrogen-bonded water molecules inside of which guest gas molecules are
52 trapped. Thermodynamic conditions prevailing in the martian subsurface
53 favor clathrate formation from 15 m down to 24 km deep (Max et al., 2011)
54 and kinetic experiments (Gainey and Elwood Madden, 2012) showed that
55 their dissociation, stimulated by pressure and temperature changes, is a fea-
56 sible mechanism for near-surface methane release.

57 Clathrates are known for their capacity to store a large amount of gas.
58 At standard pressure and temperature conditions, one cubic meter of CH₄
59 clathrates contains more than 160 cubic meters of methane and the dissocia-
60 tion of 2×10^3 m³ of such a reservoir would be required each year to explain
61 a steady-state value of 10 ppb in the martian atmosphere. Their actual pres-
62 ence depends on the availability of methane and water. Clathrate deposits
63 could therefore be very localized and occupy only a small part of the martian
64 cryosphere, which would be consistent with the observed spatial heterogene-
65 ity of methane. Near-surface reservoirs could episodically discharge methane
66 in the atmosphere following destabilizing events such as impacts, climatic
67 changes or faults and landslides generated by seismicity. On the other hand,
68 one can suggest the formation of clathrates on the south permanent polar cap
69 as a potential methane sink. However, this loss mechanism has been showed
70 not to play a significant role in the dynamics of atmospheric methane on
71 present-day Mars (Thomas et al., 2009; Herri and Chassefière, 2012).

72 The clathrate global distribution on Mars is expected to be at least as
73 complex as on Earth (Lasue et al., 2015). The formation/dissociation cycle
74 of clathrates and the thickness of their stability zone depend on many factors
75 such as subsurface composition, local heat flow, salinity of the water involved

76 in their formation (Elwood Madden et al., 2007) and climate change due to
77 seasonal and obliquity variations (Prieto-Ballesteros et al., 2006; Root and
78 Elwood Madden, 2012). This work results from the need to provide a bet-
79 ter estimation of the subsurface distribution of clathrates on Mars in order
80 to constrain depth of possible source(s) of methane. We thus define in the
81 following an upper bound to possible clathrate occurrences by determining
82 global stability regions based on temperature-pressure conditions. Moreover,
83 clathrate hydrates on Mars might be exposed to highly concentrated salt en-
84 vironments. Although its effects on the formation and stability of clathrates
85 are not well understood, salinity is an important factor to take into account
86 as it depresses the freezing point of water and thus influences the temperature
87 at which clathrates form and dissociate.

88 In the first part of this paper, a sensitivity study is performed to de-
89 termine variations of methane clathrate stability zone (CSZ) as a function
90 of thermal properties of the subsurface and heat flow. In the second part,
91 global maps of stability depth of CH₄-rich clathrate hydrates in the martian
92 soil are obtained for the first time by including in the subsurface model a top
93 layer with thermal properties that fit with the thermal inertia derived from
94 MGS TES observations (Putzig and Mellon, 2007). In addition, the spatial
95 distribution of stable methane clathrates is investigated in the presence of
96 eutectic NaCl and CaCl₂ brines. Finally, we examine the effect of gas phase
97 composition on clathrate stability and on guests abundance by considering
98 the CO₂-CH₄-N₂ and CO₂-CH₄-H₂ mixtures. The latter is examined since
99 most of the methane generation mechanisms imply H₂ production.

100 **2. Clathrate hydrate stability zone model**

101 *2.1. Clathrate hydrate model*

102 We follow an approach similar to Thomas et al. (2009) where the dis-
103 sociation pressure P_{mix}^{diss} of a multiple guest clathrate is calculated from the
104 dissociation pressures of simple guest clathrates as (Lipenkov and Istomin,
105 2001):

$$P_{mix}^{diss} = \left(\sum_G \frac{x_G}{P_G^{diss}} \right)^{-1} \quad (1)$$

106 where x_G is the molar fraction of species G in the initial gas phase. The
107 dissociation pressure P_G^{diss} of a simple clathrate of guest species G follows an
108 Arrhenius law (Miller, 1961):

$$\log(P_G^{diss}) = A + \frac{B}{T} \quad (2)$$

109 where P_G^{diss} is expressed in Pa and T is the temperature in K. The con-
110 stants A and B fit to experimental data (Deaton and Frost, 1946; Larson,
111 1955; Marshall et al., 1964; Jhaveri and Robinson, 1965; Takenouchi and
112 Kennedy, 1965; Robinson et al., 1967; Miller and Smythe, 1970; Falabella,
113 1975; Holder et al., 1980; Ng and Robinson, 1985; Makogon and Sloan, 1994;
114 Kuhs et al., 2000; Yasuda and Ohmura, 2008; Mohammadi and Richon, 2010,
115 2011; Chapoy et al., 2010) and are listed in Table 1 with the temperature
116 range in which they are valid.

117 When thermodynamic inhibitors such as salts are present in the system,
118 they change the activity of water which affects both the freezing point of
119 pure water and the clathrate stability conditions. As a consequence, when

Table 1: Parameters A and B (K) required in Eq.(2).

Molecule	A	B (K)	T range (K)
CH ₄	9.65	-895.9	< 272.9
CH ₄	21.46	-4124.4	> 272.9
CO ₂	10.13	-1116.9	< 273.1
N ₂	9.74	-702.57	< 271.9
Ar	9.24	-630.86	< 272.2
H ₂	19.70	-3080.7	> 263.85

120 salinity increases in the free water phase resulting in a decrease of the ice -
 121 liquid water transition, it moves the quadruple point (where ice, liquid wa-
 122 ter, clathrate and gaseous phases coexist) towards lower temperatures with
 123 a shift that depends on the inhibitor concentration. The dissociation curve
 124 of clathrate hydrates is thus shifted approximately parallel to the uninhib-
 125 ited phase boundary (on a plot of $\ln P$ versus T) towards lower temperature
 126 and higher pressure (Sloan and Koh, 2007). The dissociation pressure cor-
 127 responding to temperature larger than the eutectic temperature has been
 128 estimated assuming slopes comparable to the pure water system, similarly
 129 to the approach of Elwood Madden et al. (2007) for systems with eutectic
 130 NaCl and CaCl₂ brines.

