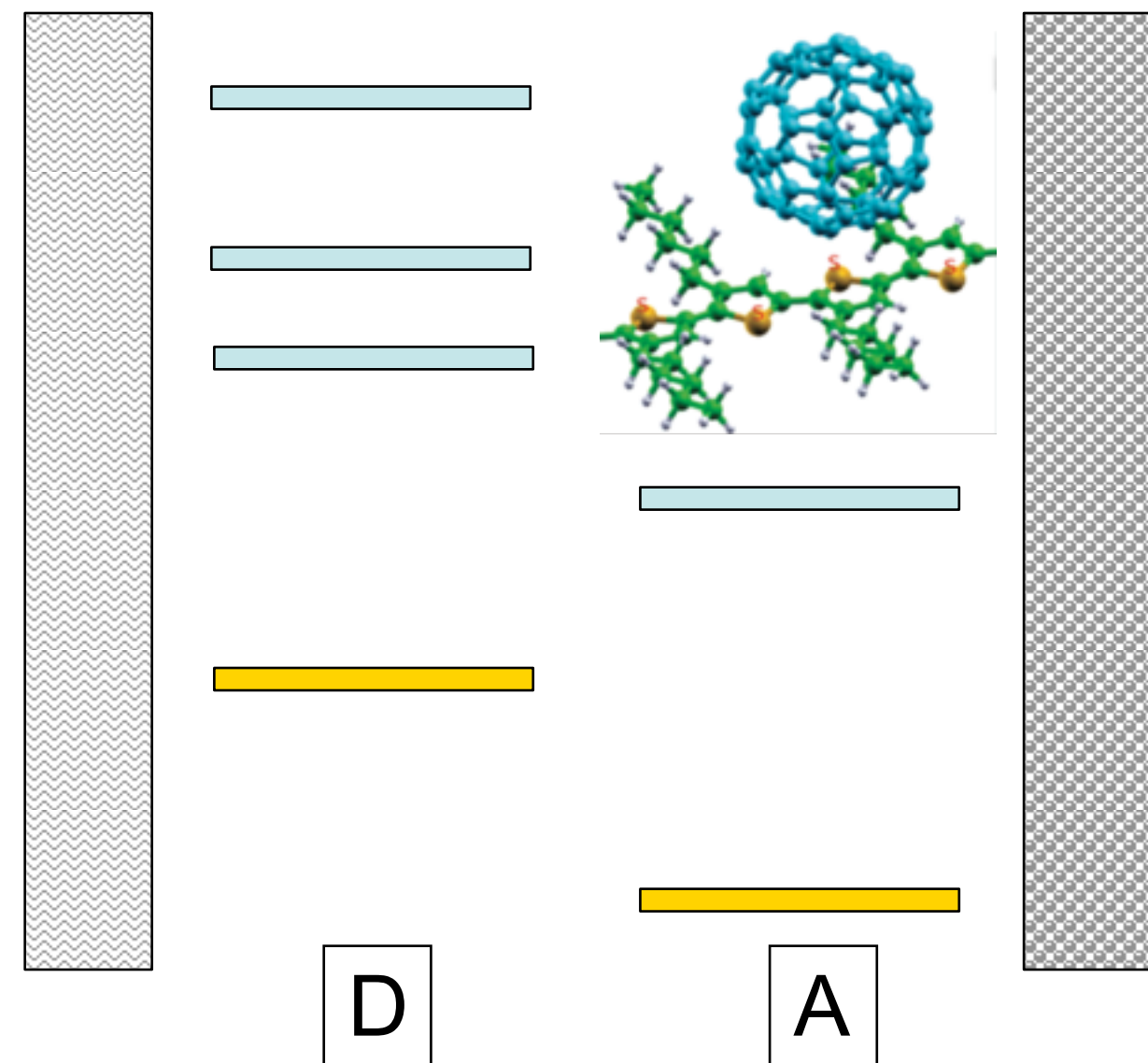
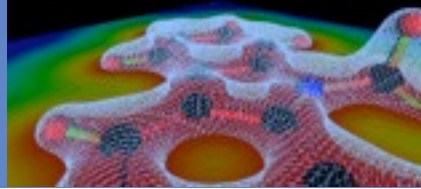


Faster G_0W_0 implementation for more accurate material design

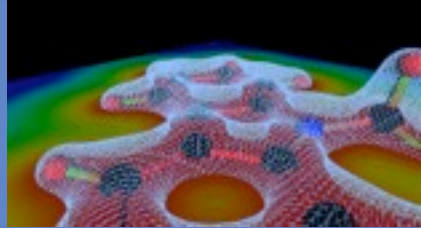
By Jonathan Laflamme Janssen

Département de physique, Université de Montréal

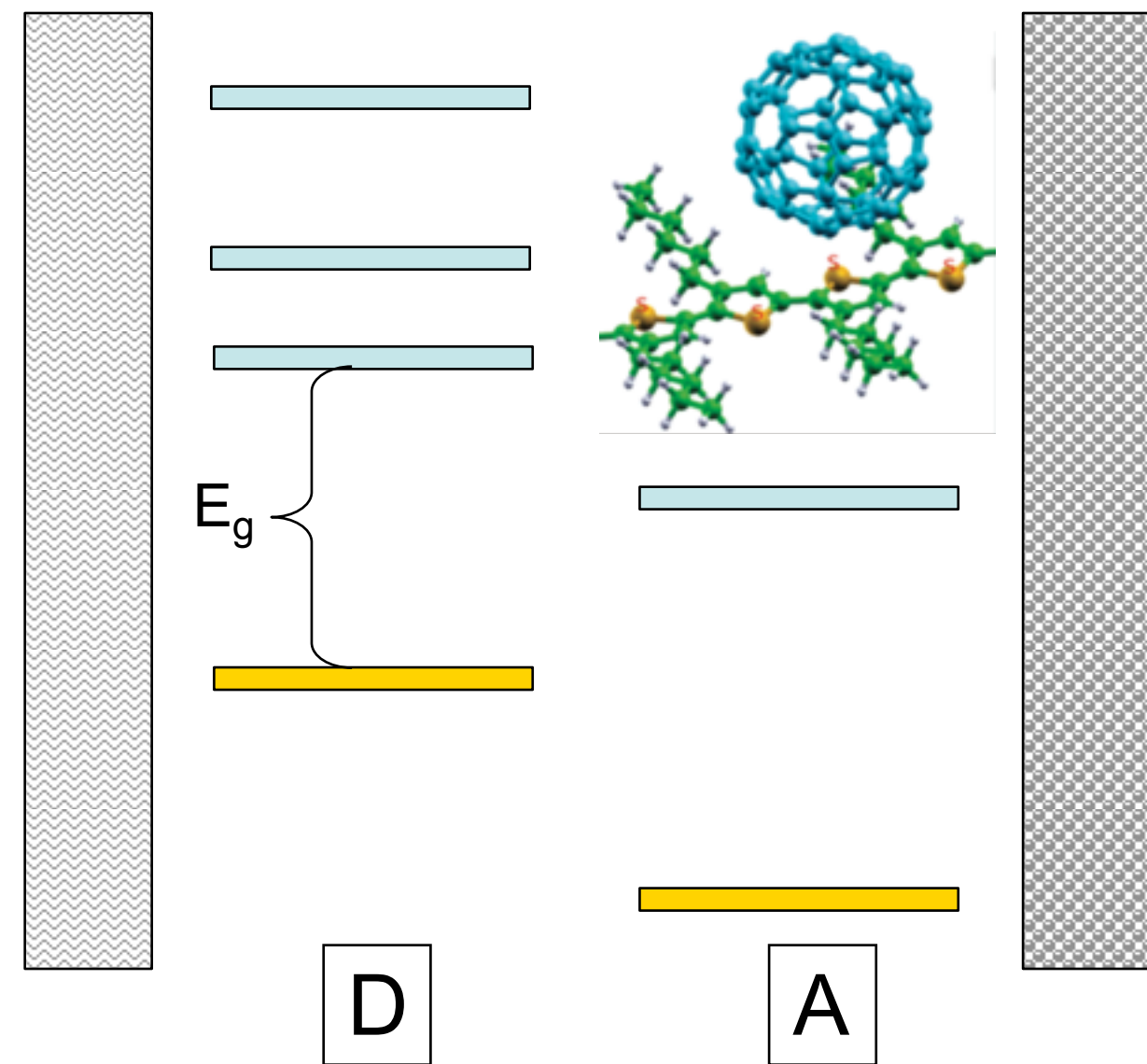
Design of polymers



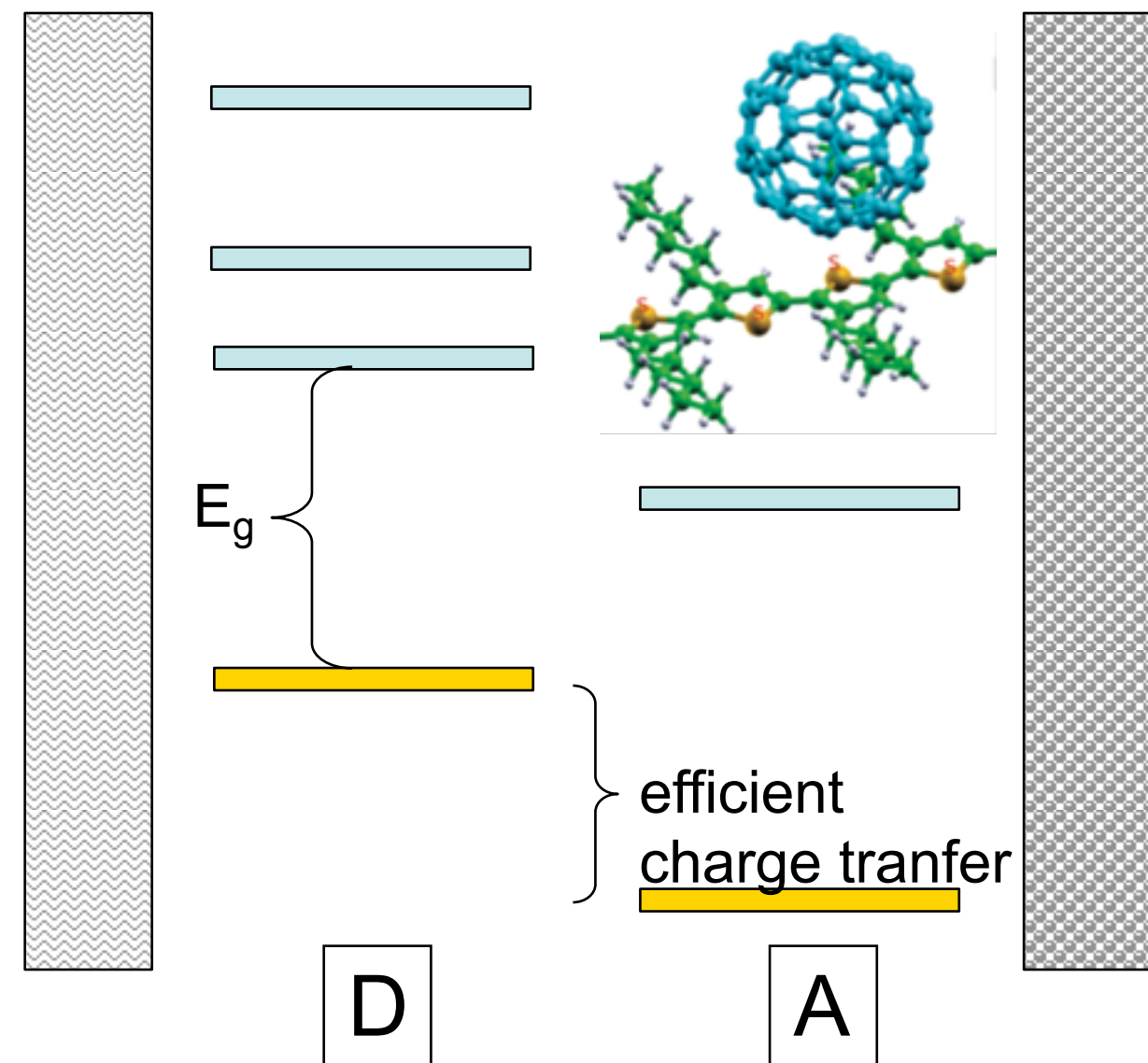
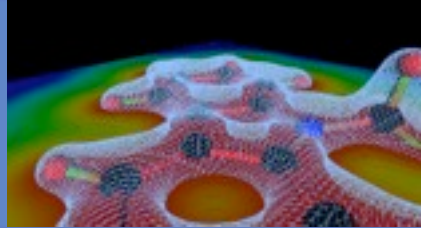
Design of polymers



