

Search for the production of ZW and ZZ boson pairs decaying into charged leptons and jets in $p\bar{p}$ collisions at $\sqrt{s} = 1.96$ TeV

T. Aaltonen,²¹ S. Amerio,^{39a,39b} D. Amidei,³¹ A. Anastassov,^{15,w} A. Annovi,¹⁷ J. Antos,¹² G. Apollinari,¹⁵ J. A. Appel,¹⁵ T. Arisawa,⁵² A. Artikov,¹³ J. Asaadi,⁴⁷ W. Ashmanskas,¹⁵ B. Auerbach,² A. Aurisano,⁴⁷ F. Azfar,³⁸ W. Badgett,¹⁵ T. Bae,²⁵ A. Barbaro-Galtieri,²⁶ V. E. Barnes,⁴³ B. A. Barnett,²³ P. Barria,^{41c,41a} P. Bartos,¹² M. Baucé,^{39b,39a} F. Bedeschi,^{41a} S. Behari,¹⁵ G. Bellettini,^{41b,41a} J. Bellinger,⁵⁴ D. Benjamin,¹⁴ A. Beretvas,¹⁵ A. Bhatti,⁴⁵ K. R. Bland,⁵ B. Blumenfeld,²³ A. Bocci,¹⁴ A. Bodek,⁴⁴ D. Bortoletto,⁴³ J. Boudreau,⁴² A. Boveia,¹¹ L. Brigliadori,^{6b,6a} C. Bromberg,³² E. Brucken,²¹ J. Budagov,¹³ H. S. Budd,⁴⁴ K. Burkett,¹⁵ G. Busetto,^{39b,39a} P. Bussey,¹⁹ P. Butti,^{41b,41a} A. Buzatu,¹⁹ A. Calamba,¹⁰ S. Camarda,⁴ M. Campanelli,²⁸ F. Canelli,^{11,dd} B. Carls,²² D. Carlsmith,⁵⁴ R. Carosi,^{41a} S. Carrillo,^{16,m} B. Casal,^{9,k} M. Casarsa,^{48a} A. Castro,^{6b,6a} P. Catastini,²⁰ D. Cauz,^{48b,48c,48a} V. Cavaliere,²² M. Cavalli-Sforza,⁴ A. Cerri,^{26,f} L. Cerrito,^{28,r} Y. C. Chen,¹ M. Chertok,⁷ G. Chiarelli,^{41a} G. Chlachidze,¹⁵ K. Cho,²⁵ D. Chokheli,¹³ A. Clark,¹⁸ C. Clarke,⁵³ M. E. Convery,¹⁵ J. Conway,⁷ M. Corbo,^{15,z} M. Cordelli,¹⁷ C. A. Cox,⁷ D. J. Cox,⁷ M. Cremonesi,^{41a} D. Cruz,⁴⁷ J. Cuevas,^{9,y} R. Culbertson,¹⁵ N. d'Ascenzo,^{15,v} M. Datta,^{15,gg} P. de Barbaro,⁴⁴ L. Demortier,⁴⁵ M. Deninno,^{6a} M. D'Errico,^{39b,39a} F. Devoto,²¹ A. Di Canto,^{41b,41a} B. Di Ruzza,^{15,q} J. R. Dittmann,⁵ S. Donati,^{41b,41a} M. D'Onofrio,²⁷ M. Dorigo,^{48d,48a} A. Driutti,^{48b,48c,48a} K. Ebina,⁵² R. Edgar,³¹ A. Elagin,⁴⁷ R. Erbacher,⁷ S. Errede,²² B. Esham,²² S. Farrington,³⁸ J. P. Fernández Ramos,²⁹ R. Field,¹⁶ G. Flanagan,^{15,t} R. Forrest,⁷ M. Franklin,²⁰ J. C. Freeman,¹⁵ H. Frisch,¹¹ Y. Funakoshi,⁵² C. Galloni,^{41b,41a} A. F. Garfinkel,⁴³ P. Garosi,^{41c,41a} H. Gerberich,²² E. Gerchtein,¹⁵ S. Giagu,^{46a} V. Giakoumopoulou,³ K. Gibson,⁴² C. M. Ginsburg,¹⁵ N. Giokaris,³ P. Giromini,¹⁷ G. Giurgiu,²³ V. Glagolev,¹³ D. Glenzinski,¹⁵ M. Gold,³⁴ D. Goldin,⁴⁷ A. Golossanov,¹⁵ G. Gomez,⁹ G. Gomez-Ceballos,³⁰ M. Goncharov,³⁰ O. González López,²⁹ I. Gorelov,³⁴ A. T. Goshaw,¹⁴ K. Goulianos,⁴⁵ E. Gramellini,^{6a} S. Grinstein,⁴ C. Grosso-Pilcher,¹¹ R. C. Group,^{51,15} J. Guimaraes da Costa,²⁰ S. R. Hahn,¹⁵ J. Y. Han,⁴⁴ F. Happacher,¹⁷ K. Hara,⁴⁹ M. Hare,⁵⁰ R. F. Harr,⁵³ T. Harrington-Taber,^{15,n} K. Hatakeyama,⁵ C. Hays,³⁸ J. Heinrich,⁴⁰ M. Herndon,⁵⁴ A. Hocker,¹⁵ Z. Hong,⁴⁷ W. Hopkins,^{15,g} S. Hou,¹ R. E. Hughes,³⁵ U. Husemann,⁵⁵ M. Hussein,^{32,bb} J. Huston,³² G. Introzzi,^{41e,41f,41a} M. Iori,^{46b,46a} A. Ivanov,^{7,p} E. James,¹⁵ D. Jang,¹⁰ B. Jayatilaka,¹⁵ E. J. Jeon,²⁵ S. Jindariani,¹⁵ M. Jones,⁴³ K. K. Joo,²⁵ S. Y. Jun,¹⁰ T. R. Junk,¹⁵ M. Kambeitz,²⁴ T. Kamon,^{25,47} P. E. Karchin,⁵³ A. Kasmi,⁵ Y. Kato,^{37,o} W. Ketchum,^{11,hh} J. Keung,⁴⁰ B. Kilminster,^{15,dd} D. H. Kim,²⁵ H. S. Kim,²⁵ J. E. Kim,²⁵ M. J. Kim,¹⁷ S. H. Kim,⁴⁹ S. B. Kim,²⁵ Y. J. Kim,²⁵ Y. K. Kim,¹¹ N. Kimura,⁵² M. Kirby,¹⁵ K. Knoepfel,¹⁵ K. Kondo,^{52,a} D. J. Kong,²⁵ J. Konigsberg,¹⁶ A. V. Kotwal,¹⁴ M. Kreps,²⁴ J. Kroll,⁴⁰ M. Kruse,¹⁴ T. Kuhr,²⁴ M. Kurata,⁴⁹ A. T. Laasanen,⁴³ S. Lammel,¹⁵ M. Lancaster,²⁸ K. Lannon,^{35,x} G. Latino,^{41c,41a} H. S. Lee,²⁵ J. S. Lee,²⁵ S. Leo,^{41a} S. Leone,^{41a} J. D. Lewis,¹⁵ A. Limosani,^{14,s} E. Lipeles,⁴⁰ A. Lister,^{18,b} H. Liu,⁵¹ Q. Liu,⁴³ T. Liu,¹⁵ S. Lockwitz,⁵⁵ A. Loginov,⁵⁵ D. Lucchesi,^{39b,39a} A. Lucà,¹⁷ J. Lueck,²⁴ P. Lujan,²⁶ P. Lukens,¹⁵ G. Lungu,⁴⁵ J. Lys,²⁶ R. Lysak,^{12,e} R. Madrak,¹⁵ P. Maestro,^{41c,41a} S. Malik,⁴⁵ G. Manca,^{27,c} A. Manousakis-Katsikakis,³ L. Marchese,^{6a,ii} F. Margaroli,^{46a} P. Marino,^{41d,41a} M. Martínez,⁴ K. Matera,²² M. E. Mattson,⁵³ A. Mazzacane,¹⁵ P. Mazzanti,^{6a} R. McNulty,^{27,j} A. Mehta,²⁷ P. Mehtala,²¹ C. Mesropian,⁴⁵ T. Miao,¹⁵ D. Mietlicki,³¹ A. Mitra,¹ H. Miyake,⁴⁹ S. Moed,¹⁵ N. Moggi,^{6a} C. S. Moon,^{15,z} R. Moore,^{15,ee,ff} M. J. Morello,^{41d,41a} A. Mukherjee,¹⁵ Th. Muller,²⁴ P. Murat,¹⁵ M. Mussini,^{6b,6a} J. Nachtman,^{15,n} Y. Nagai,⁴⁹ J. Naganoma,⁵² I. Nakano,³⁶ A. Napier,⁵⁰ J. Nett,⁴⁷ C. Neu,⁵¹ T. Nigmanov,⁴² L. Nodulman,² S. Y. Noh,²⁵ O. Norriella,²² L. Oakes,³⁸ S. H. Oh,¹⁴ Y. D. Oh,²⁵ I. Oksuzian,⁵¹ T. Okusawa,³⁷ R. Orava,²¹ L. Ortolan,⁴ C. Pagliarone,^{48a} E. Palencia,^{9,f} P. Palni,³⁴ V. Papadimitriou,¹⁵ W. Parker,⁵⁴ G. Pauletta,^{48b,48c,48a} M. Paulini,¹⁰ C. Paus,³⁰ T. J. Phillips,¹⁴ G. Piacentino,^{41a} E. Pianori,⁴⁰ J. Pilot,⁷ K. Pitts,²² C. Plager,⁸ L. Pondrom,⁵⁴ S. Poprocki,^{15,g} K. Potamianos,²⁶ A. Pranko,²⁶ F. Prokoshin,^{13,aa} F. Ptohos,^{17,h} G. Punzi,^{41b,41a} N. Ranjan,⁴³ I. Redondo Fernández,²⁹ P. Renton,³⁸ M. Rescigno,^{46a} F. Rimondi,^{6a,a} L. Ristori,^{41a,15} A. Robson,¹⁹ T. Rodriguez,⁴⁰ S. Rolli,^{50,i} M. Ronzani,^{41b,41a} R. Roser,¹⁵ J. L. Rosner,¹¹ F. Ruffini,^{41c,41a} A. Ruiz,⁹ J. Russ,¹⁰ V. Rusu,¹⁵ W. K. Sakumoto,⁴⁴ Y. Sakurai,⁵² L. Santi,^{48b,48c,48a} K. Sato,⁴⁹ V. Saveliev,^{15,v} A. Savoy-Navarro,^{15,z} P. Schlabach,¹⁵ E. E. Schmidt,¹⁵ T. Schwarz,³¹ L. Scodellaro,⁹ F. Scuri,^{41a} S. Seidel,³⁴ Y. Seiya,³⁷ A. Semenov,¹³ F. Sforza,^{41b,41a} S. Z. Shalhout,⁷ T. Shears,²⁷ P. F. Shepard,⁴² M. Shimojima,^{49,u} M. Shochet,¹¹ I. Shreyber-Tecker,³³ A. Simonenko,¹³ K. Sliwa,⁵⁰ J. R. Smith,⁷ F. D. Snider,¹⁵ H. Song,⁴² V. Sorin,⁴ R. St. Denis,¹⁹ M. Stancari,¹⁵ D. Stentz,^{15,w} J. Strologas,³⁴ Y. Sudo,⁴⁹ A. Sukhanov,¹⁵ I. Suslov,¹³ K. Takemasa,⁴⁹ Y. Takeuchi,⁴⁹ J. Tang,¹¹ M. Tecchio,³¹ P. K. Teng,¹ J. Thom,^{15,g} E. Thomson,⁴⁰ V. Thukral,⁴⁷ D. Toback,⁴⁷ S. Tokar,¹² K. Tollefson,³² T. Tomura,⁴⁹ D. Tonelli,^{15,f} S. Torre,¹⁷ D. Torretta,¹⁵ P. Totaro,^{39a} M. Trovato,^{41d,41a} F. Ukegawa,⁴⁹ S. Uozumi,²⁵ F. Vázquez,^{16,m} G. Velev,¹⁵ C. Vellidis,¹⁵ C. Vernieri,^{41d,41a} M. Vidal,⁴³ R. Vilar,⁹ J. Vizán,^{9,cc} M. Vogel,³⁴ G. Volpi,¹⁷ P. Wagner,⁴⁰ R. Wallny,^{15,k} S. M. Wang,¹