131 Once the equilibrium pressure is known, the relative abundance of a
 132 species G in a clathrate hydrate is given by:

$$x_G^H = \frac{b_s \theta_{G,s} + b_l \theta_{G,l}}{b_s \sum_J \theta_{J,s} + b_l \sum_J \theta_{J,l}} \quad (3)$$

133 where the sum at the denominator takes into account all the species in the
 134 gas phase and b_s and b_l are the number of small and large cavities per unit
 135 cell respectively. The occupancy fraction θ is calculated following the model
 136 of van der Waals and Platteeuw (1959) described in Appendix A.

137 *2.2. Subsurface model*

138 A one dimensional thermal diffusion equation for subsurface tempera-
 139 ture with depth dependent thermal conductivity, density and specific heat is
 140 solved with a Crank-Nicolson scheme on a grid with variable spacing:

$$\rho(z)c(z)\frac{\partial T(z,t)}{\partial t} = \frac{\partial}{\partial z} \left(k(z)\frac{\partial T(z,t)}{\partial z} \right) + \rho(z)H \quad (4)$$

141 where ρ is density, c is specific heat capacity, T is temperature, z is vertical
 142 coordinate, t is time, k is thermal conductivity and H is radiogenic heat pro-
 143 duction. The latter has only a modest effect for temperature determination
 144 and is set to a constant value (5×10^{-11} W kg $^{-1}$) assuming a vertically ho-
 145 mogeneous distribution of radiogenic heat-producing elements in the planet’s
 146 crust (Hahn et al., 2011).

147 In our simulations to determine global maps of the clathrate stability
 148 zone, the changes in surface temperature and pressure in latitude and lon-
 149 gitude over the martian year are given by the Mars Climate Database v5.2
 150 (“Climatology” scenario, average solar EUV conditions) (Forget et al., 1999;
 151 Millour et al., 2015). The present-day heat flow variations across the mar-
 152 tian surface are taken from the model of Parro et al. (2017) that provides
 153 a surface heat flow varying between 14 and 25 mW m $^{-2}$ with an average
 154 value of 19 mW m $^{-2}$ and the subsurface model is divided in two layers. The
 155 thermal properties of the upper layer are set to fit with the thermal inertia
 156 derived from Mars Global Surveyor Thermal Emission Spectrometer obser-
 157 vations (Putzig and Mellon, 2007). Accordingly, the regolith bulk density ρ_b
 158 (kg m $^{-3}$) of this layer is determined as a function of thermal inertia I (J m $^{-2}$
 159 K $^{-1}$ s $^{-1/2}$) following the empirical relation from Mellon and Jakosky (1993):

$$\rho_b = 150 + 100\sqrt{34.2 + 0.714I} \quad (5)$$

160 with a maximum value assumed to be 2500 kg m^{-3} . The specific heat capacity
 161 c is set to $800 \text{ J kg}^{-1} \text{ K}^{-1}$ and the thermal conductivity k is then determined
 162 using $k = I^2/(\rho_b c)$. The properties of the lower layer depend on the latitude
 163 of the studied area. Between 50°N and 50°S , the top of the lower layer is
 164 set to 1 m deep and its thermal properties are representative of dry basalt.
 165 Northward of 50°N and southward of 50°S , the lower layer starts at 50 cm deep
 166 and has thermal properties representative of ice-saturated soil. This choice is
 167 based on the subsurface water ice stability predicted by equilibrium models
 168 (Mellon et al., 2004; Schorghofer and Aharonson, 2005). The compression of
 169 pore spaces with depth is modelled following Clifford (1993):

$$\Phi(z) = \Phi_0 \exp(-z/K_0) \quad (6)$$

170 where Φ_0 is the surface porosity (40%), z is the depth and K_0 is the porosity
 171 decay constant ($\sim 2.82 \text{ km}$). The variations of the volumetric heat capacity
 172 and thermal conductivity with depth are then given by:

$$\rho c(z) = (1 - \Phi(z))\rho_{\text{dry}}c_{\text{dry}} + \Phi(z)\rho_{\text{ice}}c_{\text{ice}} \quad (7)$$

$$k(z) = k_{\text{dry}}^{(1-\Phi(z))} k_{\text{ice}}^{(\Phi(z))} \quad (8)$$

173 where properties of water ice and basalt are referred in Table 2.

174 As mentioned above, we consider a basaltic mineralogy for the second
 175 layer of our subsurface model, either in the presence of water ice or not. How-
 176 ever, local variations in the soil material could strongly affect the methane

Table 2: Typical thermophysical properties of geological materials used in this study.

	k (W m ⁻¹ K ⁻¹)	ρ (kg m ⁻³)	c (J kg ⁻¹ K ⁻¹)	I (J m ⁻² K ⁻¹ s ^{-1/2})
Water ice	2.8	920	1960	2247
Basalt	2	3000	800	2190
Ice-saturated soil	2.5	2018	1040	2290
Dry unconsolidated soil	0.045	1650	800	244

177 clathrate distribution in the martian subsurface. Therefore, in the following
 178 section, a sensitivity study is firstly performed to evaluate the impact of the
 179 soil properties, among others, on the stability field of clathrates. The lat-
 180 ter is determined using different thermal properties in our subsurface model
 181 corresponding to thermal inertia ranging from 250 to 2000 J m⁻² K⁻¹ s^{-1/2}.
 182 For these simulations, the volumetric heat capacity ρc is kept constant (1.6
 183 $\times 10^6$ J K⁻¹ m⁻³) while the thermal conductivity k is changed accordingly.

184 The pressure in pore spaces that controls stability of methane clathrate
 185 hydrates in subsurface has two limiting cases: lithostatic pressure when
 186 the pore spaces are closed and atmospheric pressure when they are open
 187 (Schuerger et al., 2013). Clathrates need to be sealed off from the martian
 188 atmosphere to remain stable, except for a very limited region corresponding
 189 to the south polar cap where surface temperatures are cold enough to enable
 190 CO₂ clathrate formation (Mousis et al., 2013). If the subsurface pressure
 191 profile is determined assuming open porosity, pressure in pore spaces is thus
 192 too low to meet the CH₄-rich clathrate stability conditions at any depth and
 193 latitude. Accordingly, all results presented in this work have been calculated
 194 assuming a lithostatic pressure related to isolated porosity.