- Gap: optimal for solar spectrum

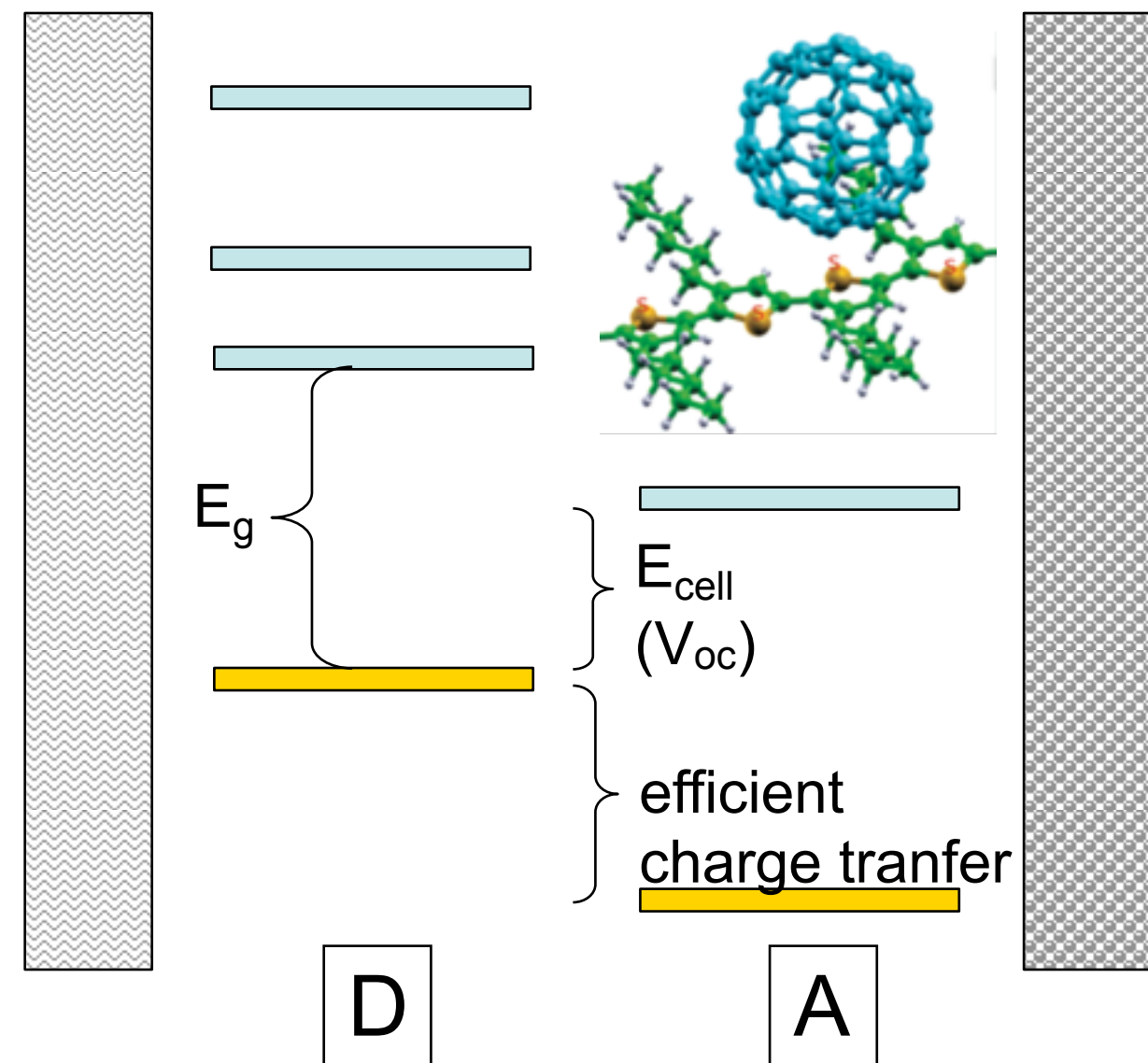
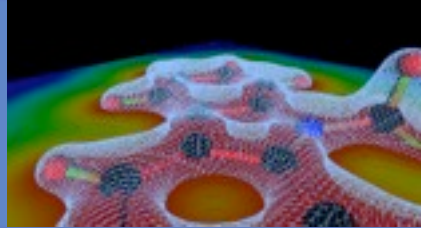


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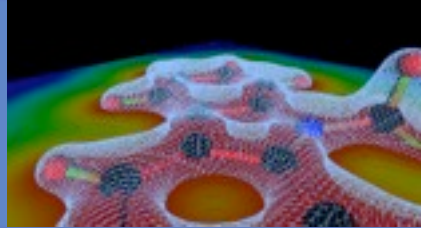
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 - (for good charge transfer)

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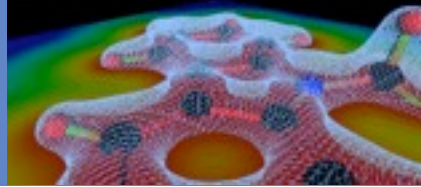


- Gap: optimal for solar spectrum
- HOMO :
 - higher than C_{60}
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 - low enough for good V_{oc} and air stability

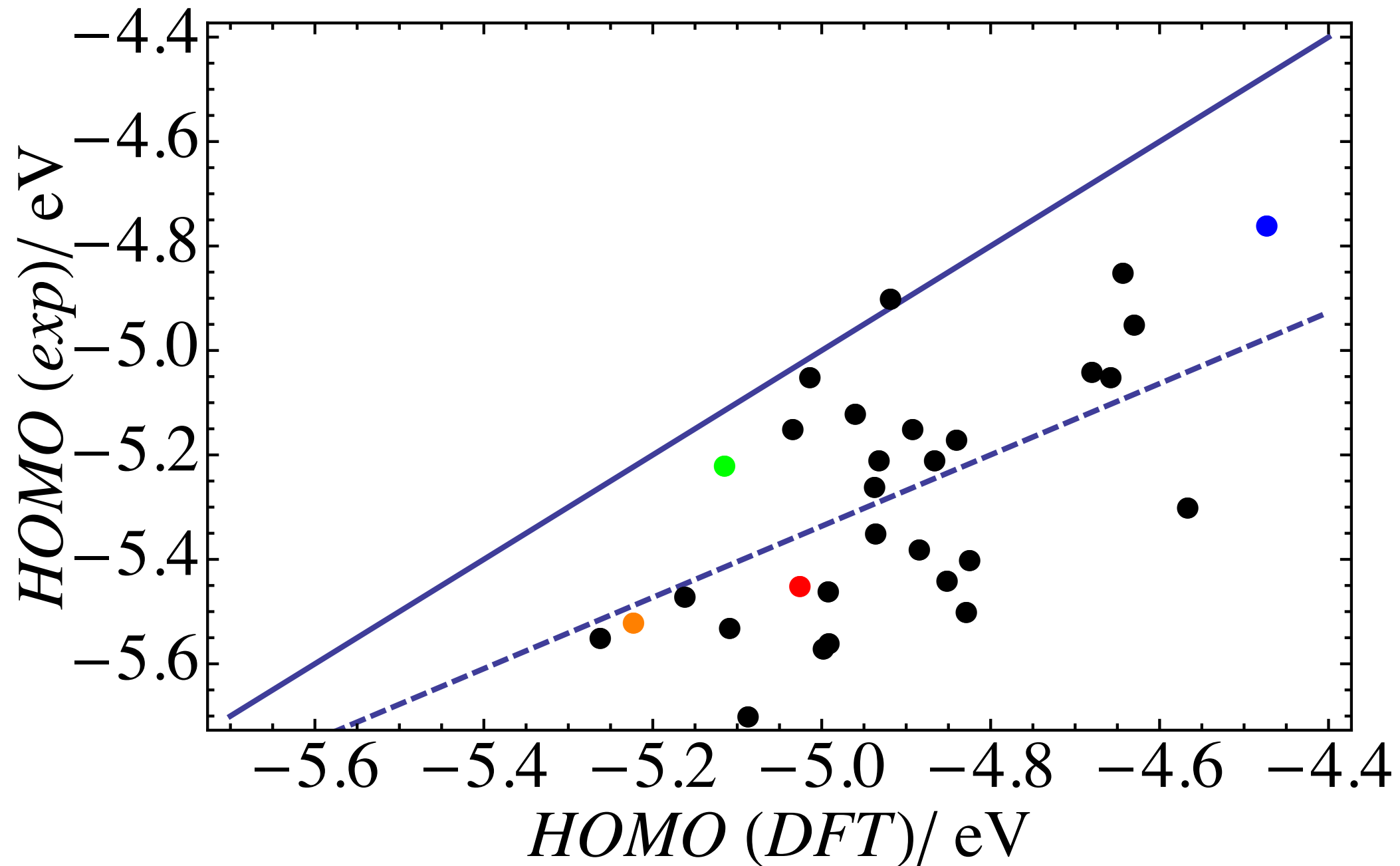
DFT to the rescue?