D. Waters,²⁸ W. C. Wester III,¹⁵ D. Whiteson,^{40,d} A. B. Wicklund,² S. Wilbur,⁷ H. H. Williams,⁴⁰ J. S. Wilson,³¹
P. Wilson,¹⁵ B. L. Winer,³⁵ P. Wittich,^{15,g} S. Wolbers,¹⁵ H. Wolfe,³⁵ T. Wright,³¹ X. Wu,¹⁸ Z. Wu,⁵ K. Yamamoto,³⁷
D. Yamato,³⁷ T. Yang,¹⁵ U. K. Yang,²⁵ Y. C. Yang,²⁵ W.-M. Yao,²⁶ G. P. Yeh,¹⁵ K. Yi,^{15,n} J. Yoh,¹⁵ K. Yorita,⁵²
T. Yoshida,^{37,1} G. B. Yu,¹⁴ I. Yu,²⁵ A. M. Zanetti,^{48a} Y. Zeng,¹⁴ C. Zhou,¹⁴ and S. Zucchelli^{6b,6a}

(CDF Collaboration)

¹*Institute of Physics, Academia Sinica, Taipei, Taiwan 11529, Republic of China*

²*Argonne National Laboratory, Argonne, Illinois 60439, USA*

³*University of Athens, 157 71 Athens, Greece*

⁴*Institut de Física d'Altes Energies, ICREA, Universitat Autònoma de Barcelona, E-08193 Bellaterra (Barcelona), Spain*

⁵*Baylor University, Waco, Texas 76798, USA*

^{6a}*Istituto Nazionale di Fisica Nucleare Bologna, I-40127 Bologna, Italy*

^{6b}*University of Bologna, I-40127 Bologna, Italy*

⁷*University of California, Davis, Davis, California 95616, USA*

⁸*University of California, Los Angeles, Los Angeles, California 90024, USA*

⁹*Instituto de Física de Cantabria, CSIC-University of Cantabria, 39005 Santander, Spain*

¹⁰*Carnegie Mellon University, Pittsburgh, Pennsylvania 15213, USA*

¹¹*Enrico Fermi Institute, University of Chicago, Chicago, Illinois 60637, USA*

¹²*Comenius University, 842 48 Bratislava, Slovakia and Institute of Experimental Physics, 040 01 Kosice, Slovakia*

¹³*Joint Institute for Nuclear Research, RU-141980 Dubna, Russia*

¹⁴*Duke University, Durham, North Carolina 27708, USA*

¹⁵*Fermi National Accelerator Laboratory, Batavia, Illinois 60510, USA*

¹⁶*University of Florida, Gainesville, Florida 32611, USA*

¹⁷*Laboratori Nazionali di Frascati, Istituto Nazionale di Fisica Nucleare, I-00044 Frascati, Italy*

¹⁸*University of Geneva, CH-1211 Geneva 4, Switzerland*

¹⁹*Glasgow University, Glasgow G12 8QQ, United Kingdom*

²⁰*Harvard University, Cambridge, Massachusetts 02138, USA*

²¹*Division of High Energy Physics, Department of Physics, University of Helsinki, FIN-00014, Helsinki, Finland;
Helsinki Institute of Physics, FIN-00014, Helsinki, Finland*

²²*University of Illinois, Urbana, Illinois 61801, USA*

²³*The Johns Hopkins University, Baltimore, Maryland 21218, USA*

²⁴*Institut für Experimentelle Kernphysik, Karlsruhe Institute of Technology, D-76131 Karlsruhe, Germany*

²⁵*Center for High Energy Physics: Kyungpook National University, Daegu 702-701, Korea
and Seoul National University, Seoul 151-742, Korea and Sungkyunkwan University, Suwon 440-746, Korea
and Korea Institute of Science and Technology Information, Daejeon 305-806, Korea
and Chonnam National University, Gwangju 500-757, Korea
and Chonbuk National University, Jeonju 561-756, Korea
and Ewha Womans University, Seoul 120-750, Korea*

²⁶*Ernest Orlando Lawrence Berkeley National Laboratory, Berkeley, California 94720, USA*

²⁷*University of Liverpool, Liverpool L69 7ZE, United Kingdom*

²⁸*University College London, London WC1E 6BT, United Kingdom*

²⁹*Centro de Investigaciones Energeticas Medioambientales y Tecnológicas, E-28040 Madrid, Spain*

³⁰*Massachusetts Institute of Technology, Cambridge, Massachusetts 02139, USA*

³¹*University of Michigan, Ann Arbor, Michigan 48109, USA*

³²*Michigan State University, East Lansing, Michigan 48824, USA*

³³*Institution for Theoretical and Experimental Physics, ITEP, Moscow 117259, Russia*

³⁴*University of New Mexico, Albuquerque, New Mexico 87131, USA*

³⁵*The Ohio State University, Columbus, Ohio 43210, USA*

³⁶*Okayama University, Okayama 700-8530, Japan*

³⁷*Osaka City University, Osaka 558-8585, Japan*

³⁸*University of Oxford, Oxford OX1 3RH, United Kingdom*

^{39a}*Istituto Nazionale di Fisica Nucleare, Sezione di Padova, I-35131 Padova, Italy*

^{39b}*University of Padova, I-35131 Padova, Italy*

⁴⁰*University of Pennsylvania, Philadelphia, Pennsylvania 19104, USA*

^{41a}*Istituto Nazionale di Fisica Nucleare Pisa, I-56127 Pisa, Italy*

^{41b}*University of Pisa, I-56127 Pisa, Italy*

^{41c}*University of Siena, I-56127 Pisa, Italy*

^{41d}*Scuola Normale Superiore, I-56127 Pisa, Italy*

- ^{41e}*INFN Pavia, I-27100 Pavia, Italy*
^{41f}*University of Pavia, I-27100 Pavia, Italy*
⁴²*University of Pittsburgh, Pittsburgh, Pennsylvania 15260, USA*
⁴³*Purdue University, West Lafayette, Indiana 47907, USA*
⁴⁴*University of Rochester, Rochester, New York 14627, USA*
⁴⁵*The Rockefeller University, New York, New York 10065, USA*
^{46a}*Istituto Nazionale di Fisica Nucleare, Sezione di Roma 1, I-00185 Roma, Italy*
^{46b}*Sapienza Università di Roma, I-00185 Roma, Italy*
⁴⁷*Mitchell Institute for Fundamental Physics and Astronomy, Texas A&M University, College Station, Texas 77843, USA*
^{48a}*Istituto Nazionale di Fisica Nucleare Trieste, I-33100 Udine, Italy*
^{48b}*Gruppo Collegato di Udine, I-33100 Udine, Italy*
^{48c}*University of Udine, I-33100 Udine, Italy*
^{48d}*University of Trieste, I-34127 Trieste, Italy*
⁴⁹*University of Tsukuba, Tsukuba, Ibaraki 305, Japan*
⁵⁰*Tufts University, Medford, Massachusetts 02155, USA*
⁵¹*University of Virginia, Charlottesville, Virginia 22906, USA*
⁵²*Waseda University, Tokyo 169, Japan*
⁵³*Wayne State University, Detroit, Michigan 48201, USA*
⁵⁴*University of Wisconsin, Madison, Wisconsin 53706, USA*
⁵⁵*Yale University, New Haven, Connecticut 06520, USA*
 (Received 30 September 2013; published 13 November 2013)

We present a measurement of the production cross section for ZW and ZZ boson pairs in final states with a pair of charged leptons, from the decay of a Z boson, and at least two jets, from the decay of a W or Z boson, using the full sample of proton-antiproton collisions recorded with the CDF II detector at the Tevatron, corresponding to 8.9 fb^{-1} of integrated luminosity. We increase the sensitivity to vector-boson decays into pairs of quarks using a neural-network discriminant that exploits the differences between the

^aDeceased.

^bVisitor from University of British Columbia, Vancouver, BC V6T 1Z1, Canada.

^cVisitor from Istituto Nazionale di Fisica Nucleare, Sezione di Cagliari, 09042 Monserrato (Cagliari), Italy.

^dVisitor from University of California Irvine, Irvine, CA 92697, USA.

^eVisitor from Institute of Physics, Academy of Sciences of the Czech Republic, 182-21, Czech Republic.

^fVisitor from CERN, CH-1211 Geneva, Switzerland.

^gVisitor from Cornell University, Ithaca, NY 14853, USA.

^hVisitor from University of Cyprus, Nicosia CY-1678, Cyprus.

ⁱVisitor from Office of Science, U.S. Department of Energy, Washington, DC 20585, USA.

^jVisitor from University College Dublin, Dublin 4, Ireland.

^kVisitor from ETH, 8092 Zürich, Switzerland.

^lVisitor from University of Fukui, Fukui City, Fukui Prefecture, Japan 910-0017.

^mVisitor from Universidad Iberoamericana, Lomas de Santa Fe, México, C.P. 01219, Distrito Federal.

ⁿVisitor from University of Iowa, Iowa City, IA 52242, USA.

^oVisitor from Kinki University, Higashi-Osaka City, Japan 577-8502.

^pVisitor from Kansas State University, Manhattan, KS 66506, USA.

^qVisitor from Brookhaven National Laboratory, Upton, NY 11973, USA.

^rVisitor from Queen Mary, University of London, London, E1 4NS, United Kingdom.

^sVisitor from University of Melbourne, Victoria 3010, Australia.

^tVisitor from Muons, Inc., Batavia, IL 60510, USA.

^uVisitor from Nagasaki Institute of Applied Science, Nagasaki 851-0193, Japan.

^vVisitor from National Research Nuclear University, Moscow 115409, Russia.

^wVisitor from Northwestern University, Evanston, IL 60208, USA.

^xVisitor from University of Notre Dame, Notre Dame, IN 46556, USA.

^yVisitor from Universidad de Oviedo, E-33007 Oviedo, Spain.

^zVisitor from CNRS-IN2P3, Paris, F-75205 France.

^{aa}Visitor from Universidad Tecnica Federico Santa Maria, 110v Valparaiso, Chile.

^{bb}Visitor from The University of Jordan, Amman 11942, Jordan.

^{cc}Visitor from Universite catholique de Louvain, 1348 Louvain-La-Neuve, Belgium.

^{dd}Visitor from University of Zürich, 8006 Zürich, Switzerland.

^{ee}Visitor from Massachusetts General Hospital, Boston, MA 02114 USA.

^{ff}Visitor from Harvard Medical School, Boston, MA 02114 USA.

^{gg}Visitor from Hampton University, Hampton, VA 23668, USA.

^{hh}Visitor from Los Alamos National Laboratory, Los Alamos, NM 87544, USA.

ⁱⁱVisitor from Università degli Studi di Napoli Federico I, I-80138 Napoli, Italy.

spatial spread of energy depositions and charged-particle momenta contained within the jet of particles originating from quarks and gluons. Additionally, we employ new jet energy corrections to Monte Carlo simulations that account for differences in the observed energy scales for quark and gluon jets. The number of signal events is extracted through a simultaneous fit to the dijet mass spectrum in three classes of events: events likely to contain jets with a heavy-quark decay, events likely to contain jets originating from light quarks, and events that fail these identification criteria. We determine the production cross section to be $\sigma_{ZW+ZZ} = 2.5^{+2.0}_{-1.0}$ pb (<6.1 pb at the 95% confidence level), consistent with the standard model prediction of 5.1 pb.

DOI: [10.1103/PhysRevD.88.092002](https://doi.org/10.1103/PhysRevD.88.092002)

PACS numbers: 14.70.-e, 12.15.-y

I. INTRODUCTION

The standard model (SM) offers precise predictions of production rates associated with self-interactions of the gauge bosons [1]. Differences between these predictions and measured diboson production cross sections may indicate the presence of non-SM physics [2,3], even specifically in hadronic final states [4]. Additionally, since hadronic final states in diboson production are similar to those from associated Higgs boson production ($p\bar{p} \rightarrow VH + X$ where $V = W, Z$), the analysis techniques used to measure diboson production in partially hadronic final states are relevant to searches for associated Higgs boson production.