195 **3. Results**

196 *3.1. Methane clathrate hydrate stability zone: a sensitivity study*

197 In this sensitivity study, unless otherwise is stated, the stability depths
198 are calculated based on the present-day mean annual surface temperatures
199 (Clifford and Parker, 2001) while the surface heat flow and the surface pres-
200 sure are constant and assumed equal to their average value, 19 mW m^{-2}
201 (Parro et al., 2017) and 610 Pa respectively. The clathrate stability zone
202 changes with seasonal variations will be addressed in the following section.

203 Regarding the subsurface composition, the type of soil directly controls
204 the geothermal conditions and therefore the depth of clathrate formation. In
205 Fig. 1, methane clathrate stability zone is determined in materials ranging
206 from ice-cemented soil (A) to dry unconsolidated soil (D).

207 Similarly to previous studies (Chastain and Chevrier, 2007; Root and
208 Elwood Madden, 2012), methane clathrate hydrates are expected to form at
209 shallower depth with increasing latitude and their stability zone is thicker
210 at high latitude. The thermal conductivity of the soil is a key parameter
211 in the determination of the clathrate stability zone. Materials with high
212 thermal conductivity such as sandstone or ice-cemented soil evacuate heat
213 more efficiently and thus maintain lower temperatures, which allows clathrate
214 formation with a stability zone of several kilometers thick. In Fig. 1A, CH_4
215 clathrate stability zone extends from about 45 m deep to 11.8 km deep at
216 the equator and from 0.6 m deep to about 22 km deep at the pole. On
217 the contrary, if thermal conductivity is not high enough to maintain the low
218 temperature required for clathrate formation, clathrate stability conditions
219 are never met in equatorial regions or only on a small subterranean layer a

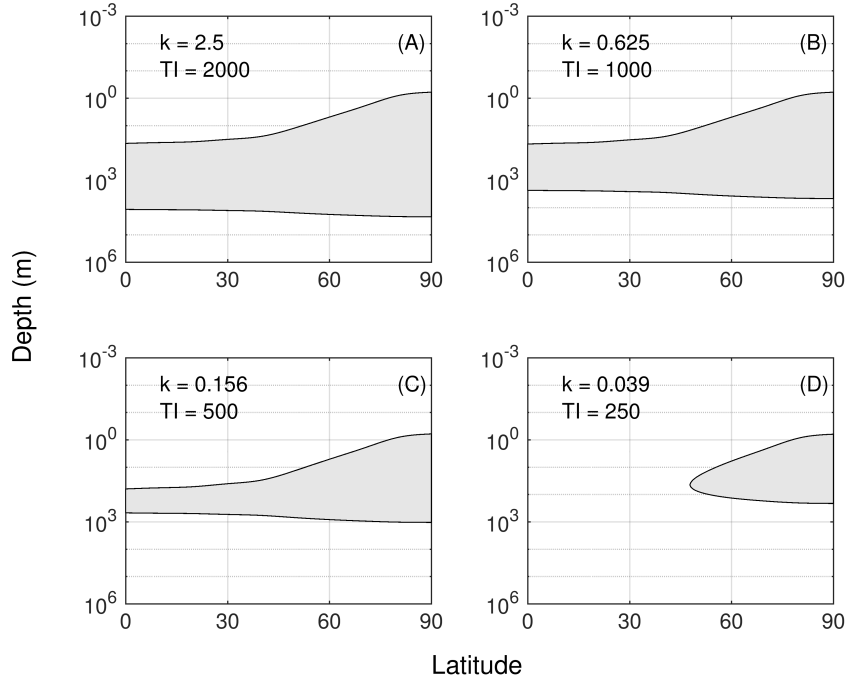


Figure 1: Stability zone of simple CH_4 clathrate in the martian crust for different thermal properties in the subsurface model corresponding to thermal inertia (TI) ranging from 250 to 2000 $\text{J m}^{-2} \text{K}^{-1} \text{s}^{-1/2}$. The volumetric heat capacity ρc is kept constant ($1.6 \times 10^6 \text{ J K}^{-1} \text{m}^{-3}$) while the thermal conductivity k is changed to 2.5 (A), 0.625 (B), 0.156 (C) and 0.039 $\text{W m}^{-1} \text{K}^{-1}$ (D) respectively.

220 few hundred meters thick as shown in Fig. 1.

221 The base of methane CSZ being located several kilometers deep in the
 222 martian subsurface, its variations in depth are highly affected by the local
 223 heat flow as shown in Fig. 2, where calculations have been made for an ice-
 224 cemented soil. As expected, the depth up to which the stability zone extends
 225 decreases with increasing heat flow. The methane clathrate stability zone

226 extends until 15 km deep at the equator when the heat flow is 14 mW m^{-2} ,
 227 while its base is 2 times shallower with a heat flow of 25 mW m^{-2} . This
 228 behavior has significant implications for clathrate formation on Early Mars
 229 where surface heat flows were quite high compared to present-day values.

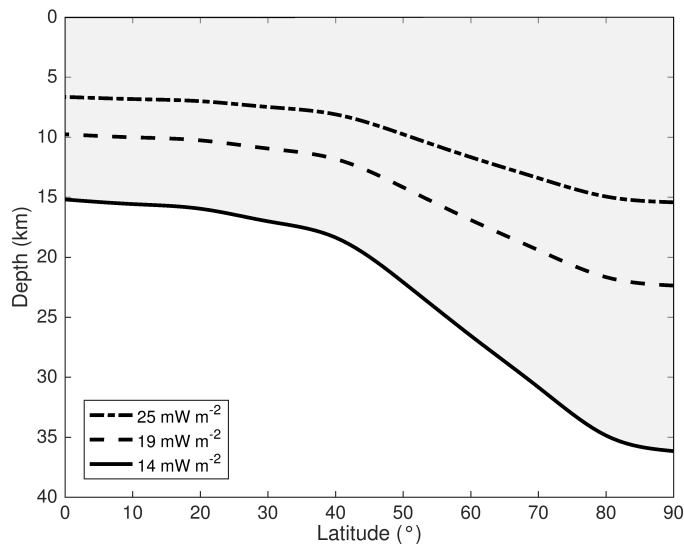


Figure 2: Depth of the base of methane clathrate stability zone as a function of latitude for different values of the heat flow (14 , 19 and 25 mW m^{-2}).