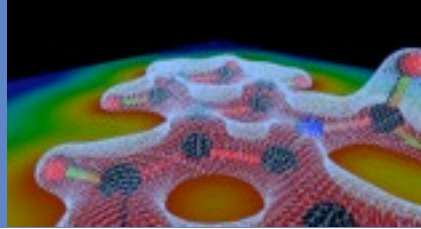
DFT to the rescue?



HOMO



The G_0W_0 : performance



G_0W_0 :

$$(T + V_{ext} + V_H + V_{xc})|\phi_n\rangle = \varepsilon_n|\phi_n\rangle$$

$$G = \sum_{n=1}^{\infty} \frac{|\phi_n\rangle\langle\phi_n|}{\omega - \varepsilon_n}$$

$$P = -iGG$$

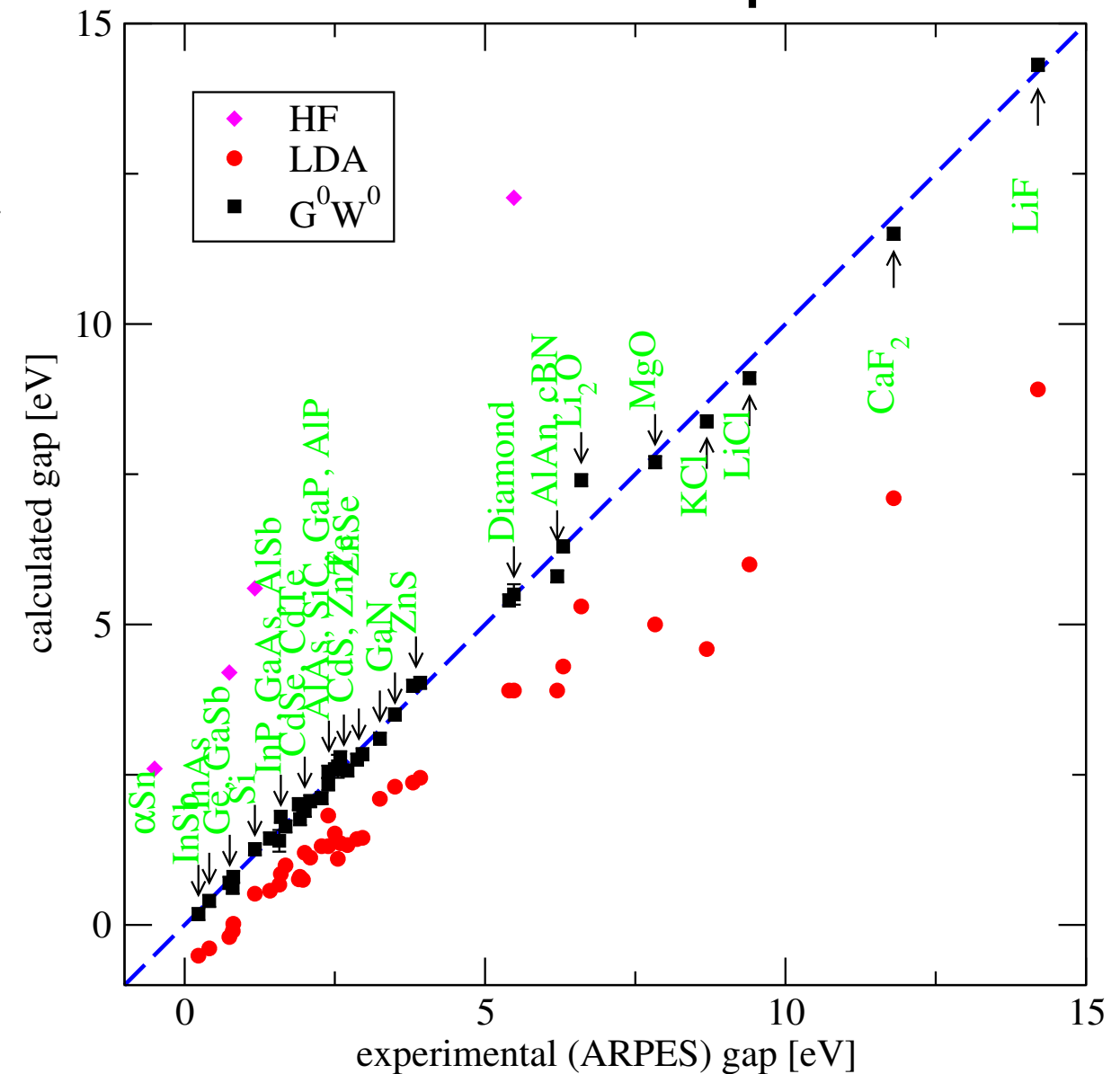
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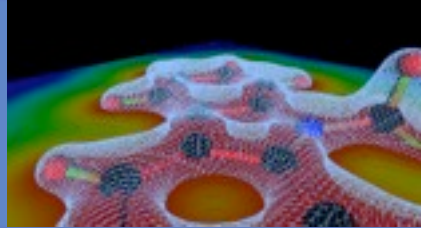
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LDA and G_0W_0 : precision



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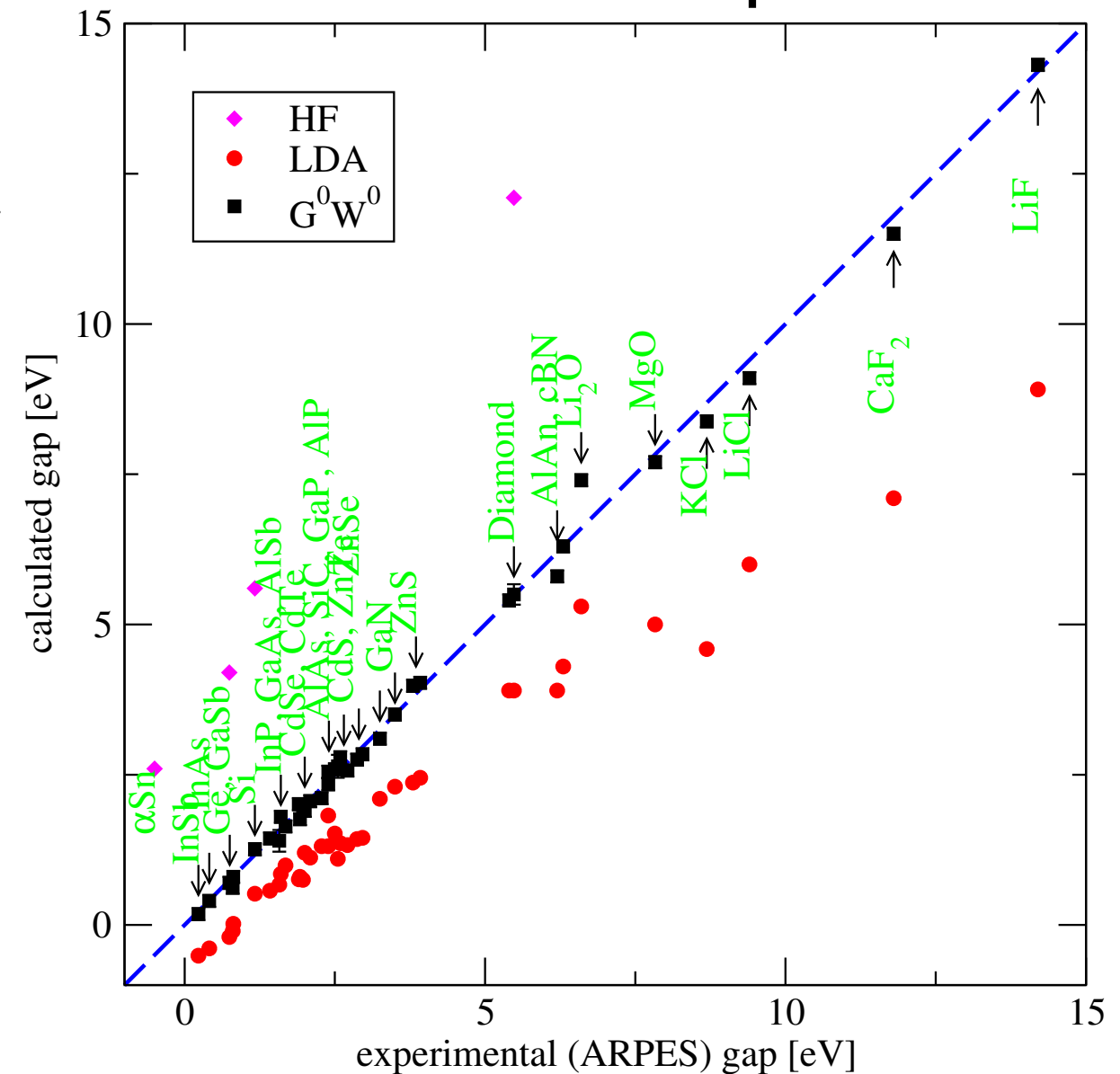
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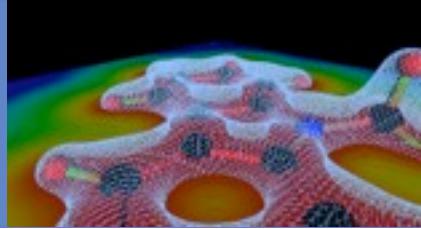
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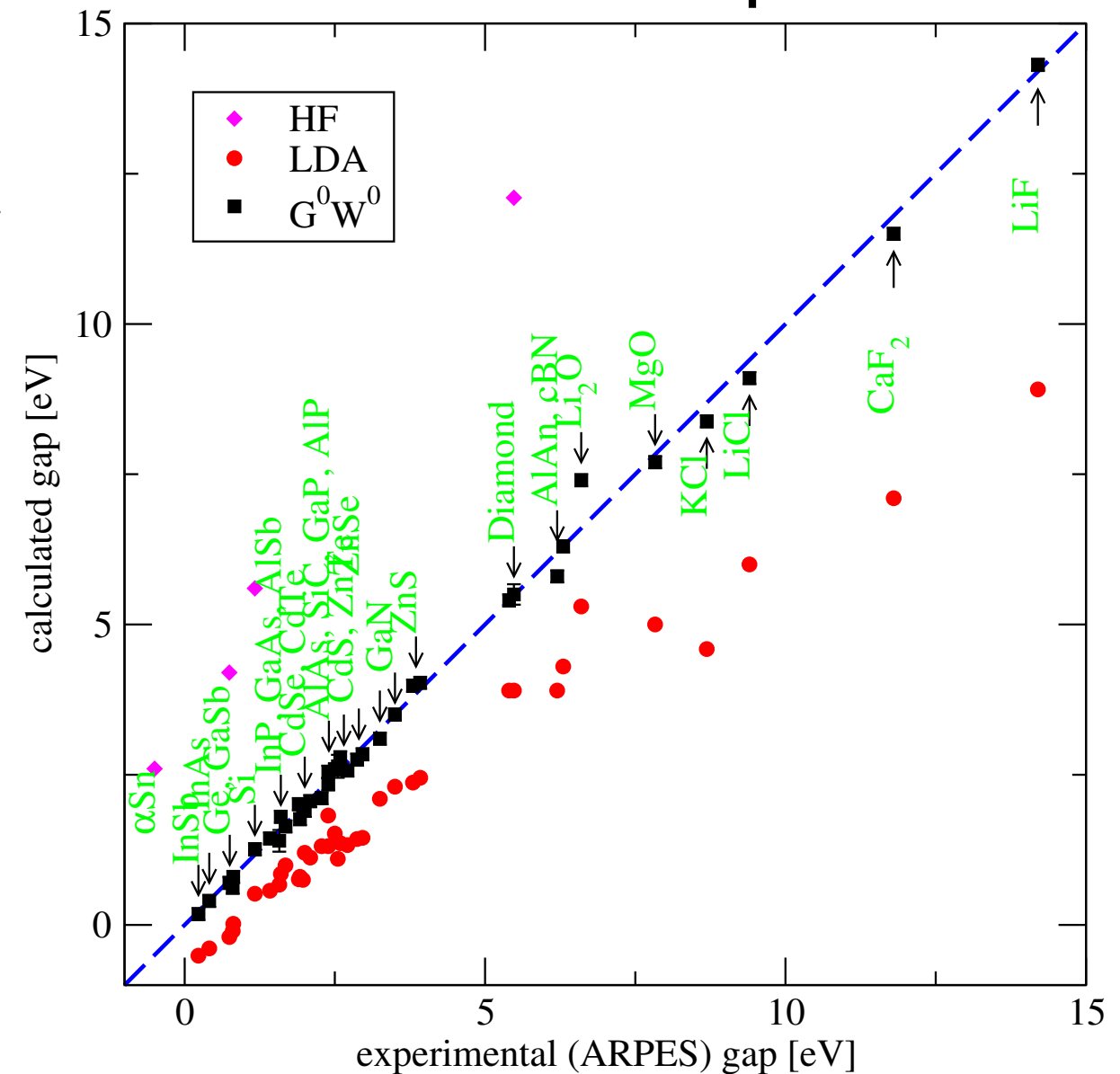
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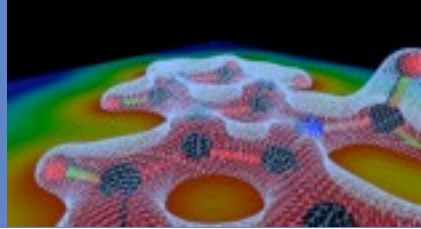