Measurements of diboson production are typically difficult due to the small production cross sections of the order of 10 pb or less [1]. Furthermore, measurements of decay channels where one W or Z boson decays hadronically are particularly challenging at hadron colliders: although expected event yields are larger than those in purely leptonic decay channels due to the higher hadronic decay ($V \rightarrow q\bar{q}'$) branching ratio, the expected backgrounds from QCD multijet processes and $V + \text{jets}$ production are also much greater. Experiments at the Tevatron have previously measured the cross sections for pair production of gauge bosons in partially hadronic decay channels [5–8], but all of these measurements have included contributions from WW production, which has a higher cross section than that for combined ZW and ZZ production. Searches using identification of b -quark decays in the final states (b -tagging) to increase sensitivity to events with $Z \rightarrow b\bar{b}$ decays have been performed [9], but have not yet provided observations of ZV production in partially hadronic decay channels.

We present a study of ZV production from a final state with two leptons and at least two jets [10]. We require the two leptons to originate from the decay of a Z boson and search for associated $V \rightarrow q\bar{q}'$ decays by performing a fit to the dijet invariant mass (m_{jj}) spectrum of the two leading- E_T [11] jets. To maximize sensitivity to diboson production, we separate events into three channels: a heavy-flavor-tagged channel, largely sensitive to $ZZ \rightarrow \ell^+ \ell^- b\bar{b}$ decays; a light-flavor-tagged channel, which uses a new artificial-neural-network-based discriminant to preferentially select events with quark-like jets over gluon-like jets; and an untagged channel, which contains the remaining events that pass the event-selection requirements.

The final fit to the m_{jj} spectra is performed simultaneously across all three channels.

This paper is organized as follows: in Sec. II we briefly describe the CDF II detector; in Sec. III we describe the data sets and event-selection requirements that are used in the ZV search; in Sec. IV we show the derivation of new jet energy corrections to Monte Carlo simulations that account for differences in the observed energy scales of quark and gluon jets; in Sec. V we provide details of a new neural-network-based discriminant that identifies jets more likely to originate from quarks than from gluons; and, in Sec. VI, we describe the signal-extraction method, and report the results of the ZV search.

II. THE CDF II DETECTOR

The CDF II detector is described in detail elsewhere [12]. The detector is cylindrically symmetric around the Tevatron beam line. Tracking detectors are installed around the interaction point to reconstruct the trajectories of charged particles (tracks). The tracking systems are located within a superconducting solenoid that produces a 1.4 T magnetic field aligned with the $p\bar{p}$ beams. Around the outside of the solenoid, calorimeter modules arranged in a projective-tower geometry measure the energies of charged and neutral particles. Drift chambers outside the calorimeter are used to detect muons, which typically deposit little energy in the calorimeter.

The central outer tracker (COT) is a 3.1-m-long open-cell drift chamber that has 96 measurement layers in the region between 0.40 and 1.37 m from the beam axis, providing full track coverage in the pseudorapidity region $|\eta| < 1.0$. Sense wires are arranged in eight alternating axial and $\pm 2^\circ$ stereo “superlayers” with 12 wires each. The position resolution of a single drift-time measurement is about 140 μm . A five-layer double-sided silicon micro-strip detector (SVX) covers the region between 2.5 to 11 cm from the beam axis. Three separate SVX barrel modules along the beam line cover a length of 96 cm, approximately 90% of the luminous beam interaction region. Three of the five layers combine an r - ϕ measurement on one side and a 90° stereo measurement on the other, and the remaining two layers combine an r - ϕ measurement with a small-angle ($\pm 1.2^\circ$) stereo measurement. The typical silicon hit resolution is 11 μm . An intermediate silicon

layer at a radius of 22 cm from the beam line in the central region links tracks in the COT to hits in the SVX. The fiducial range of the silicon detector extends to pseudorapidity magnitude $|\eta| < 2.0$.

Calorimeter modules are located outside the central tracking volume and solenoid. The inner electromagnetic layers consist of lead sheets interspersed with scintillators, while the outer hadronic layers consist of scintillators sandwiched between steel sheets. The calorimeter is split between central barrel ($|\eta| < 1.1$) and forward end-plug ($1.1 < |\eta| < 3.6$) sections. Individual towers in the central barrel subtend 0.1 in $|\eta|$ and 15° in ϕ . The sizes of the towers in the end-plug calorimeter vary with $|\eta|$, subtending 0.1 in $|\eta|$ and 7.5° in ϕ at $|\eta| = 1.1$, and 0.5 in $|\eta|$ and 15° in ϕ at $|\eta| = 3.6$. The energy resolution in the electromagnetic calorimeters is $14\%/\sqrt{E_T}$ in the central barrel and $16\%/\sqrt{E} \oplus 1\%$ in the forward end-plug section, with the energies in units of GeV. The single-particle energy resolution in the hadronic calorimeters, measured using pions, ranges from $75\%/\sqrt{E}$ in the central barrel to $80\%/\sqrt{E} \oplus 5\%$ in the forward end-plug section, with the energies expressed in units of GeV.

The hadronization of quarks and gluons produced in the interaction leads to collimated groups of high-momentum particles called jets. These jets, along with photons and electrons, leave isolated energy deposits in contiguous groups of calorimeter towers, which can be summed together into an energy cluster. Electrons and photons are identified as isolated, mostly electromagnetic clusters, and quality requirements may be placed on the presence of a high- p_T track geometrically matched to the cluster to more accurately identify electrons. Jets are identified as electromagnetic and hadronic clusters with the combined electromagnetic (EM) fraction $E_{\text{EM}}/E_{\text{total}} = E_{\text{EM}}/(E_{\text{EM}} + E_{\text{had}}) < 0.9$, clustered using the JETCLU cone algorithm [13] with a fixed cone size of $\Delta R \equiv \sqrt{(\Delta\eta)^2 + (\Delta\phi)^2} = 0.4$.

Outside the calorimeters, drift chambers detect muons. A four-layer stack of planar drift chambers detect muons with $p_T > 1.4$ GeV/c, and another four layers of drift chambers behind 60 cm of steel detect muons with $p_T > 2.0$ GeV/c. Both systems cover a region of $|\eta| < 0.6$, though they have different structure and their geometrical coverages do not overlap exactly. Muons in the region between $0.6 < |\eta| < 1.0$ pass through at least four drift layers arranged within a conic section outside of the central calorimeter. Muons are identified as either COT tracks that extrapolate to hits in the muon detectors, or as isolated tracks unmatched to hits in the muon detectors that satisfy tighter tracking-quality requirements and extrapolate to calorimeter energy depositions consistent with a minimum ionizing particle.

III. DATA SET AND EVENT SELECTION

We analyze the full data set of $p\bar{p}$ collisions collected by the CDF II detector. We require events to be collected from

periods when the tracking systems, calorimeters, and muon detectors were all functioning properly, corresponding to an integrated luminosity of 8.9 fb^{-1} . Events are selected via a number of high- E_T electron and high- p_T muon online event-selection requirements (triggers). The majority of these triggers require at least one electron (muon) with $E_T > 18$ GeV ($p_T > 18$ GeV/c). We require the events to contain two electrons (muons) with $E_T > 20$ GeV ($p_T > 20$ GeV/c) and determine the trigger selection and event reconstruction efficiencies by comparing the number of data and simulated $Z \rightarrow \ell\ell$ events containing exactly one jet with $E_T > 20$ GeV.

For the final analysis, we select events with at least two leptons, and two or more jets. In the unlikely case that more than two charged leptons are reconstructed, we select the two leptons with the highest p_T . In addition to the p_T requirements on the leptons, we require leptons associated with well-reconstructed tracks (central electrons, $|\eta| < 1$, and all muons) to be of opposite charge, and a reconstructed dilepton invariant mass, $m_{\ell\ell}$, consistent with the mass of the Z boson, $76 < m_{\ell\ell} < 106 \text{ GeV}/c^2$. We require both leading- E_T jets to have $E_T > 25$ GeV and $|\eta| < 2.0$, and to be spatially separated from the reconstructed leptons ($\Delta R > 0.4$). Additionally, the two jets must be separated by $\Delta R > 0.7$. Finally, as the final state should contain no particles that are not reconstructed in the detector, we also require that the missing transverse energy, $\cancel{E}_T$ [14], is less than 20 GeV.

After this selection, three major sources of background contribute. The dominant background comes from the production of a Z boson, decaying to an e^+e^- or $\mu^+\mu^-$ pair, in association with two jets. Simulated events generated using ALPGEN [15], interfaced with PYTHIA [16] for showering, are used to estimate this background. The production cross sections for $Z + b\bar{b}$ processes are normalized to experimental measurements [17].

Another significant background results from jets misidentified as leptons. The contributions from these lepton *fakes* are estimated using data-driven methods. For muons, we use events with same-sign muon pairs (rather than opposite-sign) that otherwise satisfy all event-selection requirements. For electrons, we derive a misidentification rate representing the likelihood for a jet to be misidentified as an electron, as a function of jet E_T and η , using jet-triggered data with minimal contributions from events with electrons. This rate is then applied to all possible electron-jet pairs in events from the high- p_T electron data set, where the jet is then treated as a second electron, and the event-selection requirements are otherwise applied normally.

While the requirement to have low $\cancel{E}_T$ reduces its total contribution, top-quark-pair production, where each top quark decays into a leptonic final state ($t\bar{t} \rightarrow W^+bW^-\bar{b} \rightarrow \ell^+\nu_\ell b\ell^-\bar{\nu}_\ell\bar{b}$), contributes events to the final event sample, especially in the heavy-flavor-tagged channel. We estimate $t\bar{t}$ contributions using PYTHIA with $\sigma_{t\bar{t}} = 7.5 \text{ pb}$ and

$m_t = 172.5 \text{ GeV}/c^2$. Finally, ZW and ZZ signal events are also modeled using PYTHIA.

IV. JET-ENERGY CALIBRATION

The energies of jets, measured in the calorimeter, are corrected to account for a number of effects that distort the true jet energy. These effects include changes in calorimeter performance as a function of $|\eta|$ and time, contributions from multiple $p\bar{p}$ interactions per beam crossing (pileup), contributions from the other partons in the interacting proton and antiproton (underlying event), the nonlinear response of the calorimeter, and energy radiated outside of the jet cone. These jet-energy-scale (JES) corrections are described in detail in Ref. [18].

These energy corrections, however, do not attempt to account for potential differences in the modeled calorimeter response to jets originating from quarks and gluons. For example, the largest corrections modify the energy scale of the jets to more accurately match that of the initial parton energies and their resulting particle jets, and are derived using PYTHIA [16] dijet Monte Carlo (MC) simulations, but independently of the initiating type of parton. Differences in the response of gluon and quark jets between simulation and data lead to differences in the measured energies for these jets that are not covered by the previously assigned systematic uncertainties on the JES [19].

We derive a data-driven correction for the response to quark and gluon jets in simulated events using two independent samples of jets with different compositions of quark- and gluon-initiated jets. In these samples, we derive a correction to the jet energy by balancing the transverse energy of the jets against particles of known transverse momentum. We use events where a jet recoils against a high- E_T photon—a sample rich in quark jets (based on simulations modeled using PYTHIA)—and utilize the significant number of $Z \rightarrow \ell^+\ell^- + \text{jet}$ events available in the full CDF data set, which have a larger fraction of gluon jets. The quark and gluon content of these two samples differ due to the difference in mass between the Z boson and the photon: because the Z boson mass is higher, the initial partons of the production process typically carry a higher fraction of the momentum of the proton than those involved in the production of high-energy photons. This leads to a difference in the quark and gluon content of these samples.

To derive a correction, we construct the *balance* of the jet with these more accurately measured reference particles,

$$K_{Z,\gamma} = (E_T^{\text{jet}}/p_T^{Z,\gamma}) - 1. \quad (1)$$

For unbiased measurements of the jet energy, the balance will equal zero. Samples of jets with nonzero balance could be corrected with a jet-energy correction factor of $1/(K_{Z,\gamma} + 1)$. However, rather than derive independent JES corrections for quark and gluon jets in data and simulation, we compare the balance in data and simulation

and derive an additional correction to be applied to simulated jets, dependent upon whether these jets are matched to quarks or gluons. The correction to simulated quark-jet energies is $+1.4 \pm 2.7\%$, while the correction to gluon-jet energies is much larger: $-7.9 \pm 4.4\%$.

A. Data set and event selection

The data set and event selection for the Z -jet balancing sample largely follow those described in Sec. III. We require two leptons consistent with resulting from the decay of a Z boson and exactly one jet with $E_T > 15 \text{ GeV}$ and no other jets with (uncorrected) $E_T > 3 \text{ GeV}$ within $|\eta| < 2.4$. Additionally, we ensure that the Z boson and jet are azimuthally opposite (back-to-back) by requiring their azimuthal separation to exceed 2.8 radians, and we require that $p_T^Z > 10 \text{ GeV}/c$.