230 *3.2. Methane clathrate hydrate stability zone: global maps*

231 Previously, the stability zone of methane clathrate hydrates in the mar-
 232 tian subsurface has been determined using present-day mean annual surface
 233 temperatures. However, oscillations of temperature due to seasonal vari-
 234 ations can propagate over several meters deep and so reach the clathrate
 235 stability zone at locations where it is close to the surface. Consequently, at
 236 high latitude, the shallow stability zone will be shifted downwards during

237 warm seasons when the soil temperatures are too high to maintain stable
238 clathrates and inversely during the colder time periods. The depth at which
239 temperature fluctuation is reduced by 63% (and thus equals $1/e$ (~ 0.37) of
240 the surface temperature oscillation) is given by the skin depth δ :

$$\delta = \sqrt{\frac{kP}{\rho c \pi}} \quad (9)$$

241 where P is the time period of variations (for seasonal changes, $P = 1$ martian
242 year). With thermal properties similar to those of Fig. 1A, $\delta = 5.4$ m. At
243 latitudes higher than 60° , the methane clathrate stability zone is shallower
244 than this annual skin depth and will therefore be influenced by seasonal vari-
245 ations. For Fig. 1C, δ is 3 times smaller, which makes the seasonal changes in
246 the clathrate stability zone significant beyond a latitude of 75° in that case.

247 Here, fluctuations of the clathrate stability depth over a martian year
248 are studied and global maps of the maximum depth reached by the top
249 of the CSZ are obtained. The latter represents therefore the depth from
250 which clathrate hydrates remain unaffected by seasonal cycles. For these
251 simulations, the subsurface model extends to 100 m deep and mixed CO_2 -
252 CH_4 - N_2 -Ar clathrates are considered with ratios between CO_2 , N_2 and Ar
253 comparable to those measured on Mars. These mixed clathrates trapping
254 the main components of the martian atmosphere together with methane are
255 more representative of those that would presently form on the Red Planet
256 compared to simple CH_4 clathrates studied in the previous section. Two
257 CH_4 mole fractions in the initial gas phase forming clathrates are considered:
258 90%, which leads to the formation of CH_4 -rich clathrates, and 1% resulting
259 in CO_2 -rich clathrates.

260 At first, in order to have an idea of the kind of variations experienced
 261 by the stability zone, we present in Fig. 3 the evolution over time of the top
 262 of the CSZ for clathrates formed from a gas phase with 90% of methane, at
 263 Phoenix landing site. At the beginning of the year, around a solar longitude
 264 (L_S) of 25° , the CSZ is the closest to the surface with a depth of ~ 2 m. Then
 265 the top of the stability zone is shifted downwards with time due to increasing
 266 surface temperature and reaches its maximum depth (~ 5 m) around $L_S =$
 267 198° . Then, the top of the clathrate stability zone approaches again the
 268 surface due to the decrease of surface temperature.

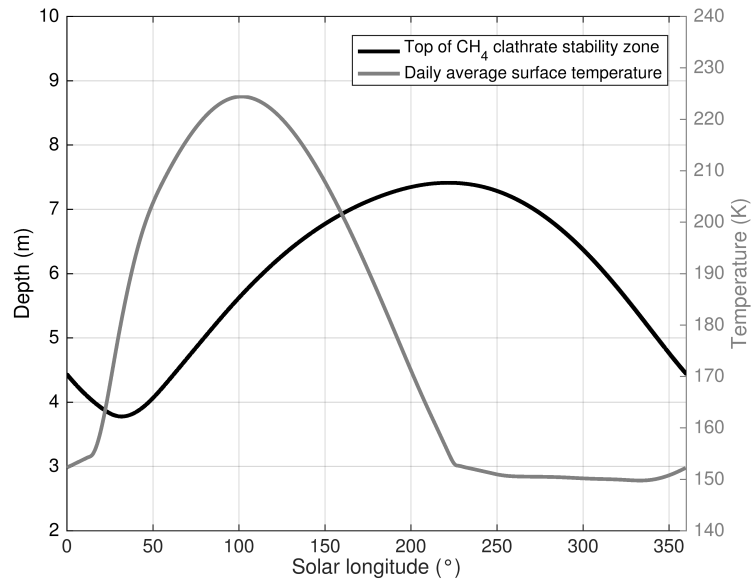


Figure 3: Variation over the year of the top of methane CSZ in the subsurface for clathrates formed from a gas phase with 90% CH₄ and daily average surface temperature at Phoenix landing site. Temperature variation is taken from MCD5.2.

269 The depth from which CH₄-rich clathrates remain unaffected by sea-

270 sonal variations in surface temperature within the martian subsurface at
271 the present day is represented in Fig. 4. The top of the CSZ has been calcu-
272 lated, following the model described in section 2, every 5 degrees and then
273 interpolated every degree. Local detections of methane are reported in black:
274 Gale Crater (Webster et al., 2015, 2018) is represented as a black star, while
275 Syrtis Major, Terra Sabae and Nili Fossae (Mumma et al., 2009) are included
276 in the black rectangle.

277 At the present day, as stability conditions of methane clathrates are met
278 in the near subsurface of Mars, the spatial variation of the top of their stabil-
279 ity zone is strongly dependent on mean annual surface temperature. Conse-
280 quently, CH₄-rich clathrate hydrates are stable closer to the surface at high
281 latitude (few meters deep) and deeper in equatorial areas (few tens of meters
282 deep) due to the larger surface temperature in these regions. By taking into
283 account the thermal inertia observed by MGS TES (Putzig and Mellon, 2007)
284 to set the thermal properties of the upper layer in our subsurface model and
285 then increasing the thermal conductivity of the second layer to correspond
286 to dry basalt, we found that the stability zone of CH₄-rich clathrates is the
287 deepest in regions where methane has been locally reported, especially in the
288 area observed by Mumma et al. (2009) where it is located ~68 m deep and
289 it extends over several kilometers in the martian subsurface. At this depth,
290 clathrates can not be affected by seasonal changes in temperature. However,
291 the addition of CO₂ in CH₄ clathrate hydrates decreases the dissociation
292 pressure in the same temperature conditions, allowing them to be stable at
293 shallower depths (Chastain and Chevrier, 2007). The spatial variation of the
294 top of the CSZ for clathrates formed from a gas phase containing 1% CH₄