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S. Faleev, PRL, 2004

The G_0W_0 : bottlenecks



- Why G_0W_0 so expensive ?

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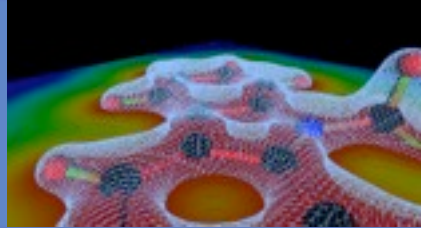
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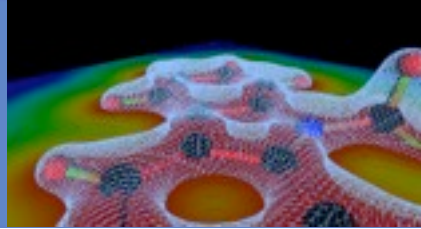
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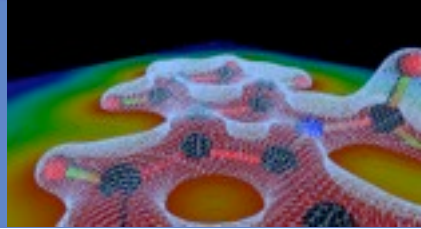
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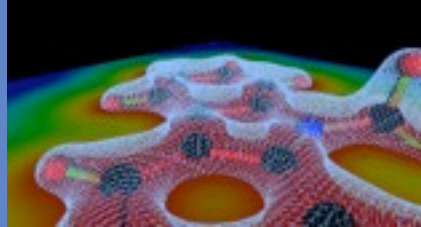
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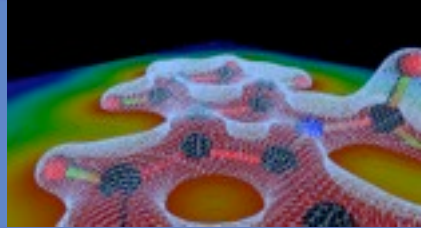
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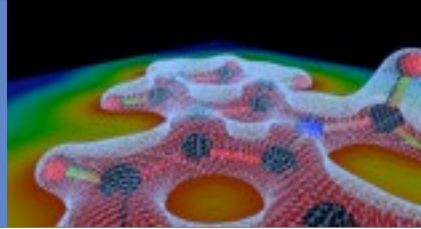
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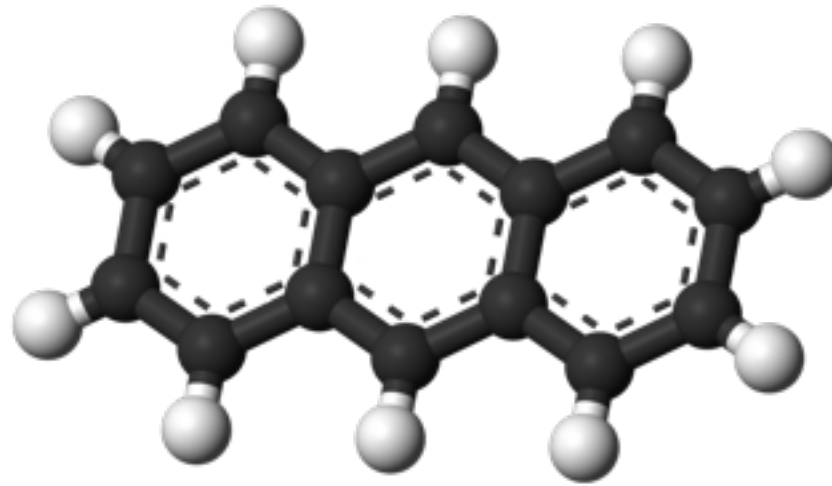
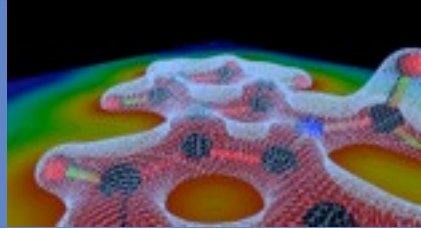
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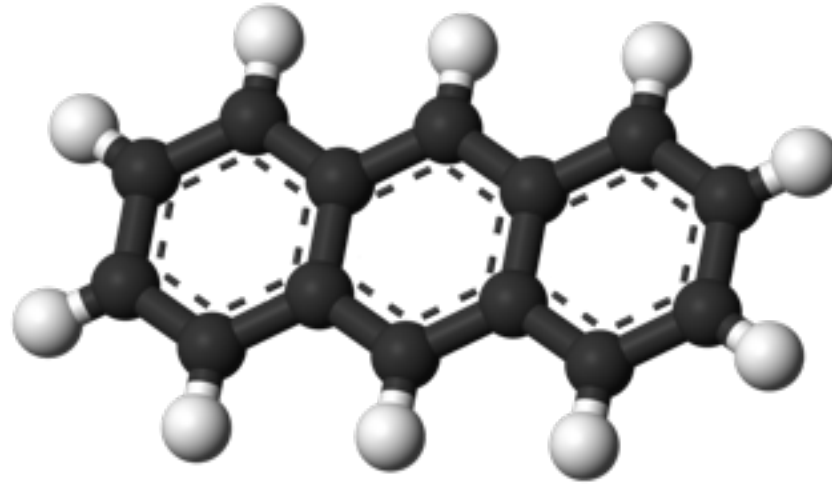
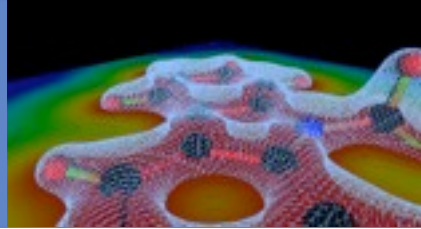
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- inversion of $\epsilon \Rightarrow N^3$ operation (N = basis size)

The case of anthracene

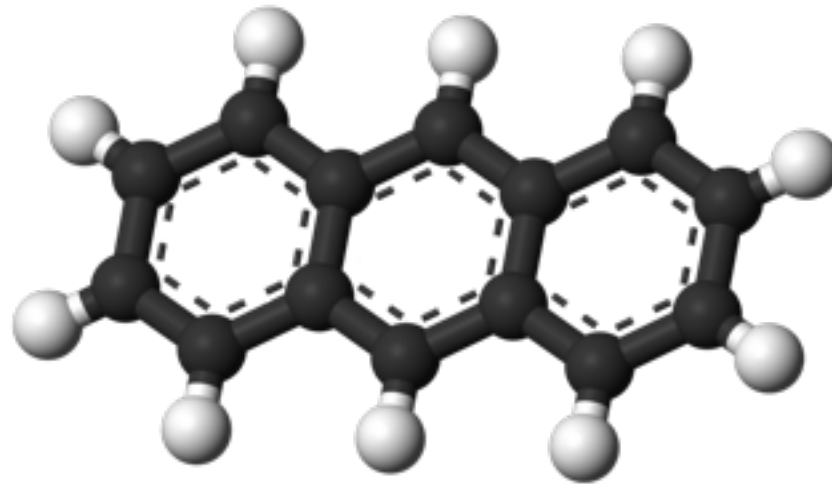
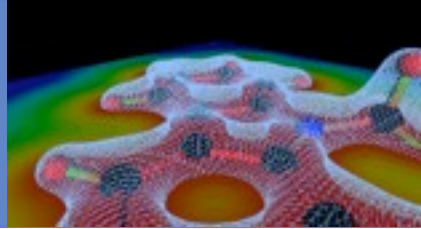