For the γ -jet balancing sample, we closely mirror the selection requirements described in Ref. [18]. We use events collected with an isolated-central-photon trigger over the same period of time as that of the high- p_T lepton samples. We compare these data to PYTHIA simulations of both $\gamma + \text{jet}$ production as well as dijet production, which also contributes to the γ -jet balancing sample.

To match the selection requirements of the isolated-central-photon trigger, we require $E_T^\gamma > 27 \text{ GeV}$ and $0.2 < |\eta_\gamma| < 0.6$ in both data and MC simulation. To decrease the contribution from dijet production, where a jet mimics the photon selection, we require the energy in the calorimeter and momentum in the tracking system contained within a cone of $R = 0.4$ around the photon to be less than 1 GeV and 2 GeV/ c , respectively. As in the Z -jet balancing sample, we require events to have exactly one measured jet with $E_T > 15 \text{ GeV}$ and no other jets with (uncorrected) $E_T > 3 \text{ GeV}$ within $|\eta| < 2.4$. We also demand the $\Delta\phi$ between the jet and photon to be larger than 3.0 radians. We further reduce contamination by vetoing events with more than one reconstructed interaction point, and by removing events with $\cancel{E}_T/E_T^\gamma > 0.8$, which likely contain activity from cosmic rays.

B. Determination of corrections

We derive separate corrections for the quark- and gluon-jet energy scales in data and simulation using the Z -jet and γ -jet balancing samples in the following way. The balances of the Z -jet and γ -jet systems (K_Z and K_γ , respectively) can be expressed as linear combinations of independent quark- and gluon-balance variables (K_q and K_g , respectively), weighted by the sample-specific quark and gluon fractions ($F_{Z,\gamma}^{q,g}$),

$$K_Z = F_Z^q K_q + F_Z^g K_g = F_Z^q K_q + (1 - F_Z^q) K_g, \quad (2)$$

$$K_\gamma = F_\gamma^q K_q + F_\gamma^g K_g = F_\gamma^q K_q + (1 - F_\gamma^q) K_g. \quad (3)$$

Rewriting these expressions by solving for K_q and K_g , we find

$$K_q = \frac{1}{F_\gamma^q - F_Z^q} [(1 - F_Z^q)K_\gamma - (1 - F_\gamma^q)K_Z], \quad (4)$$

$$K_g = \frac{1}{F_\gamma^q - F_Z^q} [F_\gamma^q K_Z - F_Z^q K_\gamma]. \quad (5)$$

These expressions apply separately to experimental data and simulated data, yielding a different balance in data and Monte Carlo simulation (K^{data} and K^{MC} , respectively) and may include a dependence on the energy of the jet, with $F_{Z,\gamma}^q = F_{Z,\gamma}^q(E_T^{\text{jet}})$ and consequently $K = K(E_T^{\text{jet}})$.

In order to solve for K_q and K_g , we require knowledge of $K_{Z,\gamma}$ and $F_{Z,\gamma}^q$. We extract a value of $K_{Z,\gamma}$ as a function of E_T^{jet} by constructing the balancing distribution, as defined in Eq. (1), in ranges of E_T^{jet} , and fit the core of the distribution with a Gaussian shape. We perform these fits separately in data and simulation and use the mean and uncertainty on the mean of the fitted Gaussian shape as the value of $K_{Z,\gamma}(E_T^{\text{jet}})$ and its uncertainty. We use this estimation of the most probable value in order to avoid effects from a small fraction of highly mismeasured jets, which may strongly bias the mean and median of the distribution.

The distributions of K_Z and K_γ in data and simulated data are shown in Fig. 1. Not only are jets measured poorly (the balance does not average to zero), but in the Z-jet balancing sample, largely dominated by gluon jets, there is significant disagreement between the correction factors for simulated jets and those in data. We do not see a similar disagreement in the γ -jet balancing sample, indicating that the simulation models the behavior of the jets in this quark-jet dominated sample accurately.

We determine $F_{Z,\gamma}^q$ from simulation by matching jets to their originating partons. In the γ -jet balancing sample, the quark fraction is about 85% at $E_T^{\text{jet}} \approx 30$ GeV and drops to about 71% at $E_T^{\text{jet}} \approx 70$ GeV. In the Z-jet balancing sample, these fractions are roughly 38% and 49%, respectively, in the same E_T^{jet} regions. In data, it is not possible to

directly match jets to their originating parton, and we must therefore rely on simulation to extract values of $F_{Z,\gamma}^q(E_T^{\text{jet}})$. Because we are trying to correct for discrepancies in the reconstruction of quark and gluon jets between data and simulation, we cannot simply use the simulation-derived $F_{Z,\gamma}^q$ values from each jet E_T bin. Rather, we parametrize $F_{Z,\gamma}^q$ from simulation as a function of $p_T^{Z/\gamma}$ and determine $F_{Z,\gamma}^q$ in each jet E_T bin of the data based on the p_T^Z or p_T^γ distribution in that bin.

Using Eqs. (4) and (5), we construct distributions of K_q and K_g as functions of the jet E_T , as shown in Fig. 2. We see good agreement between data and simulation in K_q but poorer agreement in K_g , where the data correction appears consistently lower than that for simulation. This suggests that the MC simulation is systematically overestimating gluon jet energies, relative to the data.

Using the distributions of K_q and K_g , we determine the corrections that need to be applied to simulated jets in order to best match the energy scale of the data. These MC simulation corrections are defined as $(K_q^{\text{data}} + 1)/(K_q^{\text{MC}} + 1)$ for quark jets and $(K_g^{\text{data}} + 1)/(K_g^{\text{MC}} + 1)$ for gluon jets, as shown in Fig. 3. Due to the photon trigger used to select the γ -jet balancing sample, reliable balancing information is not available for jets with transverse energies smaller than 27.5 GeV in that sample, limiting the range of applicability of the corrections. Since we are interested in jets of energies extending down to 20 GeV, we extrapolate the quark-jet energy correction derived for jets with $E_T > 27.5$ GeV to lower jet energies and use the Z-jet balancing sample to extract a gluon correction assuming this extrapolated quark correction.

The quark and gluon corrections' dependence on jet energy are accurately modeled by a constant for jets with $E_T > 15$ GeV. Quark jet energies in simulation should be increased by approximately 1.4% to more accurately match the data, while gluon jet energies should be decreased by approximately 7.9%.

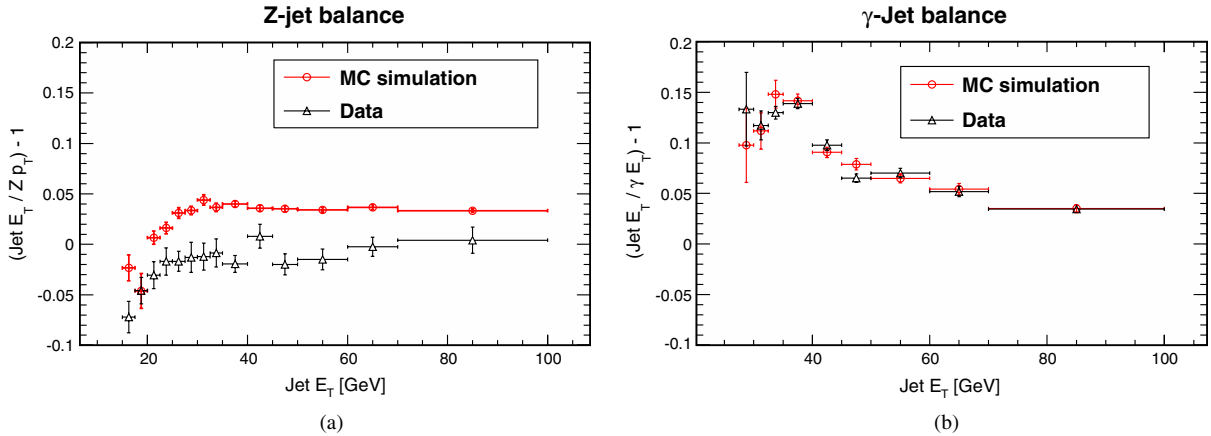


FIG. 1 (color online). Balancing distributions, (a) K_Z and (b) K_γ , in data and MC simulation as a function of E_T^{jet} . The uncertainties include solely the contribution from the fluctuations in the mean of a Gaussian fit to the balancing distributions in bins of E_T^{jet} .

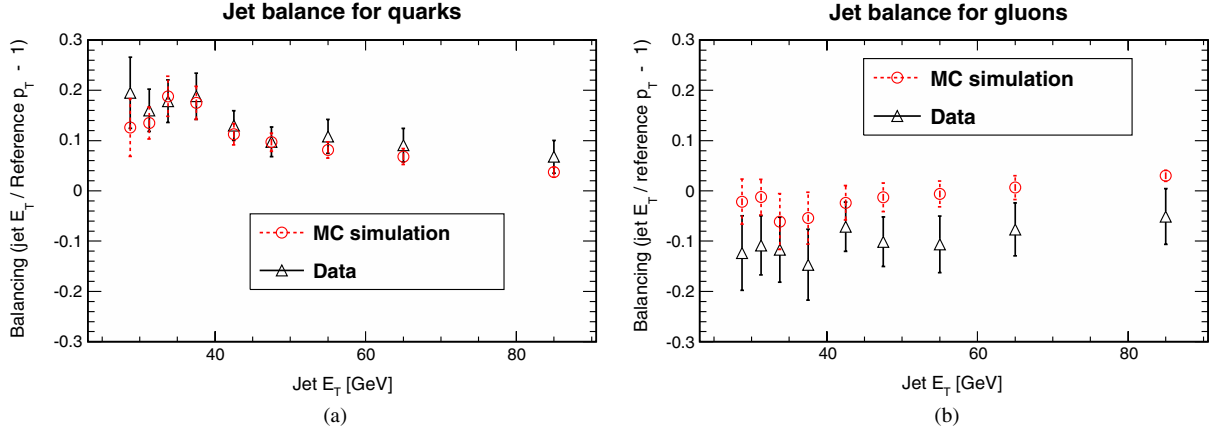


FIG. 2 (color online). Derived balancing variable for (a) quark jets, K_q , and (b) gluon jets, K_g , in data and MC simulation as a function of E_T^{jet} . The uncertainties on each point are from the uncertainties from the mean of the Gaussian fit and the uncertainties on the quark fractions, added in quadrature.

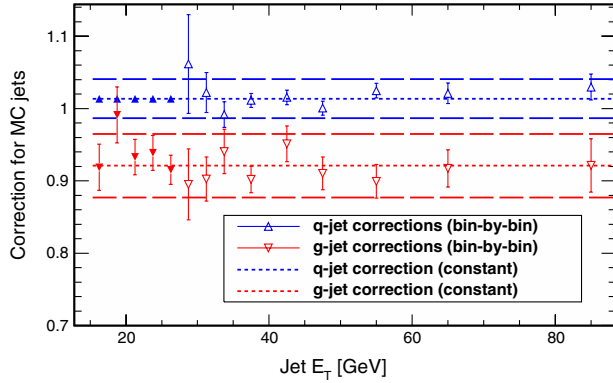


FIG. 3 (color online). Derived correction for simulated quark jets and gluon jets as a function of E_T^{jet} . The open triangles represent corrections derived using both γ -jet and Z-jet balancing samples, while the filled triangles represent the assumed uniform correction for quarks and the corresponding correction for gluons calculated from the Z-jet balancing sample alone. The uncertainties shown are statistical only. The short dashed lines are the fits of the correction to a constant across jet E_T (as opposed to the corrections in each bin of E_T^{jet}), and the long dashed lines represent the total systematic uncertainty bands on that constant correction, further described in Sec. IV C.

C. Uncertainties on simulated jet energy corrections

We consider the following sources of uncertainty on the correction factors.

- (1) *Statistical uncertainty*: We use the standard deviation of the corrections to the simulation for each jet E_T bin around the assumed uniform correction function. This is an uncertainty of $\pm 2.0\%$ for quark-jet energies, and $\pm 2.5\%$ for gluon-jet energies (the uncertainties on quark- and gluon-jet energy corrections are anticorrelated).
- (2) F_Z^q : We compare the distribution of a quark-gluon discriminant parameter (described in detail in Sec. V) in data and simulation, and fit the data distribution

using quark and gluon templates from simulation. We take the average deviation of the value determined for this quark-gluon discriminant from the nominal MC simulation value as a systematic uncertainty on F^q , constant across jet E_T . This uncertainty is approximately 10%. Here we vary the calculated quark fraction in the data Z-jet balancing sample by $\pm 10\%$ and recalculate the corrections for quark and gluon jets. This translates to an uncertainty of $\pm 0.6\%$ for quark-jet energies and $\pm 2.1\%$ for gluon-jet energies.