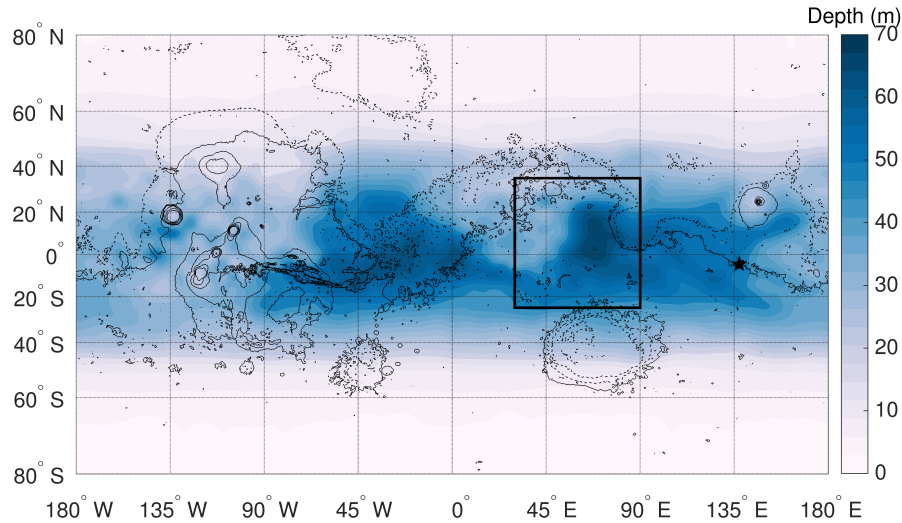


Figure 4: Depth (m) of the top of the clathrate stability zone in present-day martian subsurface for CH_4 -rich clathrates formed from a gas phase with 90% of methane. The first layer of the thermal model has properties that fit with the thermal inertia derived from MGS TES observations (Putzig and Mellon, 2007), while the second layer has properties representative of dry basalt or ice-cemented soil depending on the latitude. Local detections of methane are reported in black: Gale Crater (Webster et al., 2015, 2018) is represented as a black star, while Syrtis Major, Terra Sabae and Nili Fossae (Mumma et al., 2009) are included in the black rectangle.

295 (the main component is therefore CO_2) is shown in Fig. 5. We observe the
 296 same pattern as in Fig. 4 but the stability field is closer to the surface, the
 297 largest stability depth being 28 m. Note that if the thermal properties of
 298 the martian soil are kept identical to those derived from TES MGS observa-
 299 tions in the second layer of the subsurface model, the conditions for CH_4 -rich
 300 clathrate stability are never met between 50°N and 50°S , as it can be deduced

301 from Fig. 1D.

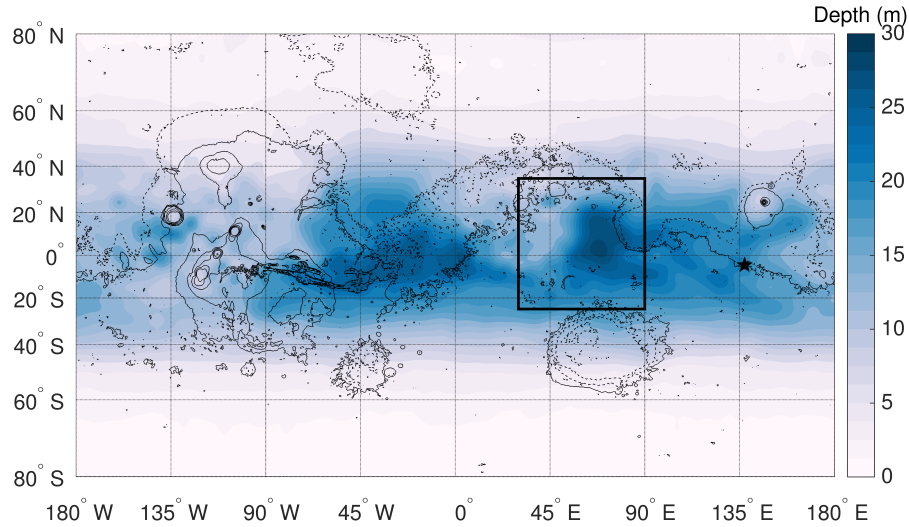


Figure 5: Same as Fig. 4 but for clathrates formed from a gas phase with 1% of methane.

302 As mentioned earlier, the presence of salts moves the stability curve of
303 methane clathrate to lower temperature and higher pressure. This induces
304 an upward shift of the base of the clathrate stability zone, which is therefore
305 significantly narrowed (Elwood Madden et al., 2007). Here, we evaluate this
306 shift across the martian surface by generating global maps for the base of the
307 methane CSZ in the presence of eutectic NaCl and CaCl₂ brines. Results,
308 presented in Fig. 6, show that the stability field of methane clathrate in the
309 crust of Mars extends down to 8-12 km in equatorial regions and 14-21 km
310 in high latitude regions in the presence of pure water. The base of the CSZ
311 is strongly affected by the heat flow, as discussed in the previous section,

312 and is therefore deeper in the northern hemisphere and the Hellas Planitia
313 region. Interestingly, it is closest to the surface in regions where methane
314 was reported during the northern summer 2003 (Mumma et al., 2009). In
315 the presence of eutectic NaCl and CaCl₂ brines, the base of the methane CSZ
316 gets closer to the martian surface and reaches depths of 6-17 km and 2-13 km
317 respectively, shallowest depths corresponding to the region including Syrtis
318 Major and Terra Sabae.

319 *3.3. Variability of methane trapping*

320 Methane currently generated at depth in the martian crust would migrate
321 upwards to reach the local clathrate stability zone and be enclathrated. The
322 trapping of the different species present in the initial gas phase depends on
323 the interactions between the guest molecules and the clathrate cages, as well
324 as the partial pressures of the gases present during and after formation. Fig. 7
325 shows the fraction of methane in mixed CO₂-CH₄-N₂-Ar clathrates as a func-
326 tion of the CH₄ initial abundance in the gas phase for structures I and II. In
327 the presence of CO₂, the formation of CH₄-rich clathrate hydrates is possible
328 only if the initial gas phase is enriched in methane (Thomas et al., 2009).
329 Moreover, methane trapping is slightly enhanced when formation tempera-
330 ture increases. At 270 K, to form clathrates trapping 50 % of methane in
331 their cavities, the initial gas phase must contain about 68 % CH₄ and 57 %
332 CH₄ for structures I and II respectively.