The case of anthracene



$$N_v = 33$$

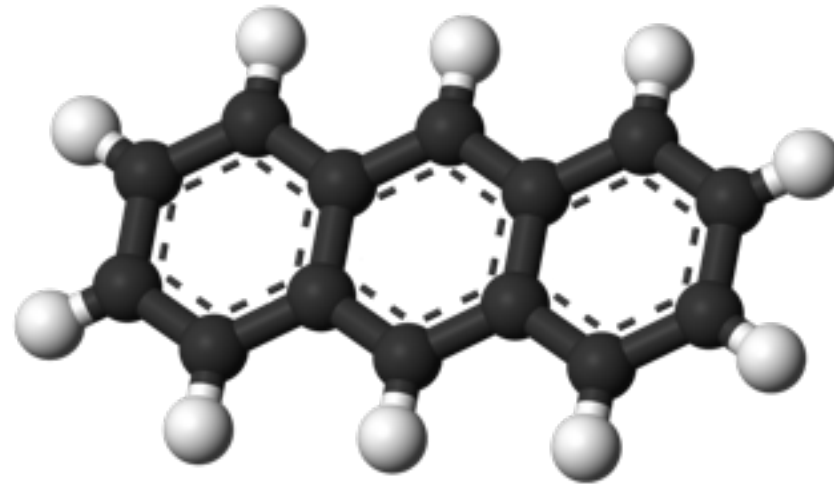
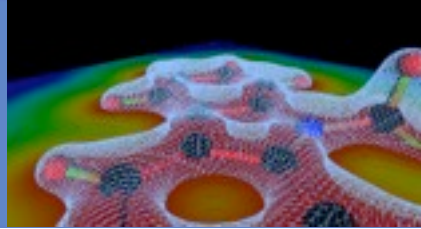
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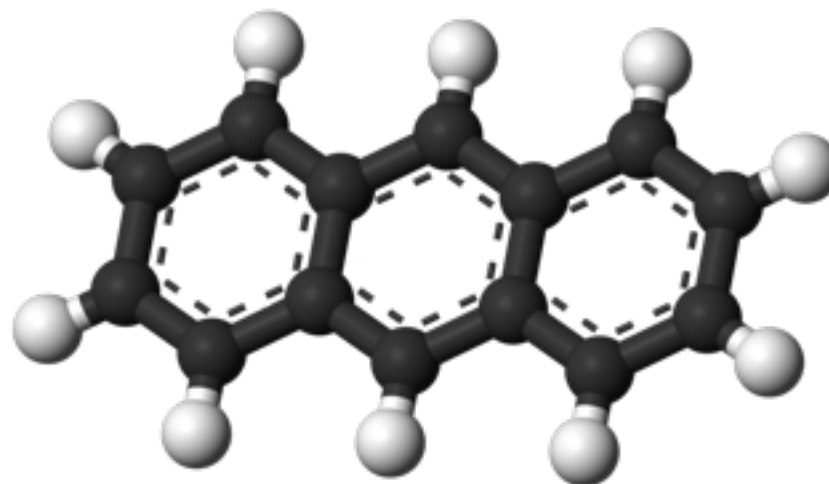
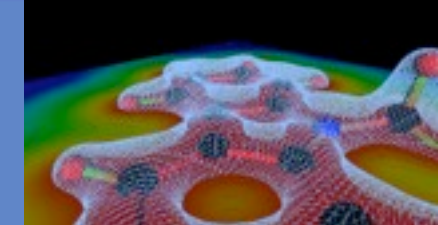


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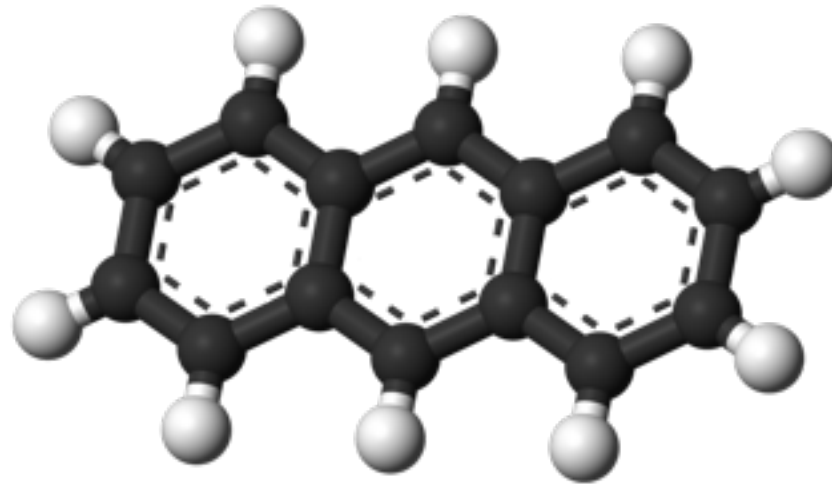
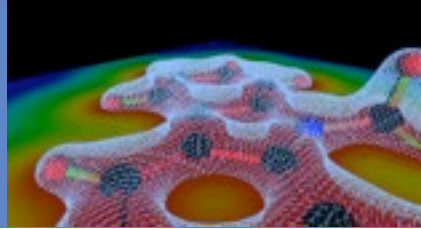
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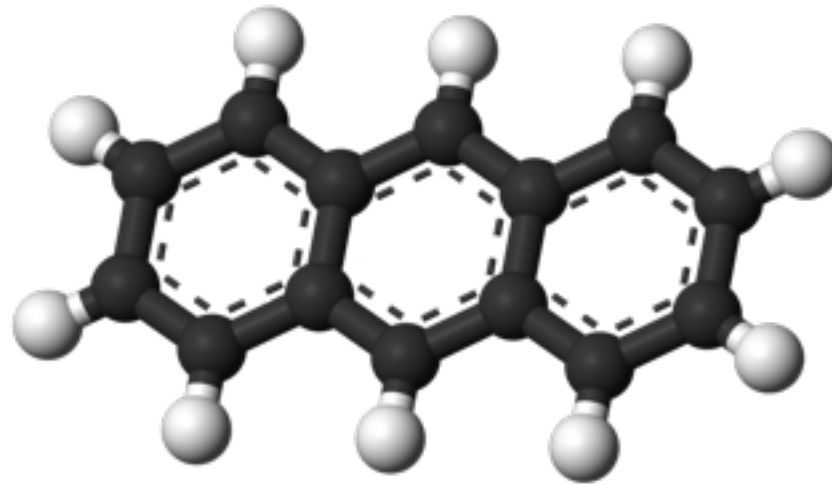
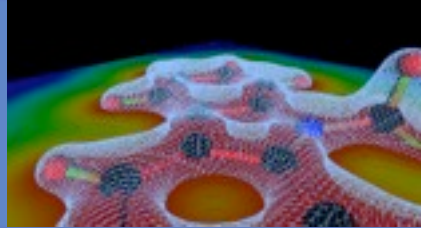
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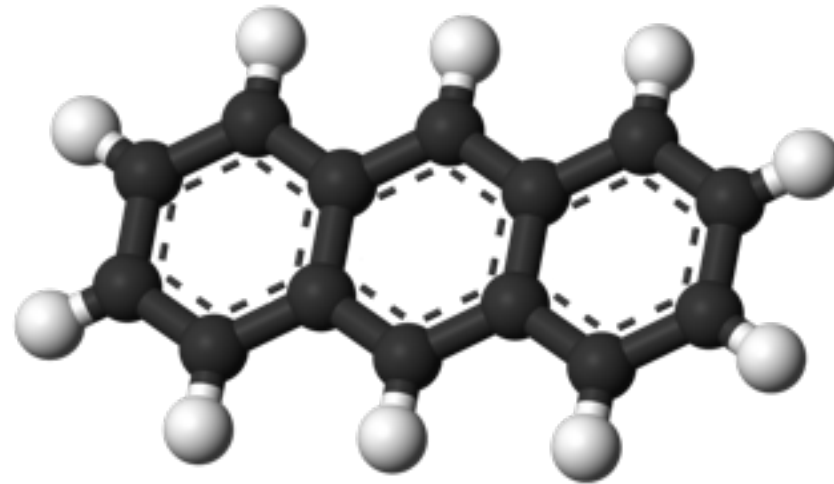
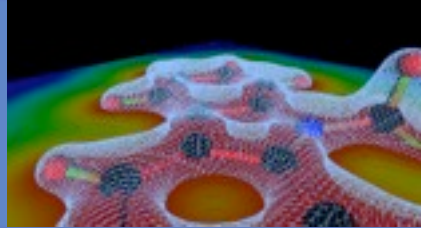
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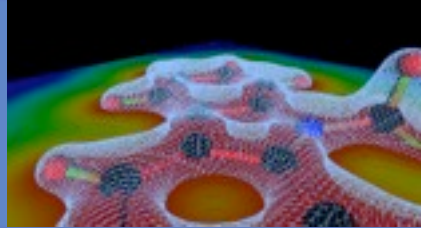
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The plan



$$P = -i \textcircled{G} G$$

$$\epsilon = 1 - vP$$

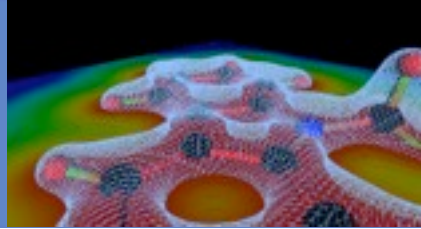
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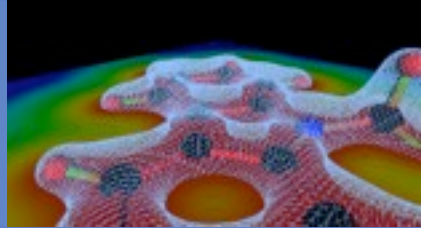
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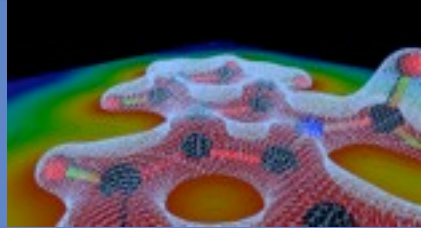
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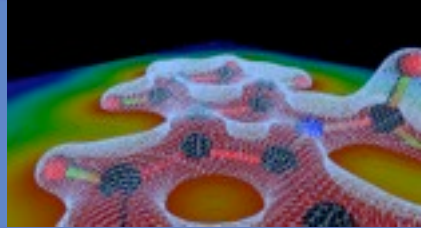
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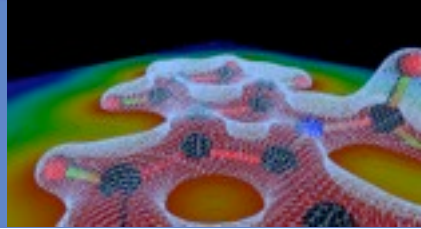
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- summations $\rightarrow$ Sternheimer's equations