- (3) F_γ^q : We follow a similar procedure of fitting the quark-gluon discriminant parameter in the γ -jet sample and obtain a similar uncertainty of $\pm 10\%$ on the quark fraction. This translates to an uncertainty of $\pm 1.8\%$ for quark-jet energies, and $\pm 2.7\%$ for gluon-jet energies.
- (4) *Low E_T extrapolation*: We check the dependence of the gluon-jet energy corrections on the assumed quark-jet corrections for low- E_T jets by varying the quark-jet E_T for these jets by $\pm 2\%$. We see a small change in the gluon-energy corrections, which translates to $\pm 0.4\%$ of the gluon-jet energy.
- (5) *Number of interaction vertices dependence*: The γ -jet balancing sample incorporates a requirement on the number of reconstructed interaction vertices to reduce contamination from pileup. The Z-jet balancing sample does not have such a requirement, due to a much smaller background contribution and in order to retain as many events as possible. We check for any bias in the corrections resulting from the effect of this requirement by looking for any shift in the corrections when the requirement is placed on the Z-jet balancing sample. We see a change in the quark-jet energies of $\pm 0.2\%$, and the gluon-jet energies of $\pm 1.2\%$.

The uncertainties are summarized in Table I. Because the corrections shift the energy response in the simulation to more accurately match the data, the quark-jet and gluon-jet

TABLE I. Summary of the additional jet energy corrections applied to simulated jets and their uncertainties.

Jet-energy correction		Quark jets 1.014	Gluon jets 0.921
Uncertainty	Statistical	0.020	0.025
	$F_{Q\text{-jet}}^{Z\text{-jet}}$	0.006	0.021
	$F_{Q\text{-jet}}^{\eta\text{-jet}}$	0.018	0.027
	Low E_T extrapolation		0.004
	N_{vert} difference	0.002	0.012
Total uncertainty		± 0.027	∓ 0.044

energy correction uncertainties are anticorrelated. The uncertainties are similar in magnitude to the default CDF jet energy scale uncertainties [18].

V. ARTIFICIAL NEURAL NETWORK QUARK-TO-GLUON DISCRIMINANT

In this analysis, we search for two high- p_T leptons from the decay of a Z boson and two jets from a $W \rightarrow q\bar{q}'$ or $Z \rightarrow q\bar{q}$ decay. Thus, the two signal jets are quark jets. Conversely, the dominant background, two jets produced in association with a $Z \rightarrow \ell^+ \ell^-$ decay, contains a significant fraction of gluon jets (of the order of 50%). The ability to separate quark jets from gluon jets is therefore useful for increasing sensitivity to ZW and ZZ production.

For a given energy, gluon jets, due to their higher color charge, tend to feature a higher particle multiplicity and be spatially broader in the detector than light-quark (u , d , and s) jets. We attempt to quantify the spatial spread of jets using a collection of artificial neural networks (NNs) trained to separate gluon jets from light-flavor quark jets. We refer to the output of the final NN as the jet quark-to-gluon discriminant value (or jet QG value). We calibrate the response of the final NN in MC simulation to match the response in data based on a $W \rightarrow \ell\nu + 1$ jet event sample. The tagging efficiency and mistag rate associated to a requirement on the jet QG value are obtained from two independent event samples: $W \rightarrow \ell\nu + 2$ jets events, which are representative of the Z + jets background; and $t\bar{t} \rightarrow b\bar{b}\ell\nu q\bar{q}'$ events, which contain two non-heavy-flavor jets from the hadronic decay of a W boson, similar to the diboson signal.

A. Jet QG definition

A total of three NNs contribute to the final QG discriminant. The initial two networks separate quark and gluon jets by exploiting distinctive features in the distribution of energies reconstructed in calorimeter towers and momenta of charged particles reconstructed in the tracking chambers. Thus, every jet is assigned a *tower* NN value and *track* NN value, which are the outputs of these networks. These two NN values are then used as inputs to a third NN.

Each of the NNs is trained using simulated samples of jets matched to either a light-flavor quark or gluon with

$p_T > 20$ GeV/c within $\Delta R = 0.4$ of the center of the jet and further requiring that no additional partons with transverse momenta exceeding 8 GeV/c are present within $\Delta R = 0.7$. These jets are selected from a $Z \rightarrow \mu^+ \mu^- + 2$ parton ALPGEN sample, interfaced with PYTHIA showering. Each NN is a feed-forward multilayer perceptron with a hyperbolic-tangent-like response function [20]. The networks are trained on 100 000 quark and gluon jets and tested for biases in overtraining on samples containing 500 000 quark and gluon jets. Gluon-jet distributions are reweighted to match the E_T and η distributions of the quark jets to remove any discrimination power coming solely from these variables.

For each jet we obtain a list of the calorimeter towers within a cone of $\Delta R = 0.7$. Each tower has a location coordinate, (η, ϕ) , and energy deposition E associated with it. We construct a distribution of the distance, ΔR , between all pairs of towers within the jet and weight each tower pair by its relevance in terms of energy to obtain a distribution that characterizes the spatial spread of the energy within each jet. The weight is given by

$$\frac{E_i E_j}{0.5[(\Sigma E)^2 - \Sigma E^2]},$$

where E_i and E_j are the energies detected in the two towers of the pair, ΣE is the sum of the energy in all towers within a cone of $\Delta R = 0.7$ around the jet, and ΣE^2 is the sum of the square of the energies of each tower in that same cone. The denominator is chosen to normalize the sum of the weights of all tower pairs to unity. We sample this distribution in 56 intervals (bins) of size $\Delta R_{\text{bin}} = 0.025$ for $0.0 < \Delta R < 1.4$, where the contents of the first three bins are empty due to the segmentation of the calorimeter. Typical distributions of the weighted ΔR between tower pairs for quark and gluon jets are shown in Fig. 4, using a larger bin size. The outputs of the tower NN for quark and gluon jets using the training and testing samples are shown in Fig. 5.

We follow a similar prescription using tracks within a cone of $\Delta R = 0.7$ around each jet, using the tracks' locations in (η, ϕ) (with respect to the primary vertex) and momenta p to obtain a distribution of the distance between pairs of tracks (in ΔR), with each pair weighted by the momentum carried by that pair, or

$$\frac{p_i p_j}{0.5[(\Sigma p)^2 - \Sigma p^2]},$$

where p_i and p_j are the magnitude of the momenta of the charged particles in the pair, Σp is the scalar sum of the momenta carried by all charged particles within a cone of $\Delta R = 0.7$ around the jet, and Σp^2 is the sum of the square of the momenta of each charged particle within that same cone. We require all contributing charged particles to come from the primary vertex and have $p_T > 0.4$ GeV/c. We split the ΔR between track pairs distribution into the same

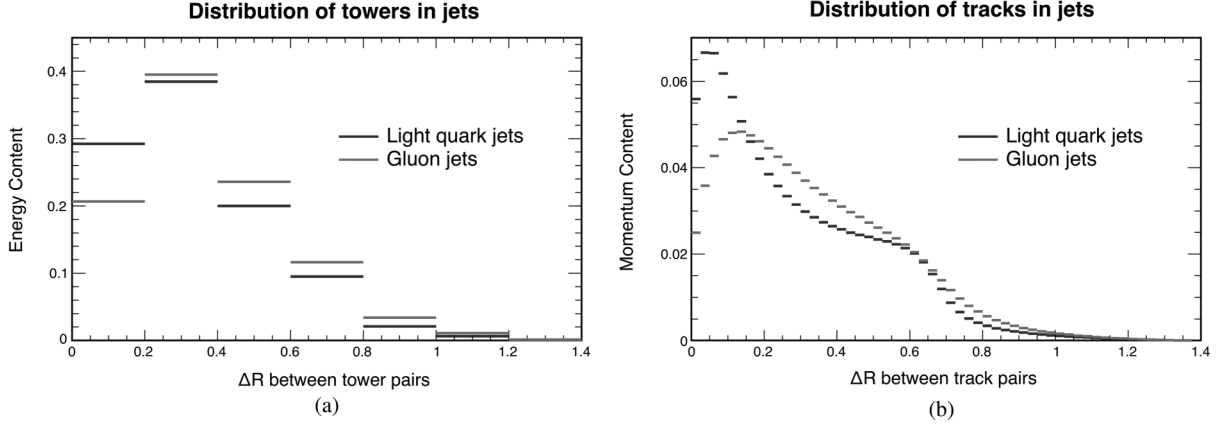


FIG. 4. Typical distributions of energy (momentum) content of jets as a function of the ΔR (a) between pairs of towers and (b) between pairs of tracks in light-flavor quark and gluon jets in simulation.

56 intervals (bins) as used in the tower NN, and the content of each bin is used as an input into the track NN.

Typical distributions of ΔR between track pairs for quark and gluon jets are shown in Fig. 4. Light-flavor quark jets tend to peak at low ΔR , indicating that they are rather collimated, while gluon jets tend to have a higher mean-valued ΔR distribution. The bin contents of these ΔR distributions are used as inputs into NNs that discriminate between quark and gluon jets.

The outputs of the track NN for quark and gluon jets using the training and testing samples are shown in Fig. 5. Higher NN scores indicate jets that are more quark-like. We see good performance in both the tower and track NNs. The cusps in the track NN distribution are associated to jets containing only two charged-particle tracks located inside a cone of $\Delta R = 0.7$, and thus have only one nonzero bin in their distributions of ΔR between track pairs.

The final NN uses both the tower and track NN values as inputs, along with other jet variables that provide discrimination power between quark jets and gluons: the ratio of ΣE associated to towers within a cone of

$\Delta R = 0.4$ to ΣE associated to towers within a cone of $\Delta R = 0.7$; the ratio of Σp associated to charged particles within a cone of $\Delta R = 0.4$ to Σp associated to charged particles within a cone of $\Delta R = 0.7$; the number of contributing towers with nonzero energy in cones of $\Delta R = 0.4$ and 0.7 ; the number of contributing charged-particle tracks in cones of $\Delta R = 0.4$ and 0.7 ; and the jet EM fraction. Additionally, other variables that affect the shape of the ΔR distributions, independent of whether the jet originates from a quark or gluon, are included: the jet E_T , the jet η , and the number of reconstructed interaction vertices in the event. The output of this final NN is shown in Fig. 6 for light-flavor quark and gluon jets using the training and testing samples. In simulated jets, we see significant separation between quark and gluon jets using this discriminant.

B. Jet QG calibration and performance

The response of the NN quark-to-gluon discriminant may differ between data and MC simulation, especially since uncorrected tower energies are used in the construction of

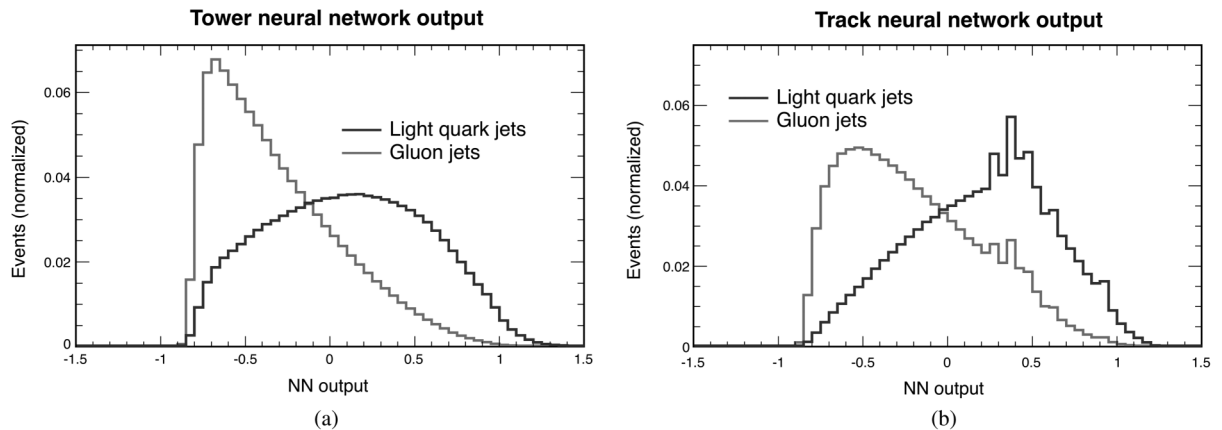


FIG. 5. Distributions of the outputs of the NNs processing (a) tower information and (b) track information in light-flavor quark and gluon jets in simulation. Higher NN scores indicate jets that are more quark-like.