333 Here, we investigate the effect of gas phase composition on clathrate sta-
334 bility and on guests abundance by considering the CO₂-CH₄-N₂ and CO₂-
335 CH₄-H₂ mixtures. The CH₄ production mechanisms often involve H₂ which
336 can be made available locally via serpentinization, while a gas phase enriched

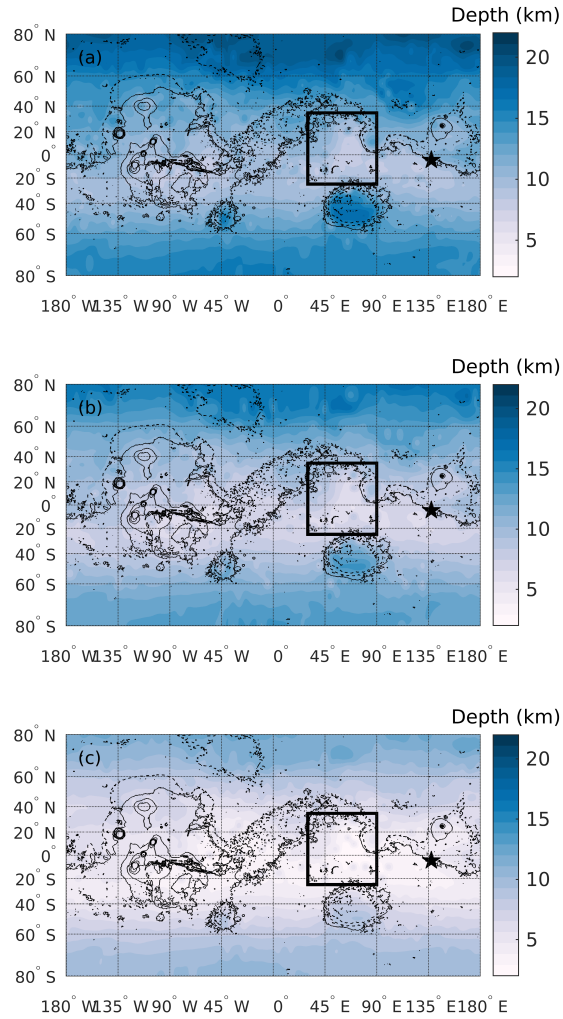


Figure 6: Estimated depth to the base of the clathrate stability zone for methane clathrate in the presence of pure water (a), eutectic NaCl brine (b), and eutectic CaCl₂ brine (c).

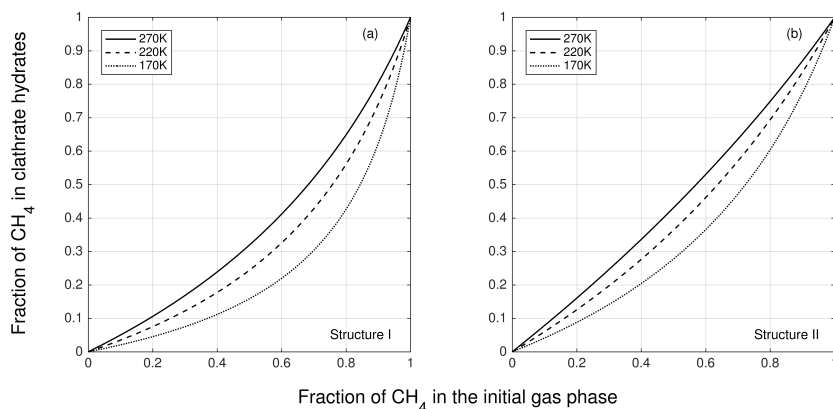


Figure 7: Fraction of CH_4 incorporated in mixed CO_2 - CH_4 - N_2 -Ar clathrate hydrates as a function of the CH_4 fraction in the initial gas phase for different temperatures. The ratios between CO_2 , N_2 and Ar are similar to those measured in the present martian atmosphere. Calculations have been made at the dissociation pressure of clathrates.

337 in N_2 is thought to be more representative of Early Mars atmosphere. Fig. 8
 338 to 9 represent the formation pressure and the evolution of guests abundance,
 339 as a function of CH_4 mole fraction in the gas phase, in binary CO_2 - CH_4 and
 340 ternary CO_2 - CH_4 - N_2 and CO_2 - CH_4 - H_2 sI clathrates at 270 K, respectively.
 341 The ratio between CO_2 and N_2/H_2 in the system is assumed to be 1. Al-
 342 though this ratio is probably unrealistic, the purpose here is to highlight the
 343 general trend for methane trapping in clathrates formed from a gas phase en-
 344 riched in hydrogen or nitrogen. Firstly, it can be seen that the incorporation
 345 of nitrogen or hydrogen in mixed CO_2 - CH_4 clathrates increases the formation
 346 pressure. For low CH_4 fraction (<0.1) in the gaseous phase, the formation
 347 pressure is multiplied by about 1.9 and 2 due to the presence of N_2 and
 348 H_2 respectively. It can also be seen that nitrogen and hydrogen are weakly

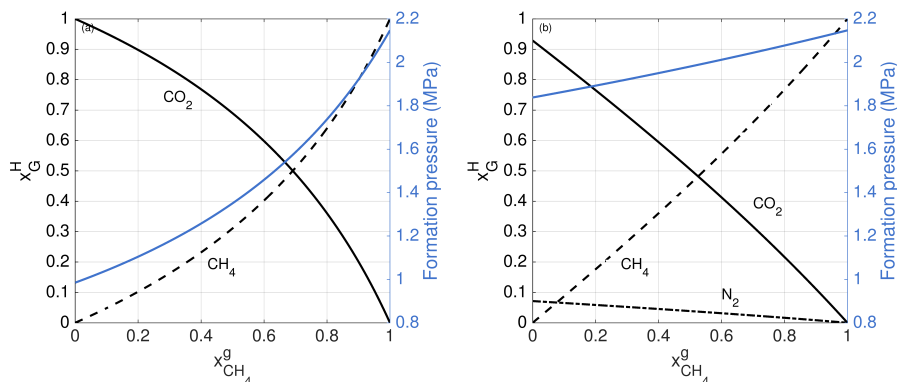


Figure 8: Formation pressure and fraction of the different guests x_G^H in (a) binary CO₂-CH₄ and (b) ternary CO₂-CH₄-N₂ sI clathrates at 270 K and as a function of CH₄ mole fraction in the gas phase. The ratio between CO₂ and N₂ in the gas phase is assumed to be 1.