- planewaves basis $\rightarrow$ Lanczos basis

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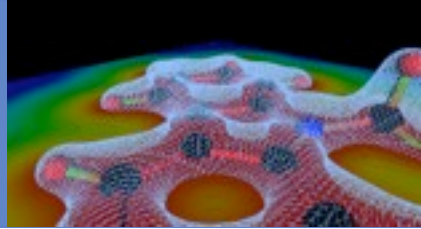
Sternheimer's equation



$$\bullet P|\psi\rangle = \sum_v |v\rangle \left(\sum_c |c\rangle \frac{1}{\omega - (\varepsilon_c - \varepsilon_v)} - \frac{1}{\omega + (\varepsilon_c - \varepsilon_v)} \langle c| \langle v| \right) |\psi\rangle$$

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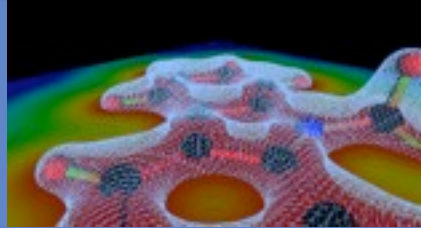
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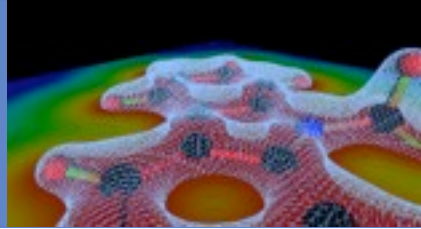
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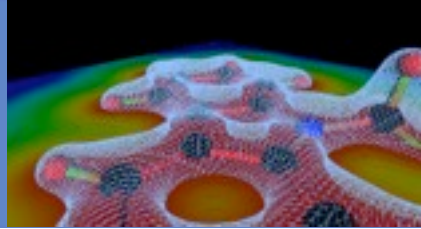
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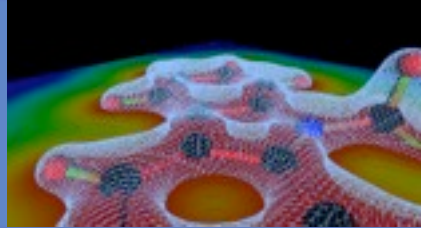
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Sternheimer's equation



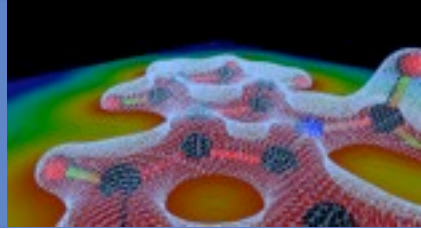
$$\bullet P|\psi\rangle = \sum_v |v\rangle \left(\sum_c |c\rangle \frac{1}{\omega - (\varepsilon_c - \varepsilon_v)} - \frac{1}{\omega + (\varepsilon_c - \varepsilon_v)} \langle c| \langle v| \right) |\psi\rangle$$

• We define :

$$|\phi_v^-\rangle \equiv \sum_c \frac{|c\rangle \langle c|}{\omega - (\varepsilon_c - \varepsilon_v)} |v\rangle |\psi\rangle = \frac{1}{\omega - H + \varepsilon_v} \mathcal{P}_c |v\rangle |\psi\rangle$$

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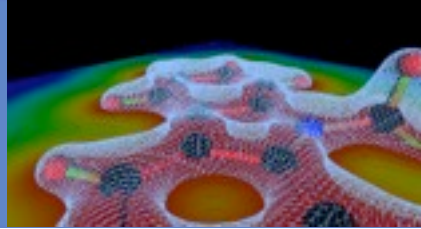
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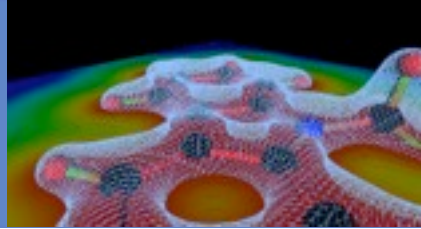
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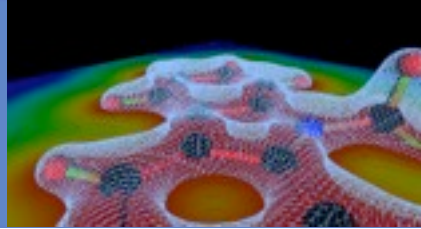
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Sternheimer's equation



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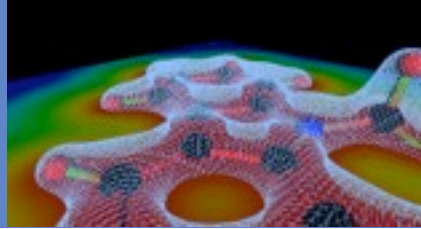
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$$(H - \epsilon_v - \omega) |\phi_v^-\rangle = -\mathcal{P}_c |v\rangle |\psi\rangle \quad \text{Sternheimer equation}$$

$$1) P = -iGG$$

Sternheimer's equation



- $P|\psi\rangle = \sum_v |v\rangle \left(\sum_c |c\rangle \frac{1}{\omega - (\epsilon_c - \epsilon_v)} - \frac{1}{\omega + (\epsilon_c - \epsilon_v)} \langle c| \langle v| \right) |\psi\rangle$

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$$\sum_c$$

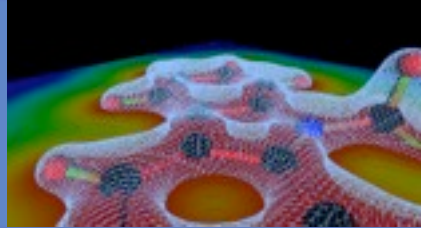


Solving

$$A|x\rangle = |b\rangle$$

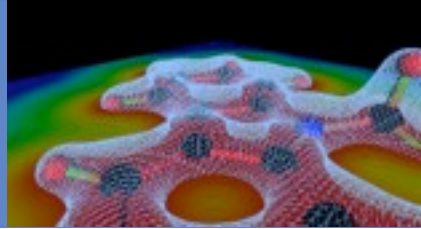
$$2) \Sigma = i\mathbf{G}\mathbf{W}$$

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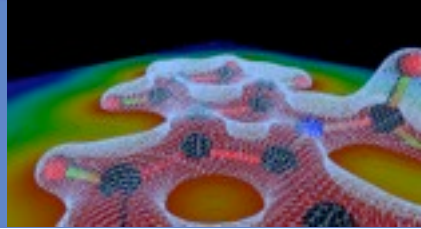


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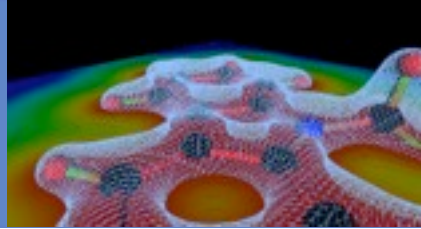


- We want : $\epsilon_m \approx \epsilon_m + \langle m | \Sigma_x + \Sigma_c - V_{xc} | m \rangle$



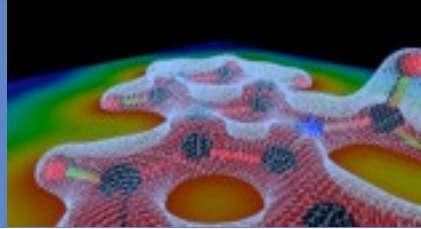
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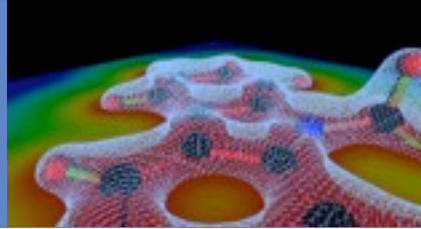
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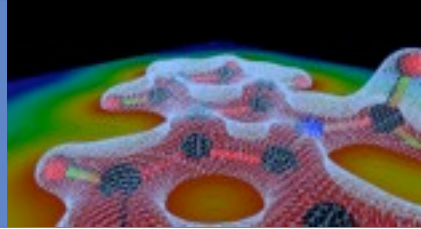
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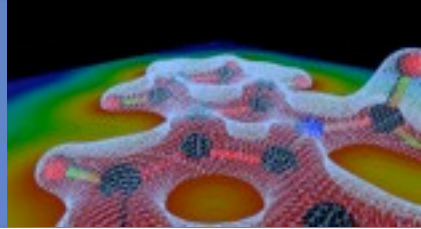
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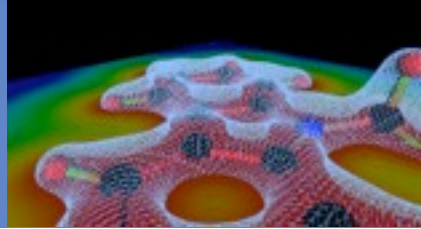
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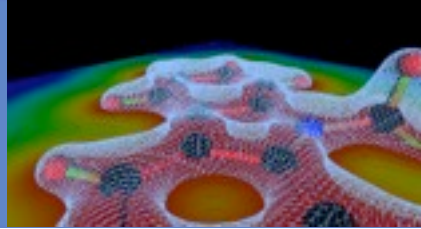
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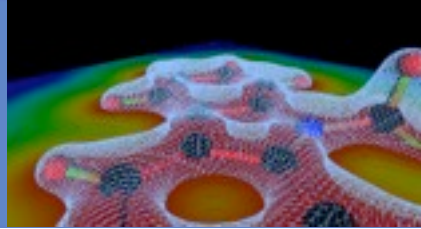
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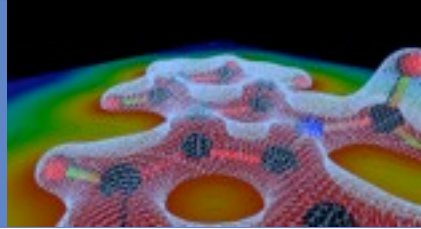
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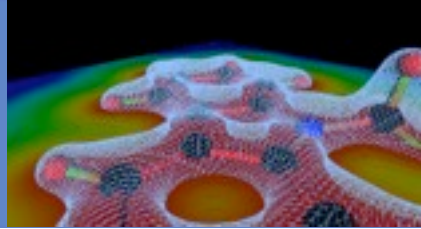


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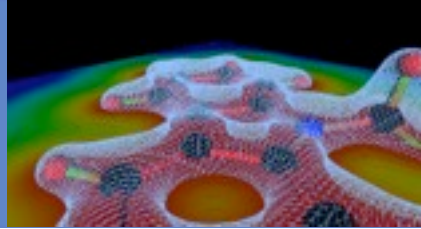
$$2) \Sigma = iGW$$