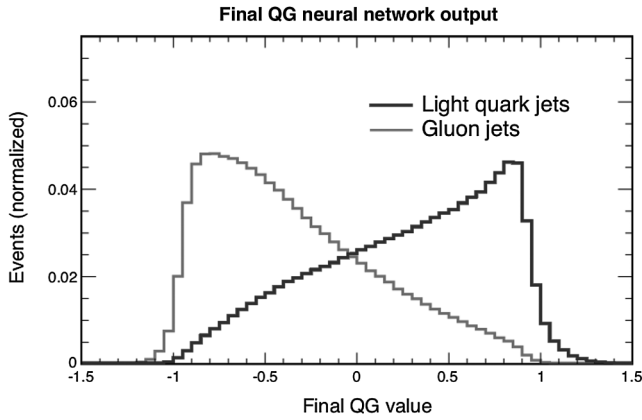


FIG. 6. Distribution of the output of the final NN for light-flavor quark and gluon jets in simulation. Higher NN scores indicate jets that are more quark-like.

the tower NN. Since the signal and most backgrounds are modeled with simulated data, we calibrate the simulation response to match the response in data. We use a control region of independent events with a jet composition similar to that of the final state, $W \rightarrow \ell \nu + 1$ jet events. We then validate the calibration and establish uncertainties on the modeling using control samples of data with features similar to the signal and samples enriched in the dominant backgrounds: $t\bar{t}$ decays in lepton + jets final states and $W \rightarrow \ell \nu + 2$ jet events, respectively.

To form the $W + 1$ jet calibration sample, we choose data collected with the standard high- E_T (p_T) central electron (muon) triggers and select events with exactly one central ($|\eta| < 1.0$) electron (muon) with E_T (p_T) > 20 GeV/ c . To select events consistent with a $W \rightarrow \ell \nu$ decay, we also require a significant missing transverse energy, $\cancel{E}_T > 25$ GeV, and a reconstructed transverse mass [21] consistent with leptonic W -boson decays, $m_T > 25$ GeV/ c^2 . To further suppress any contributions from multijet events where a jet mimics the lepton + $\cancel{E}_T$ signature, we require that the $\cancel{E}_T$ is not aligned with any reconstructed jet [$\Delta\phi(\cancel{E}_T, \text{jet}) > 0.2$ radians] and that the $\cancel{E}_T$ significance—a dimensionless quantity comparing the observed $\cancel{E}_T$ against the energy resolution of jets, soft unclustered particles, and the event topology (see Ref. [5])—be larger than four (one) for events with electrons (muons). We also require that the events in this calibration sample have exactly one jet with $E_T > 20$ GeV and $|\eta| < 2.0$.

We consider various processes that contribute to this sample, listed in Table II, and model them using a combination of the PYTHIA, ALPGEN, and MADGRAPH [22] event generators interfaced with PYTHIA for showering. The dominant contribution is $W \rightarrow \ell \nu$ production in association with one jet, which is modeled using ALPGEN. As we are largely concerned with the agreement in shapes between data and simulation, we scale the simulation distributions to match the data. Additionally, we reweight the

TABLE II. Number of events in the $W + 2$ jets and $t\bar{t}$ lepton + jets region, showing only the uncertainties assigned on the normalization of each sample. The $W +$ jets samples are rescaled to match data in the number of events observed after the $W +$ jets selection. The distinction between b and q jets in the $t\bar{t}$ sample refers to the lower two b ness jets: events where both jets are matched to non- b -quark jets are labeled q jets, while if one of the jets is matched to a b jet, the event is labeled b jets.

	$W +$ jets selection	$t\bar{t}$ selection
$W +$ jets	21500 ± 2200	38.7 ± 3.9
$W + b$ jets	940 ± 380	13.8 ± 5.5
$Z +$ jets	1250 ± 130	3.1 ± 0.3
$Z + b$ jets	86 ± 34	1.4 ± 0.6
$WW + WZ$	1386 ± 83	5.9 ± 0.4
Single top quark	767 ± 77	19.6 ± 2.0
$t\bar{t}$	1378 ± 83	469 ± 28
$t\bar{t}$ (b jets)		108 ± 7
$t\bar{t}$ (q jets)		361 ± 22
Total expected	27300 ± 2200	551 ± 30
Data	27 319	579

simulation to match the jet E_T and η distributions in data to remove these variables as a possible causes for mismodeling of the jet QG value.

We observe poor modeling of the tower NN values, where the jets in data appear more gluon-like than those in simulated events. The fact that jets in data appear more spatially spread than jets in simulation is consistent with the observed differences in jet energy scales for data and simulation, described in Sec. IV: the fraction of the jet energy contained within a cone of $\Delta R = 0.4$ is higher in simulated gluon jets than in gluon jets in data. We correct for these discrepancies by applying a linear shift to the tower NN values observed in simulation in order to match with data using the $W + 1$ jet sample. We apply different linear shifts for jets in the central and plug calorimeters, and for jets in events with different levels of pileup. We apply further corrections to the response of the final NN to more accurately match the correlations of these calibrated tower-NN values with other jet quantities: the number of towers in the jets and the ratio of ΣE in a cone of $\Delta R = 0.4$ to ΣE in a cone of $\Delta R = 0.7$. The modeling is more accurate in the track NN than in the tower NN, though we still introduce a similar linear shift in simulated track NN values to more accurately match data. The calibrated variables are input directly into the final NN, without retraining the network.

We further validate the response of the jet QG value by comparing data and MC simulation in a $W \rightarrow \ell \nu + 2$ jets event sample and in an event sample dominated by $t\bar{t}$ production where two quark jets originate from the hadronic decay of a W boson. Table III summarizes the requirements used to select these two samples: the $W + 2$ jet sample is similar to the previously described $W + 1$ jet sample, except for modified jet selections to

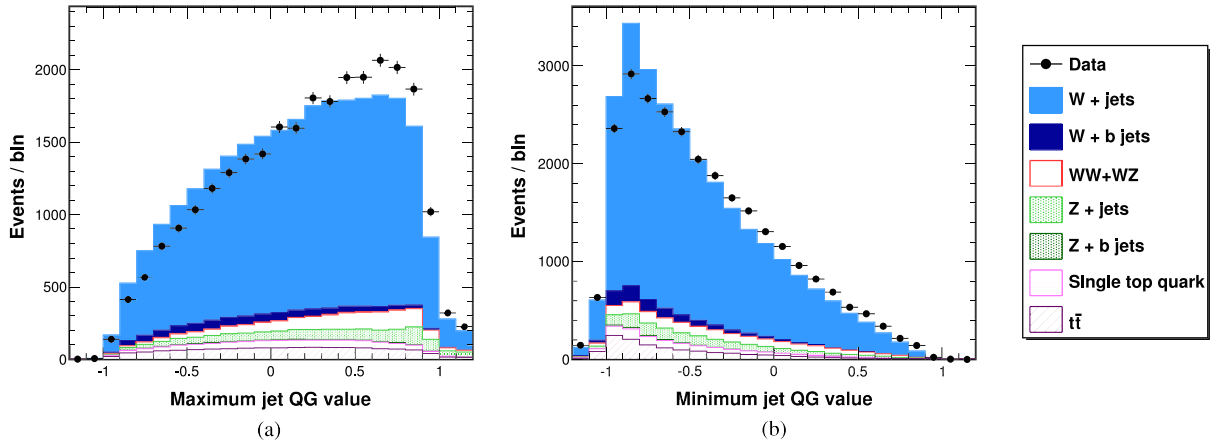
TABLE III. Summary of event-selection requirements for the $t\bar{t}$ lepton + jets selection and the $W + 2$ jets selection, used to understand the modeling of events in the QG discriminant. Requirements in the center are shared requirements in the two samples.

$W + \text{jets selection}$		$t\bar{t}$ selection
	Central e or μ , $p_T > 20 \text{ GeV}/c$ $\cancel{E}_T > 25 \text{ GeV}$	
$\Delta\phi(\cancel{E}_T, \text{nearest jet}) > 0.4 \text{ rad}$ $\cancel{E}_T \text{ sig} > 4$ (e only) $m_T(W) > 25 \text{ GeV}/c^2$ (e only)		$\Delta\phi(\cancel{E}_T, \text{nearest jet}) > 0.2 \text{ rad}$
$N_{\text{jets}}(E_T > 20 \text{ GeV}) = 2$		Sum $E_T > 300 \text{ GeV}$ $N_{\text{jets}}(E_T > 20 \text{ GeV}) = 4$
1st/2nd jet $E_T > 25 \text{ GeV}$		2nd highest $b\text{ness}$ jet $b\text{ness} > -0.5$ 2nd highest $b\text{ness}$ jets $E_T > 20 \text{ GeV}$ 2 lowest $b\text{ness}$ jets $E_T > 25 \text{ GeV}$
	Jet $ \eta < 2.0$ ΔR between jets > 0.7	

match those used in the signal region of the $ZV \rightarrow \ell\ell jj$ search. The $t\bar{t}$ selection eschews the $\cancel{E}_T$ significance and m_T requirements, used to reduce multijet backgrounds, in favor of requirements on the minimum scalar sum of transverse quantities (jets' E_T , $\cancel{E}_T$, and the charged lepton's p_T) in the event, which is effective in removing both multijet and $W + \text{jets}$ backgrounds. Because we are interested in selecting the two jets in the $t\bar{t}$ candidate events that come from the decay of a W boson, as opposed to the b jets produced in the $t \rightarrow Wb$ decays, we make use of the jet- $b\text{ness}$ tagger [23]. This multivariate b -jet identification algorithm exploits properties of individual charged-particle tracks within a jet, looking at properties characteristic of charged particles originating from B -hadron decays. The final score, the output of a NN discriminant that ranges between -1 and 1 , is called the jet $b\text{ness}$, where higher scores identify jets that are more likely to originate from B -hadron decays. We classify the two jets with the highest $b\text{ness}$ scores in the event as the two b jets, and the remaining two jets as those resulting from a $W \rightarrow q\bar{q}'$ decay.

Because we are looking for jet QG shape differences between data and simulation that induce acceptance uncertainties when a requirement on the jet QG value is applied, we scale the number of $W + \text{jet}$ events in simulation to match the yield observed in the $W + 2$ jets data. The number of events in each sample is shown in Table II. The distributions of the maximum and minimum QG values of the two jets are shown in Figs. 7 and 8. We see fairly good modeling in the $t\bar{t}$ sample, but poorer modeling in the $W + 2$ jet sample where, even after calibrations, the jets in simulation appear more gluon-like than the jets in data. We account for this remaining discrepancy between data and simulation below.

We enhance the sensitivity to the signal when forming a *light-flavor-tagged* channel where the minimum jet QG value is greater than 0.0 . We determine a probability for a quark jet to meet this requirement (efficiency), and for a gluon jet to be misidentified as a quark jet (mistag rate), with the $t\bar{t}$ and $W + 2$ jet samples. The efficiency measured in data, ϵ_D is a function of the QG requirement, q , and may be expressed as

FIG. 7 (color online). Distribution of (a) the maximum and (b) minimum jet QG values of the two jets in the $W + 2$ jet sample.