349 trapped in mixed CO₂-CH₄ sI clathrates, their maximum relative abundance
 350 reaching 0.07 and $3.45 \cdot 10^{-3}$, respectively, when no methane is present in
 351 the gas phase. However, their incorporation in clathrate hydrates allows a
 352 slightly better trapping of methane. For example, in Fig. 9, CO₂-CH₄-H₂ sI
 353 clathrates trapping 50% of methane in their cavities are formed from a gas
 354 phase consisting of 53% CH₄, while 69% are required without the presence of
 355 H₂. In addition, if the structure II is the one that crystallizes, clathrate will
 356 trap a greater fraction of methane, similarly to what is observed in Fig. 7.

357 4. Discussion and conclusions

358 If they exist, methane clathrates on Mars could provide a viable source of
 359 sporadic releases as their stability conditions are currently met in the shallow
 360 subsurface. These crystalline compounds can be stable very near the surface

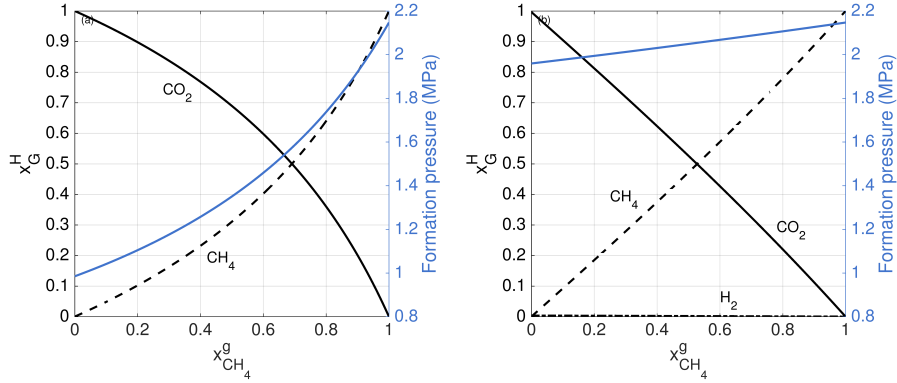


Figure 9: Formation pressure and fraction of the different guests x_G^H in (a) binary CO_2 - CH_4 and (b) ternary CO_2 - CH_4 - H_2 sI clathrates at 270 K and as a function of CH_4 mole fraction in the gas phase. The ratio between CO_2 and H_2 in the gas phase is assumed to be 1.

361 at high latitudes, and can be as close as 20 m to the surface in the tropics.
 362 However, the top of their stability zone is the deepest (~ 68 m) in regions
 363 where methane has been locally reported, especially in the area observed by
 364 Mumma et al. (2009). This depth is too great for clathrates to decompose due
 365 to seasonal changes in temperature. However, destabilizing events, such as
 366 seismic activity or impacts (Mousis et al., 2013), could dissociate clathrates
 367 locally and sporadically in agreement with methane observations. Indeed,
 368 these processes could release overburden pressure, increase local soil temper-
 369 atures or open cracks and fractures that would therefore connect the clathrate
 370 reservoir to the surface leading to its destabilization and the methane release
 371 in the atmosphere.

372 Regarding the base of the methane clathrate stability zone on Mars, we
 373 showed that it is the shallowest (~ 8 km) where Mumma et al. (2009) reported

374 observations of local methane plumes. In addition, when methane clathrates
375 form in the presence of eutectic NaCl and CaCl₂ brines, the base of the
376 stability zone gets closer to the martian surface, reaching in that particular
377 region a depth of 6 km and 2 km respectively. The clathrate decomposition
378 could be triggered at the base of the CSZ by interactions with high salinity
379 fluids (Elwood Madden et al., 2007). In regions where CH₄ plumes have
380 been observed, methane released by those deep clathrates would thus be
381 transported through the crust along a shorter pathway to be either trapped
382 again in subsurface reservoirs or discharged in the martian atmosphere.

383 As shown in section 3, the clathrate stability zone strongly depends on
384 the heat flow values. If the methane production has stopped early in the
385 martian history when surface heat flows were quite high, then formation of
386 CH₄ clathrate reservoirs would have been limited to the stability zone associ-
387 ated to this time period. Thereafter, although the stability field of methane
388 clathrate would have continued to expand with the cooling of the planet,
389 new clathrate formations would not have occurred without the addition of
390 methane in the system. Future studies specifically taking into account the
391 thermal evolution of Mars would thus be required to further constrain depths
392 of ancient clathrate reservoirs.

393 The formation of CH₄-rich clathrate hydrates from a gas phase contain-
394 ing carbon dioxide requires an initial gas phase enriched in methane. If
395 those mixed CO₂-CH₄ clathrates form in contact with a significant fraction
396 of nitrogen, as it may have been the case on Early Mars, or hydrogen, for
397 example related to methane generation mechanisms, they can trap a larger
398 amount of methane in their water cavities compared to clathrates formed

399 from carbon dioxide and methane exclusively. However, although N_2 and H_2
400 are weakly enclathrated, their presence strongly increases the formation pres-
401 sure of clathrate hydrates and therefore their stability depth. These ternary
402 clathrate hydrates trapping methane, if present in the near subsurface, would
403 be located deeper than binary CO_2-CH_4 clathrates. Their destabilization by
404 surface processes is therefore more challenging.