Sternheimer's equation

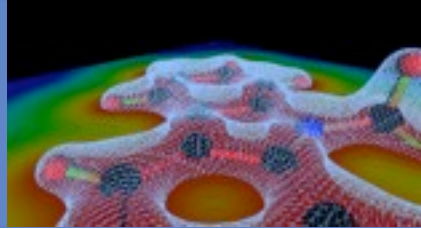


2) $\Sigma = i\textcolor{red}{G}W$

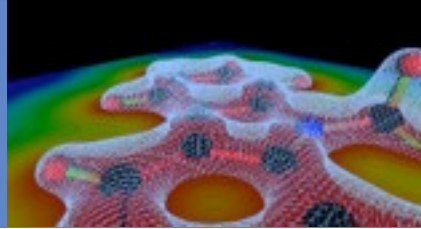
Sternheimer's equation



$$\langle m | \Sigma_c | m \rangle = \frac{i}{2\pi} \int_{-\infty}^{+\infty} d\omega \sum_q \langle q | [\epsilon^{-1} - 1] v \Phi_m^T \frac{1}{\omega - H^T + \varepsilon_m} \Phi_m^{\dagger T} | q \rangle$$

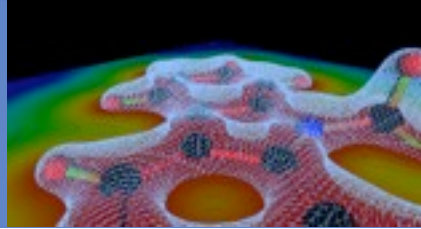


$$\begin{aligned}\langle m | \Sigma_c | m \rangle &= \frac{i}{2\pi} \int_{-\infty}^{+\infty} d\omega \sum_q \langle q | [\epsilon^{-1} - 1] v \Phi_m^T \frac{1}{\omega - H^T + \varepsilon_m} \Phi_m^{\dagger T} | q \rangle \\ &= \frac{i}{2\pi} \int_{-\infty}^{+\infty} d\omega \sum_q \langle q | [\epsilon^{-1} - 1] v \Phi_m^T | x_q \rangle\end{aligned}$$

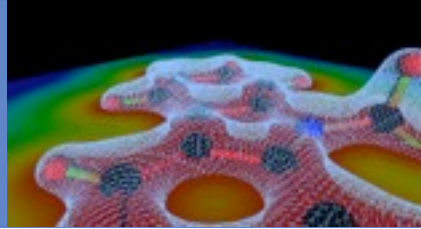


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 \langle m | \Sigma_c | m \rangle &= \frac{i}{2\pi} \int_{-\infty}^{+\infty} d\omega \sum_q \langle q | [\epsilon^{-1} - 1] v \Phi_m^T \left[\frac{1}{\omega - H^T + \epsilon_m} \Phi_m^{\dagger T} | q \rangle \right] \\
 &= \frac{i}{2\pi} \int_{-\infty}^{+\infty} d\omega \sum_q \langle q | [\epsilon^{-1} - 1] v \Phi_m^T | x_q \rangle
 \end{aligned}$$

A blue box highlights the term $\frac{1}{\omega - H^T + \epsilon_m} \Phi_m^{\dagger T} | q \rangle$ in the first line, and a blue arrow points from this box to the circled $| x_q \rangle$ in the second line.

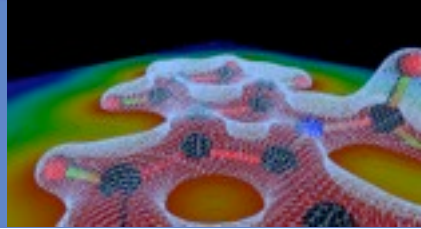


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 &= \frac{i}{2\pi} \int_{-\infty}^{+\infty} d\omega \sum_q \langle q | [\epsilon^{-1} - 1] v \Phi_m^T | \textcolor{blue}{x}_q \rangle \leftarrow \\
 |x_q\rangle &\equiv \frac{1}{\omega - H^T + \varepsilon_m} \Phi_m^{\dagger T} | q \rangle
 \end{aligned}$$



$$\begin{aligned}
 \langle m | \Sigma_c | m \rangle &= \frac{i}{2\pi} \int_{-\infty}^{+\infty} d\omega \sum_q \langle q | [\epsilon^{-1} - 1] v \Phi_m^T \left[\frac{1}{\omega - H^T + \epsilon_m} \Phi_m^{\dagger T} | q \rangle \right] \\
 &= \frac{i}{2\pi} \int_{-\infty}^{+\infty} d\omega \sum_q \langle q | [\epsilon^{-1} - 1] v \Phi_m^T | x_q \rangle \\
 | x_q \rangle &\equiv \frac{1}{\omega - H^T + \epsilon_m} \Phi_m^{\dagger T} | q \rangle
 \end{aligned}$$

Diagrammatic annotations: A blue box highlights the fraction $\frac{1}{\omega - H^T + \epsilon_m} \Phi_m^{\dagger T} | q \rangle$ in the first line. A blue arrow points from this box to the state $|x_q\rangle$ in the second line, which is circled in blue. Another blue arrow points from the denominator $\omega - H^T + \epsilon_m$ in the third line to the same state $|x_q\rangle$.



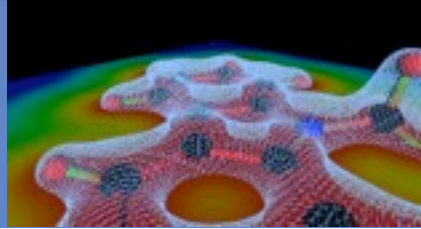
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$$\leftarrow |x_q\rangle \equiv \frac{1}{\boxed{\omega - H^T + \varepsilon_m}} \Phi_m^{\dagger T} | q \rangle$$

$$\left(\omega - H^T + \varepsilon_m \right) |x_q\rangle = \Phi_m^{\dagger T} | q \rangle \quad \text{Sternheimer equation}$$

$$2) \Sigma = i \mathbf{G} \mathbf{W}$$

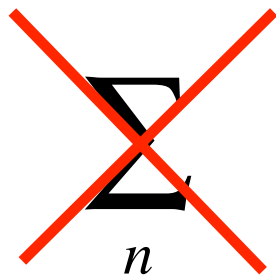
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$$\begin{aligned} \langle m | \Sigma_c | m \rangle &= \frac{i}{2\pi} \int_{-\infty}^{+\infty} d\omega \sum_q \langle q | [\epsilon^{-1} - 1] v \Phi_m^T \left[\frac{1}{\omega - H^T + \epsilon_m} \Phi_m^{\dagger T} | q \rangle \right] \\ &= \frac{i}{2\pi} \int_{-\infty}^{+\infty} d\omega \sum_q \langle q | [\epsilon^{-1} - 1] v \Phi_m^T | x_q \rangle \end{aligned}$$

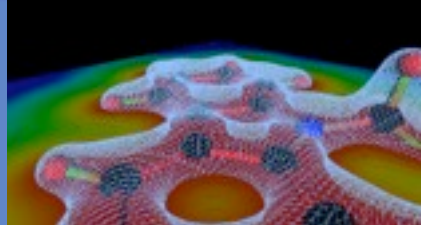
$$|x_q\rangle \equiv \frac{1}{\omega - H^T + \epsilon_m} \Phi_m^{\dagger T} | q \rangle$$

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$$2) \Sigma = iGW$$

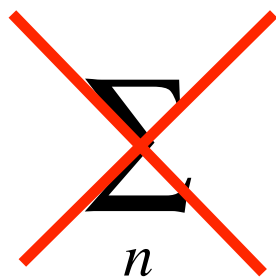
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$$\begin{aligned} \langle m | \Sigma_c | m \rangle &= \frac{i}{2\pi} \int_{-\infty}^{+\infty} d\omega \sum_q \langle q | [\epsilon^{-1} - 1] v \Phi_m^T \frac{1}{\omega - H^T + \epsilon_m} \Phi_m^{\dagger T} | q \rangle \\ &= \frac{i}{2\pi} \int_{-\infty}^{+\infty} d\omega \sum_q \langle q | [\epsilon^{-1} - 1] v \Phi_m^T | x_q \rangle \end{aligned}$$

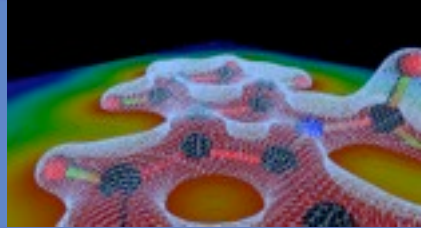
$$|x_q\rangle \equiv \frac{1}{\omega - H^T + \epsilon_m} \Phi_m^{\dagger T} | q \rangle$$

$$(\omega - H^T + \epsilon_m) |x_q\rangle = \Phi_m^{\dagger T} | q \rangle \quad \text{Sternheimer equation}$$

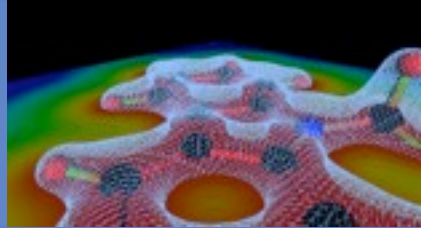


Solving

$$A|x\rangle = |b\rangle$$

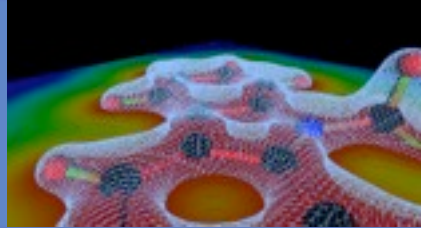


$$\langle m | \Sigma^c | m \rangle = \frac{i}{2\pi} \int_{-\infty}^{+\infty} d\omega \sum_q \langle \textcolor{blue}{q} | [\epsilon^{-1} - 1] (\dots) | \textcolor{blue}{q} \rangle$$



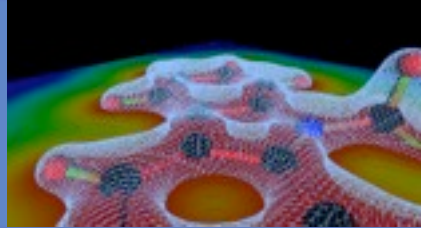
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- Need a basis: $\{|\textcolor{blue}{q}\rangle\}$



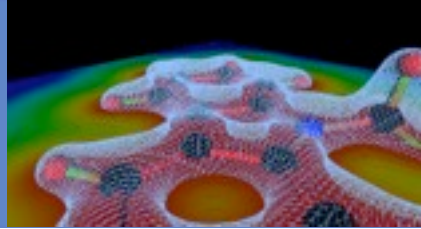
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- Need a basis: $\{ | \textcolor{blue}{q} \rangle \}$
- The ideal basis:



$$\langle m | \Sigma^c | m \rangle = \frac{i}{2\pi} \int_{-\infty}^{+\infty} d\omega \sum_q \langle q | [\epsilon^{-1} - 1] (\dots) | q \rangle$$