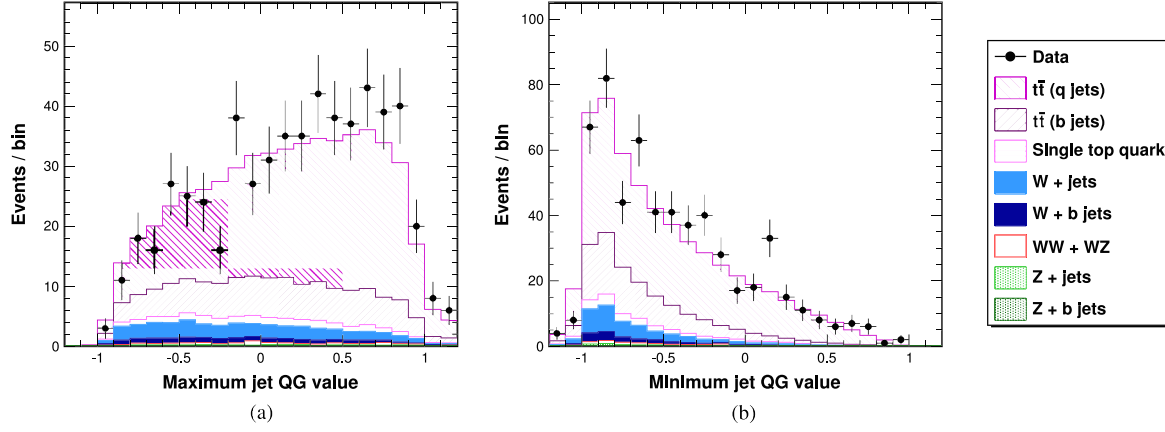


FIG. 8 (color online). Distribution of (a) the maximum and (b) minimum jet QG values of the two jets in the $t\bar{t}$ sample. The distinction between q and b jets refers to the lower two b ness jets: events where both jets are matched to non- b quark jets are labeled q jets, while if one of the jets is matched to a b jet, the event is labeled b jets.

$$e_D(q) = \frac{e_{\text{raw}}(q) - s_m(q)m_{\text{MC}}(q)f_g}{1 - f_g}, \quad (6)$$

where e_{raw} is the fraction of data events passing the QG requirement; m_{MC} is the mistag rate for gluons, as measured in simulation; s_m is a scale factor on the mistag rate in simulated jets to match the mistag rate measured in data; and f_g is the fraction of gluon jets in the sample. We can write a similar expression for the mistag rate from

$$m_D(q) = \frac{m_{\text{raw}}(q) - s_e(q)e_{\text{MC}}(q)f_q}{1 - f_q}, \quad (7)$$

where m_{raw} is the fraction of data events meeting the QG requirements; e_{MC} is the efficiency for quarks to pass the requirement, as measured in simulation; s_e is a scale factor on the efficiency in simulated jets to match the mistag rate measured in data; and f_q is the fraction of quark jets in the sample. Squared uncertainties on these quantities may be expressed as

$$\sigma_e^2(q) = \frac{1}{(1 - f_g)^2} \left[\frac{e_{\text{raw}}(1 - e_{\text{raw}})}{N_D} + (\sigma_m f_g)^2 \right] + \sum_X \frac{\sigma_X^2}{[N_{\text{MC}}(1 - f_g)]^2} \times [(e + s_m m_{\text{MC}})(f_g - f_g^X) + f_q^X(e_{\text{MC}} - e_X)]^2, \quad (8)$$

where N_D and N_{MC} are the number of data and simulated events, respectively, and where the X represents the various simulated subsamples, and σ_m is the uncertainty on the mistag rate, which may be expressed in an analogous fashion. The uncertainty includes a statistical uncertainty on the data, uncertainties on the mistag rate and efficiency, and uncertainties on the relative difference in the contributions from the simulation. We take the uncertainties on the normalizations of the $t\bar{t}$, single top quark, diboson, $V + \text{jets}$, and $V + b$ jets to be 6%, 10%, 6%, 10%, and 40%, respectively, based on the uncertainties in their production cross sections.

We measure the efficiency in the $t\bar{t}$ sample, where the fraction of gluon jets is small, and measure the mistag rate in the $W + 2$ jets sample, where the gluon fraction is much larger and similar to the fraction in the $Z + 2$ jets signal region. The efficiency, mistag rate, and their uncertainties are determined using an iterative procedure. We first calculate the mistag rate in data assuming that the efficiency in data equals the efficiency in simulation. We then calculate the efficiency in data assuming that value for the mistag rate and proceed to update the mistag rate assuming the new value for the efficiency from data. We observe rapid convergence on robust values for the efficiency and mistag rate. Table IV shows the efficiency and mistag rate for the given requirement of minimum QG > 0.0 , measured in both data and MC simulation. The simulation underestimates the rate

TABLE IV. Efficiency and mistag rates for the chosen jet QG requirements, as evaluated in data and MC simulation, along with the necessary change in the jet QG threshold for the simulation to model the proper rates and the uncertainties on them.

	MC	Data	MC revised jet QG requirement (-1σ , nom., $+1\sigma$)
Efficiency	0.241	0.295 ± 0.034	$(-0.0325, -0.09, -0.14)$
Mistag rate	0.088	0.087 ± 0.027	$(0.09, -0.0175, -0.11)$

for quark jets to meet the jet QG requirement, while correctly predicting the observed mistag rate.

We implement a correction to the MC simulation by varying the requirement on the minimum QG value in order to reproduce the efficiency and mistag rate observed in data. The uncertainties on these quantities are also obtained by varying the jet QG requirement. The alternate thresholds used for simulated quark and gluon jets are listed in Table IV.

VI. SIGNAL EXTRACTION AND RESULTS

We extract the number of signal events using a binned χ^2 -minimization fit to data, using the techniques described in Ref. [24]. We create histogram templates for both signal and background samples. The templates, along with the uncertainties we assign to their normalization in the fit procedure, are listed below.

- (1) ZV signal: We allow the normalization of the signal template to float unconstrained in the fit. We assume that each signal process contributes proportionally to its predicted SM cross section: 3.6 pb for ZW and 1.5 pb for ZZ [1].
- (2) $Z + \text{jets}$: This is the largest background. We allow its normalization to float in the fit, unconstrained.
- (3) $Z + b$ jets: We constrain the normalization of this significant background to be within $\pm 40\%$ of its nominal value.
- (4) $t\bar{t}$: We use a production cross section of $\sigma_{t\bar{t}} = 7.5$ pb, and assign an uncertainty of 6.5% to the normalization of this template, based on the theoretical cross-section uncertainty [25].
- (5) Misidentified leptons: We use the method described in Sec. III to construct templates for the contribution from jets mimicking one or two leptons. We assign an uncertainty of 50% to the misidentification rate,

based on studies using different trigger thresholds for the jet data used to obtain these rates.

We perform a simultaneous fit to data using independent templates for each of three channels. For events passing the basic signal selection requirements described in Sec. III, we first construct a heavy-flavor tag (HF-tag) channel composed of events passing a minimum jet b -ness requirement (jet b -ness > 0), using the jet- b -ness tagger [23]. For events failing this requirement, we then select events passing the minimum jet-QG value requirement described in Sec. V to form a light-flavor tag (LF-tag) channel. Events failing this requirement are then placed in the third, untagged channel, which has a lower signal fraction than the two tagged channels, but still includes a significant amount of signal due to the tight tagging requirements. The predicted and observed numbers of events are shown in Table V.

Additional systematic uncertainties on both the normalization and shapes of the templates used in the fit are also considered. We estimate shape and normalization uncertainties due to mismodeling between data and MC simulation in the jet energy scale (as described in Sec. IV C) and the jet energy resolution, the modeling of the tagging variables, and the lepton energy scale and resolution. Additional shape uncertainties on the $Z + \text{jets}$ backgrounds are considered by increasing and decreasing the renormalization and factorization scale, Q^2 , from the default value in the simulation of $m_Z^2 + p_{T,Z}^2$. We also consider the effect on the shape of the dijet invariant mass when increasing or decreasing initial- and final-state QCD radiation (ISR/FSR) in the ZV signal model. These systematic uncertainties, along with the normalization constraints described above, are treated as nuisance parameters in the fit, and are included in the χ^2 -minimization procedure [24]. They are summarized in Table VI.

TABLE V. Predicted and observed numbers of events in the final event selection, where the numbers of events are rounded to the appropriate significant figures given the uncertainties. The uncertainties incorporate all systematic uncertainties summarized in Table VI and include an additional 10% uncertainty on the normalization of $Z + \text{jets}$ events and a 6% uncertainty on the normalization of $ZW + ZZ$ events, from the theoretical uncertainties on the production cross sections for those processes.

	All events	Heavy-flavor-tagged	Light-flavor-tagged	Untagged
$Z + \text{jets}$	8700 ± 1100	93 ± 14	1520 ± 310	7100 ± 970
$Z + b$ jets	710 ± 300	111 ± 48	55 ± 26	550 ± 230
$t\bar{t}$	9.2 ± 0.9	3.3 ± 0.4	0.7 ± 0.1	5.1 ± 0.6
Misidentified leptons	330 ± 170	4.8 ± 2.4	41 ± 20	280 ± 140
Predicted background	9700 ± 1200	212 ± 55	1620 ± 330	7900 ± 1100
$ZW + ZZ$	313 ± 29	12.8 ± 1.6	89 ± 12	212 ± 22
$N_{\text{signal}}/N_{\text{background}}$	3.2%	6.0%	5.2%	2.6%
Total predicted events	10000 ± 1300	225 ± 55	1710 ± 330	8100 ± 1100
Data events	9 846	172	1 724	7 950

TABLE VI. Summary of the systematic uncertainties considered in the fit of the dijet mass distribution. Uncertainties that change both the shape and rate of templates used in the fit are treated in a correlated fashion.

Source	Channel	ZV	Z + jets	Z + b jets	$t\bar{t}$	Mis-ID leptons
Cross section/normalization	All	Unconstrained	Unconstrained	$\pm 40\%$	$\pm 6.5\%$	$\pm 50\%$
Jet-energy resolution	HF-tag	$\pm 0.8\%$	$\pm 0.3\%$	$\pm 1.0\%$	$\pm 0.2\%$	
	LF-tag	$\pm 1.0\%$	$\pm 0.7\%$	$\pm 1.5\%$	$\pm 6.2\%$	
	Untagged	$\pm 0.6\%$	$\pm 0.9\%$	$\pm 0.7\%$	$\pm 1.1\%$	
Jet energy scale	HF-tag	$\pm 4.0\%$	$\pm 4.4\%$	$\pm 3.8\%$	$\pm 4.0\%$	
	LF-tag	$\pm 1.5\%$	$\pm 0.3\%$	$\pm 0.6\%$	$\pm 3.0\%$	
	Untagged	$\pm 1.9\%$	$\pm 5.7\%$	$\pm 3.8\%$	$\pm 1.9\%$	
Q^2	All	none	Shape only	Shape only	none	
ISR/FSR	All	Shape only	none	none	none	
b ness tag	HF-tag	$\pm 7.8\%$	$\pm 7.8\%$	$\pm 9.2\%$	$\pm 7.6\%$	
	LF-tag	$\pm 0.2\%$	$\pm 0.0\%$	$\pm 1.2\%$	$\pm 2.8\%$	
	Untagged	$\pm 0.4\%$	$\pm 0.1\%$	$\pm 1.8\%$	$\pm 4.5\%$	
QG tag	LF-tag	$\pm 10\%$	$\pm 16\%$	$\pm 2.0\%$	$\pm 15\%$	
	Untagged	$\pm 4.3\%$	$\pm 3.5\%$	$\pm 2.0\%$	$\pm 2.0\%$	
Lepton-energy scale	All	$\pm 0.5\%$	$\pm 0.5\%$	$\pm 0.5\%$	$\pm 1.5\%$	
Lepton-energy resolution	All	$\pm 0.1\%$	$\pm 0.1\%$	$\pm 0.0\%$	$\pm 2.7\%$	

Figure 9 shows the dijet mass distributions in data with the fit results overlaid. Table VII shows the number of events of each class determined by the fit. We fit for approximately 50% of the expected

signal normalization and observe good agreement between data and simulation in the final fit for each of the three channels, with a total $\chi^2/\text{d.o.f.} = 59.8/55$.