405 As shown in our sensitivity study, the clathrate stability zone is very re-
406 sponsive to the composition of the subsurface material. In our simulations
407 to establish global maps of the clathrate stability zone, we mainly consid-
408 ered a basaltic composition, which is fairly representative of the martian
409 crust. However, local variations in the soil material could strongly affect the
410 methane clathrate distribution displayed on our global maps and even pre-
411 vent clathrate formation at low latitude if the thermal conductivity of the
412 material is quite small. In order to further constrain the clathrate stability
413 field on Mars, geological observations coupled to a comprehensive sensitiv-
414 ity analysis with respect to soil parameters is thus required. In addition, our
415 global maps do not show metastable zones. Indeed, some clathrate reservoirs
416 could be localized outside their present stability field in areas associated to
417 ancient stability zones. These reservoirs could represent essential sources in
418 the study of methane outgassing scenarios since, being unstable, they would
419 gradually release methane from depths possibly very close to the surface. Ex-
420 perimental work on kinetic of clathrate formation and dissociation should be
421 used together with a subsurface thermal model taking into account obliquity
422 variations to evaluate the potential extent and locations of these metastable
423 clathrate reservoirs. Finally, in order to give a better accurate estimation of

424 the methane clathrate occurrences on Mars, a physical model such as the one
 425 presented in this work need to be coupled with topography/geological ana-
 426 lysis to determine possible regions within the clathrate stability zone where
 427 sufficient CH₄ supplies could exist.

428 **Appendix A. Thermodynamic modelling of clathrates**

429 Thermodynamic equilibrium implies the minimization of the Gibbs en-
 430 ergy and the equality of chemical potentials of water μ_w in the ice phase α
 431 and in the clathrate hydrate phase H :

$$\mu_w^H = \mu_w^\alpha \quad (\text{A.1})$$

432 For convenience, this equalization can be rewritten in terms of chemical
 433 potential differences by introducing a hypothetical phase β that corresponds
 434 to a clathrate with empty cavities:

$$\left. \begin{aligned} \Delta\mu_w^{\beta-H} &= \mu_w^\beta - \mu_w^H \\ \Delta\mu_w^{\beta-\alpha} &= \mu_w^\beta - \mu_w^\alpha \end{aligned} \right\} \rightarrow \Delta\mu_w^{\beta-H} = \Delta\mu_w^{\beta-\alpha} \quad (\text{A.2})$$

435 where μ_w^β is the chemical potential of water in the empty clathrate. $\Delta\mu_w^{\beta-\alpha}$
 436 is determined from classical thermodynamics, while $\Delta\mu_w^{\beta-H}$ is evaluated via
 437 statistical thermodynamics. Following the model of van der Waals and Plat-
 438 teeuw (1959), this term can be written as a function of the occupancy fraction
 439 $\theta_{G,i}$ of a species G in a given type i of cage ($i = \text{small or large}$) for a given
 440 type of clathrate structure (I or II) as:

$$\Delta\mu_w^{\beta-H} = -RT \sum_i \nu_i \ln(1 - \sum_G \theta_{G,i}) \quad (\text{A.3})$$

441 where R is the universal gas constant, T is the thermodynamic temperature
 442 and ν_i is the number of cavities of type i per water molecule in the clathrate
 443 unit cell. The occupancy fraction is defined as:

$$\theta_{G,i} = \frac{C_{G,i}(T)f_G(T,P)}{1 + \sum_J C_{J,i}(T)f_J(T,P)} \quad (\text{A.4})$$

444 where $f_G(T,P)$ is the fugacity of the guest G in the gas or liquid phase
 445 and is calculated using the Peng-Robinson equation of state. The Langmuir
 446 constant $C_{G,i}(T)$ characterizes the attractiveness of the cavity i for a species
 447 G . If the cavity is assumed to be perfectly spherical and the guest molecule
 448 can freely rotate in the cage, the Langmuir constant can be written as:

$$C_{G,i}(T) = \frac{4\pi}{k_B T} \int_0^{R_c} \exp\left(\frac{-w_{G,i}(r)}{k_B T}\right) r^2 dr \quad (\text{A.5})$$

449 where k_B is the Boltzmann constant, R_c is the radius of the spherical cavity
 450 and $w_{G,i}(r)$ is the spherically averaged potential energy between the guest
 451 molecule and the cavity. This interaction potential is often described by the
 452 Kihara potential (McKoy and Sinanoğlu, 1963):

$$w_{G,i}(r) = 2z\epsilon \left[\frac{\sigma^{12}}{R_c^{11}r} \left(\delta^{10} + \frac{a}{R_c} \delta^{11} \right) - \frac{\sigma^6}{R_c^5 r} \left(\delta^4 + \frac{a}{R_c} \delta^5 \right) \right] \quad (\text{A.6})$$

453 with

$$\delta^N = \frac{1}{N} \left[\left(1 - \frac{r}{R_c} - \frac{a}{R_c} \right)^{-N} - \left(1 + \frac{r}{R_c} - \frac{a}{R_c} \right)^{-N} \right] \quad (\text{A.7})$$

454 where z is the coordination number of the cavity. The parameters R_c and z
 455 are specific to the type of cage and the clathrate structure and are listed in

456 Table A.1 for sI and sII clathrates. The Kihara parameters ϵ , a and σ are
 457 specific to the guest molecule and are determined by a linear regression fitted
 458 to experimental data. The set of Kihara parameters used in the present work
 459 is given in Table A.2. These parameters have been optimized by Herri and
 460 Chassefière (2012) in a study where they fitted the model of van der Waals
 461 and Platteeuw (1959) to compare the deviation from experimental data of
 462 pure clathrate hydrates at temperatures relevant to Mars.

Table A.1: Clathrate hydrate structure parameters (Sloan and Koh, 2007).

	sI		sII	
	Small	Large	Small	Large
Cavity type	Small	Large	Small	Large
Number of cages per unit cell	2	6	16	8
Average cavity radius ^a (Å)	3.95	4.33	3.91	4.73
Coordination number ^b	20	24	20	28

^a The cavity radius will change with temperature, pressure and guest composition.

^b The number of water molecules per cavity.

Table A.2: Kihara parameters used in the present study. These parameters are derived from Herri and Chassefière (2012) for CH₄, CO₂, N₂ and Ar and from Strobel et al. (2009) for H₂.

Molecule	ϵ/k_B (K)	σ (Å)	a (Å)
CH ₄	166.36	3.0500	0.3834
CO ₂	178.21	2.8730	0.6805
N ₂	133.13	3.0993	0.3526
Ar	174.14	2.9434	0.1840
H ₂	80.424	3.07838	0.1973

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