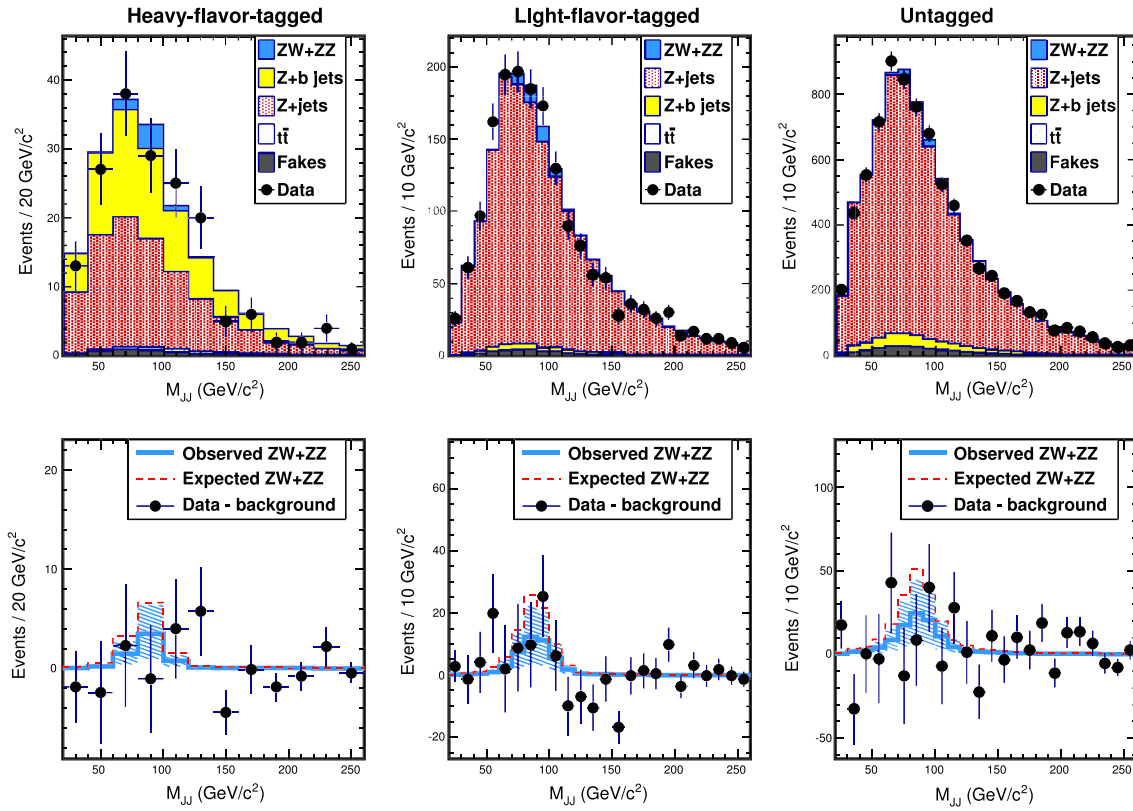


FIG. 9 (color online). Invariant dijet mass distributions with fit results overlaid for the ZW + ZZ process in the dilepton + dijet selection in the heavy-flavor-tagged channel (left panels), light-flavor-tagged channel (center panels), and untagged channel (right panels). The top row shows the output from the fit compared to the data, while the bottom row shows the background subtracted from data, compared to the expected (dashed line) and fitted (solid line, with uncertainties in bands) signal contributions.

TABLE VII. Number of events in each class from the best fit to the data.

Process	$N_{\text{events, HF-tag}}$	$N_{\text{events, LF-tag}}$	$N_{\text{events, Untagged}}$
$Z + \text{jets}$	91.9 ± 8.3	1605 ± 50	7200 ± 600
$Z + b \text{ jets}$	71 ± 14	37 ± 10	360 ± 100
$t\bar{t}$	3.2 ± 0.4	0.7 ± 0.1	5.3 ± 0.4
Misidentified leptons	4.6 ± 2.3	39 ± 20	270 ± 140
Total background	171 ± 14	1681 ± 36	7840 ± 600
$ZW + ZZ$	6.3 ± 4.4	45 ± 30	106 ± 72
Total events	177 ± 14	1726 ± 40	7940 ± 610
Data events	172	1 724	7 950

We do not see significant evidence for $ZW + ZZ$ production in this decay channel. Hence, we set upper limits on the production cross section using likelihood-ratio ordering [26], where we analyze the distribution of observed cross sections in pseudoexperiments generated with a variety of scale factors on the input-signal cross section. When generating pseudoexperiments, we consider additional systematic uncertainties that affect the acceptance, assigning a 2% uncertainty on the signal template from limited knowledge of the parton distribution functions, and 2.5% and 6% uncertainties due to the uncertainties on the lepton-scale factor and integrated luminosity, respectively. The set of input cross sections in the pseudoexperiments range from 0.0 to 2.9 times the expected cross section, with a step size of 0.1.

Figure 10 shows the resulting confidence band. Using the 1σ bands, we determine $\sigma(p\bar{p} \rightarrow ZW + ZZ) = 2.5^{+2.0}_{-1.0}$ pb, compared to the standard model prediction of $\sigma_{\text{SM}} = 5.1$ pb. We do not exclude the no-signal hypothesis, and establish a limit of $\sigma_{ZW+ZZ} < 6.1$ pb ($1.25 \times \sigma_{\text{SM}}$) at the 95% C.L.

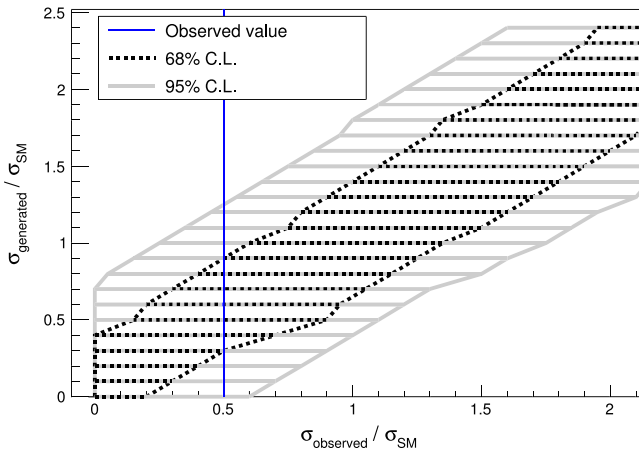


FIG. 10 (color online). Confidence bands showing the expected range of observed cross sections as a function of the true cross section, with 68% C.L. (black dashed region) and 95% C.L. (solid gray region).

VII. SUMMARY

In this paper we described a search for ZW - and ZZ -boson pair production from a final state with two charged, high-transverse-momentum electrons or muons and two hadronic jets. We increased the sensitivity by tagging events with jets likely originating from heavy- and light-flavor quarks and classifying them in separate analysis channels using neural-network-based taggers. These taggers benefited from the large sample of events containing top quarks collected by CDF, allowing a data-driven estimate of the efficiency and mistag rates for jets passing tagging requirements. We also improved the modeling of the Monte Carlo simulations, especially those that describe the $Z + \text{jets}$ background, by deriving and incorporating improved energy corrections for simulated jets to more accurately reproduce the phenomenology of jets originating from quarks and gluons in the data.

Using the full CDF Run II proton-antiproton collisions data set, which corresponds to an integrated luminosity of 8.9 fb^{-1} , we fit for the normalization of $ZW, ZZ \rightarrow \ell^+ \ell^- + q\bar{q}'$ events using the dijet invariant mass distribution. We incorporated many of the systematic uncertainties associated with the modeling of signal and background processes as nuisance parameters in the dijet mass fit. We measured a cross section of $\sigma_{ZW+ZZ} = 2.5^{+2.0}_{-1.0}$ pb, which is nonzero at the 1.75σ level of significance. We also obtained a limit on the cross section of $\sigma_{ZW+ZZ} < 6.1$ pb at the 95% C.L.

ACKNOWLEDGMENTS

We thank the Fermilab staff and the technical staffs of the participating institutions for their vital contributions. This work was supported by the U.S. Department of Energy and National Science Foundation; the Italian Istituto Nazionale di Fisica Nucleare; the Ministry of Education, Culture, Sports, Science and Technology of Japan; the Natural Sciences and Engineering Research Council of Canada; the National Science Council of the Republic of China; the Swiss National Science Foundation; the A.P. Sloan Foundation; the Bundesministerium für Bildung und

Forschung, Germany; the Korean World Class University Program, the National Research Foundation of Korea; the Science and Technology Facilities Council and the Royal Society, United Kingdom; the Russian Foundation for Basic Research; the Ministerio de Ciencia

e Innovación, and Programa Consolider-Ingenio 2010, Spain; the Slovak R&D Agency; the Academy of Finland; the Australian Research Council (ARC); and the EU community Marie Curie Fellowship Contract No. 302103.

-
- [1] J. M. Campbell and R. K. Ellis, *Phys. Rev. D* **60**, 113006 (1999).
- [2] K. Hagiwara, R. Peccei, D. Zeppenfeld, and K. Hikasa, *Nucl. Phys.* **B282**, 253 (1987).
- [3] M. Kober, B. Koch, and M. Bleicher, *Phys. Rev. D* **76**, 125001 (2007).
- [4] E. J. Eichten, K. Lane, and A. Martin, *Phys. Rev. Lett.* **106**, 251803 (2011).
- [5] T. Aaltonen *et al.* (CDF Collaboration), *Phys. Rev. Lett.* **103**, 091803 (2009).
- [6] T. Aaltonen *et al.* (CDF Collaboration), *Phys. Rev. Lett.* **104**, 101801 (2010).
- [7] T. Aaltonen *et al.* (CDF Collaboration), *Phys. Rev. Lett.* **106**, 171801 (2011).
- [8] V. M. Abazov *et al.* (D0 Collaboration), *Phys. Rev. Lett.* **108**, 181803 (2012).
- [9] T. Aaltonen *et al.* (CDF Collaboration), *Phys. Rev. D* **85**, 012002 (2012).
- [10] W. Ketchum, Ph.D. thesis, University of Chicago, 2012, Report No. FERMILAB-THESIS-2012-36.
- [11] The coordinate system used to describe the detector has its origin located at the detector's center. The proton beam direction is defined as the $+z$ direction. The polar angle, θ , is measured from the origin with respect to the z axis, and ϕ is the azimuthal angle. Pseudorapidity, transverse energy, and transverse momentum are defined as $\eta = -\ln \tan(\theta/2)$, $E_T = E \sin \theta$, and $p_T = p \sin \theta$, respectively. The rectangular coordinates x and y point radially outward and vertically upward from the Tevatron ring, respectively.
- [12] A. Abulencia *et al.* (CDF Collaboration), *J. Phys. G* **34**, 2457 (2007).
- [13] F. Abe *et al.* (CDF Collaboration), *Phys. Rev. D* **45**, 1448 (1992).
- [14] We define the missing transverse momentum $\cancel{E}_T \equiv -\sum_i E_T^i \mathbf{n}_i$, where $\mathbf{n}_i$ is the unit vector in the azimuthal plane that points from the beamline to the i th calorimeter tower. We call the magnitude of this vector the missing transverse energy, $\cancel{E}_T$.
- [15] M. L. Mangano, M. Moretti, F. Piccinini, R. Pittau, and A. D. Polosa, *J. High Energy Phys.* **07** (2003) 001.
- [16] T. Sjostrand, S. Mrenna, and P. Z. Skands, *J. High Energy Phys.* **05** (2006) 026.
- [17] T. Aaltonen *et al.* (CDF Collaboration), *Phys. Rev. D* **79**, 052008 (2009).
- [18] A. Bhatti *et al.*, *Nucl. Instrum. Methods Phys. Res., Sect. A* **566**, 375 (2006).
- [19] A. Paramonov, F. Canelli, M. D'Onofrio, H. Frisch, and S. Mrenna, *Nucl. Instrum. Methods Phys. Res., Sect. A* **622**, 698 (2010).
- [20] A. Hoecker, P. Speckmayer, J. Stelzer, J. Therhaag, E. von Toerne, and H. Voss, *Proc. Sci., ACAT2007* (2007) 040.
- [21] We define the transverse mass as $m_T = \sqrt{2p_T^{\text{lepton}} \cancel{E}_T (1 - \cos \Delta\phi_{\ell, \cancel{E}_T})}$, where $\Delta\phi_{\ell, \cancel{E}_T}$ is the difference in ϕ in the direction of the lepton $\vec{p}_T$ and $\vec{\cancel{E}}_T$ vectors.
- [22] J. Alwall, P. Demin, S. de Visscher, R. Frederix, M. Herquet, F. Maltoni, T. Plehn, D. L. Rainwater, and T. Stelzer, *J. High Energy Phys.* **09** (2007) 028.
- [23] J. Freeman, J. Lewis, W. Ketchum, S. Poprocki, A. Pronko, V. Rusu, and P. Wittich, *Nucl. Instrum. Methods Phys. Res., Sect. A* **663**, 37 (2012).
- [24] T. Aaltonen *et al.* (CDF Collaboration), *Phys. Rev. D* **82**, 112005 (2010).
- [25] U. Langenfeld, S. Moch, and P. Uwer, *Phys. Rev. D* **80**, 054009 (2009).
- [26] G. J. Feldman and R. D. Cousins, *Phys. Rev. D* **57**, 3873 (1998).