- Need a basis: $\{|q\rangle\}$
- The ideal basis:
 - Is small, e.g.
 $\Rightarrow$ NOT planewaves

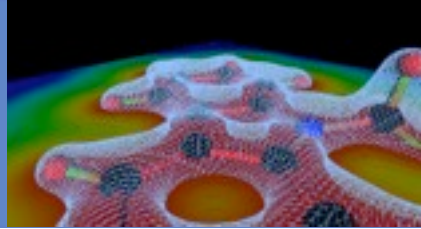


$$\langle m | \Sigma^c | m \rangle = \frac{i}{2\pi} \int_{-\infty}^{+\infty} d\omega \sum_q \langle q | [\epsilon^{-1} - 1] (\dots) | q \rangle$$

- Need a basis: $\{|q\rangle\}$
- The ideal basis:
 - Is small, e.g.
 $\Rightarrow$ NOT planewaves
 - Is easy to compute
 $\Rightarrow$ NOT $\{|n\rangle\}$

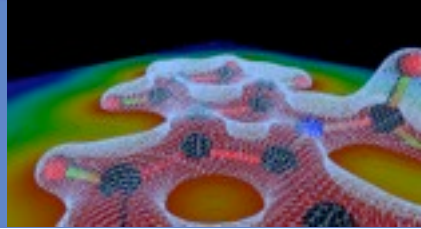
$$2) \Sigma = iGW$$

Lanczos algorithm



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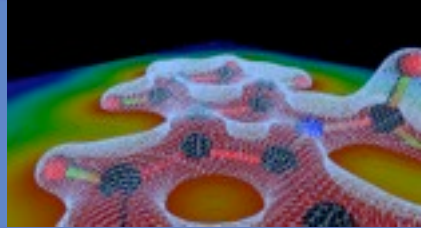
Lanczos algorithm



- Idea :

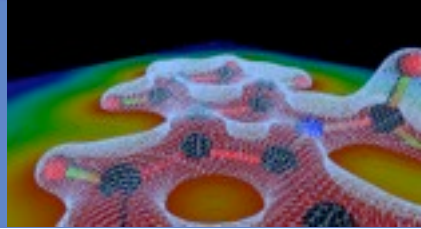
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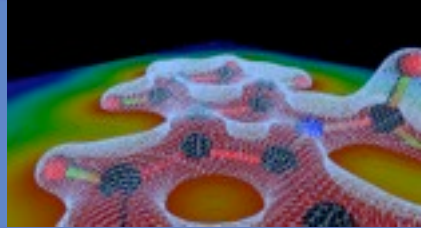
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- efficient sampling of big eigenvalues of : $\epsilon^{-1} - 1 \approx 1 - \epsilon = vP$



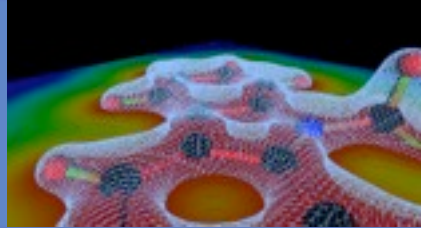
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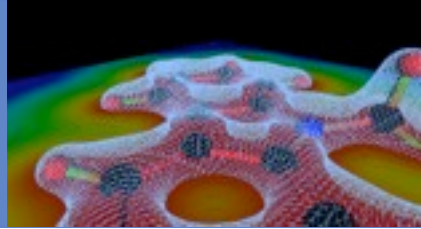
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 - biggest eigenvalues pop out

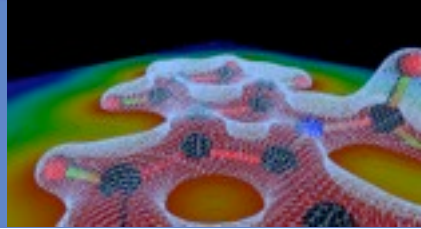


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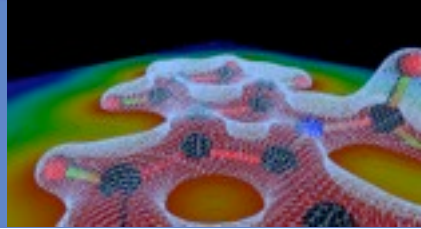
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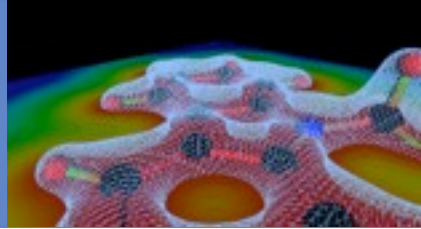
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- Lanczos procedure: same $\{|q\rangle\}$
 - Don't pay all orthogonalization
 - Obtain ϵ for free

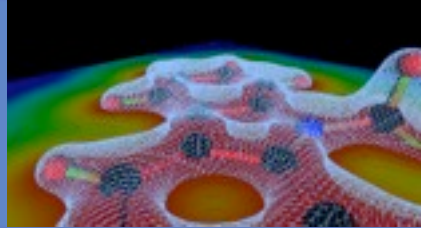
$$3) W = \epsilon^{-1} v$$

Lanczos algorithm



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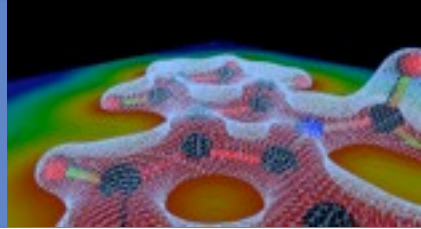
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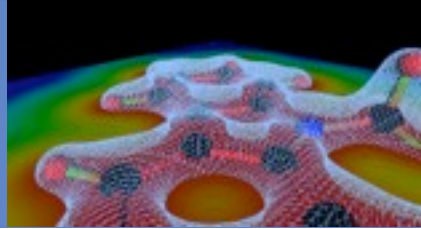
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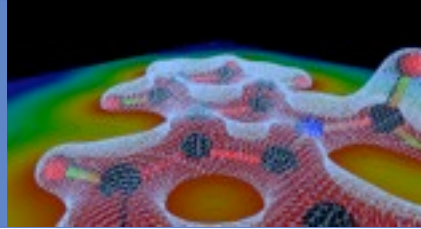
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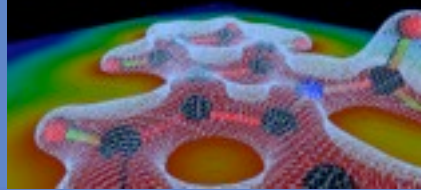


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- We have a small $\{|q\rangle\}$

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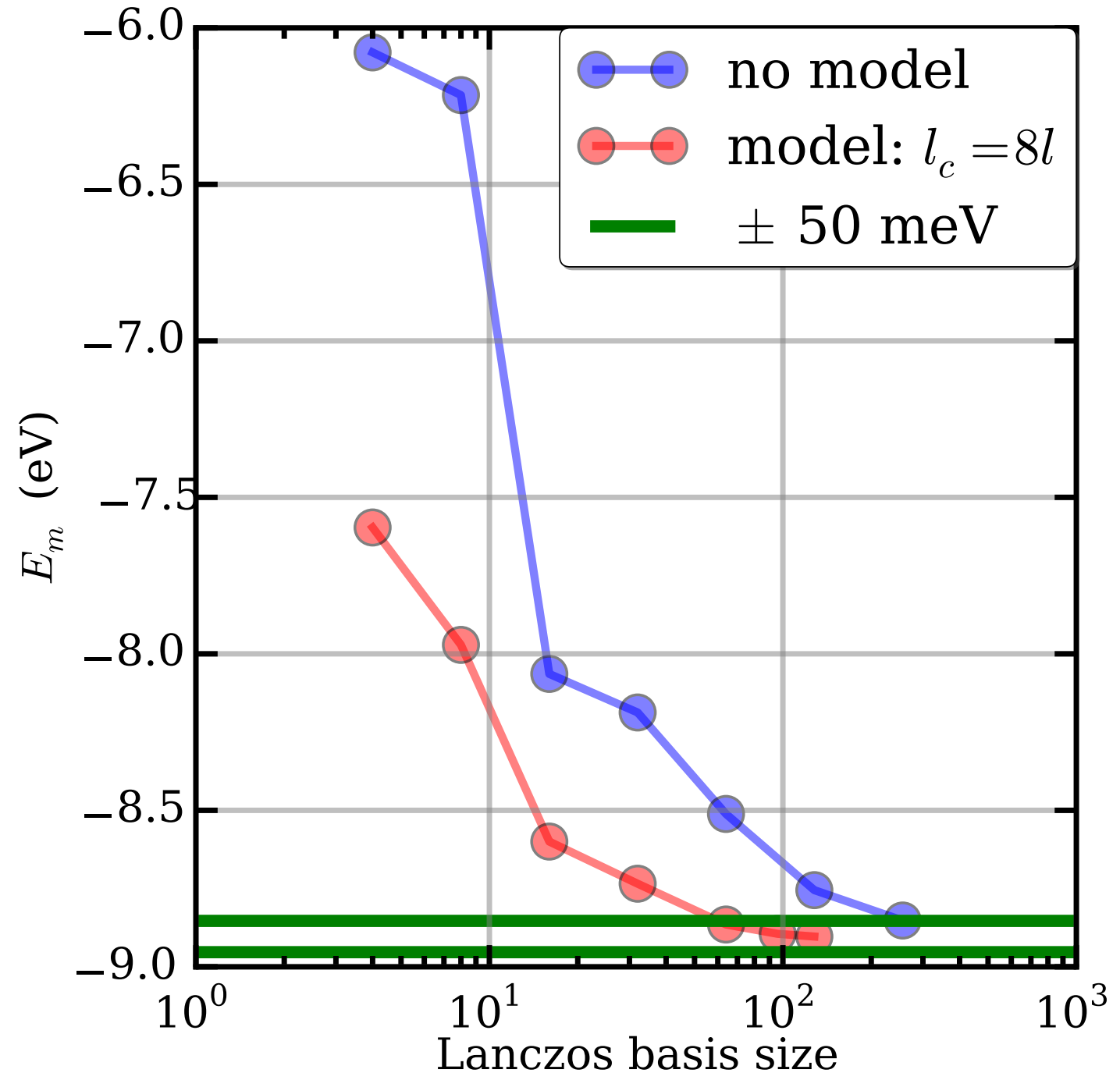
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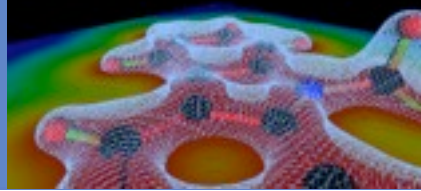
Benzene

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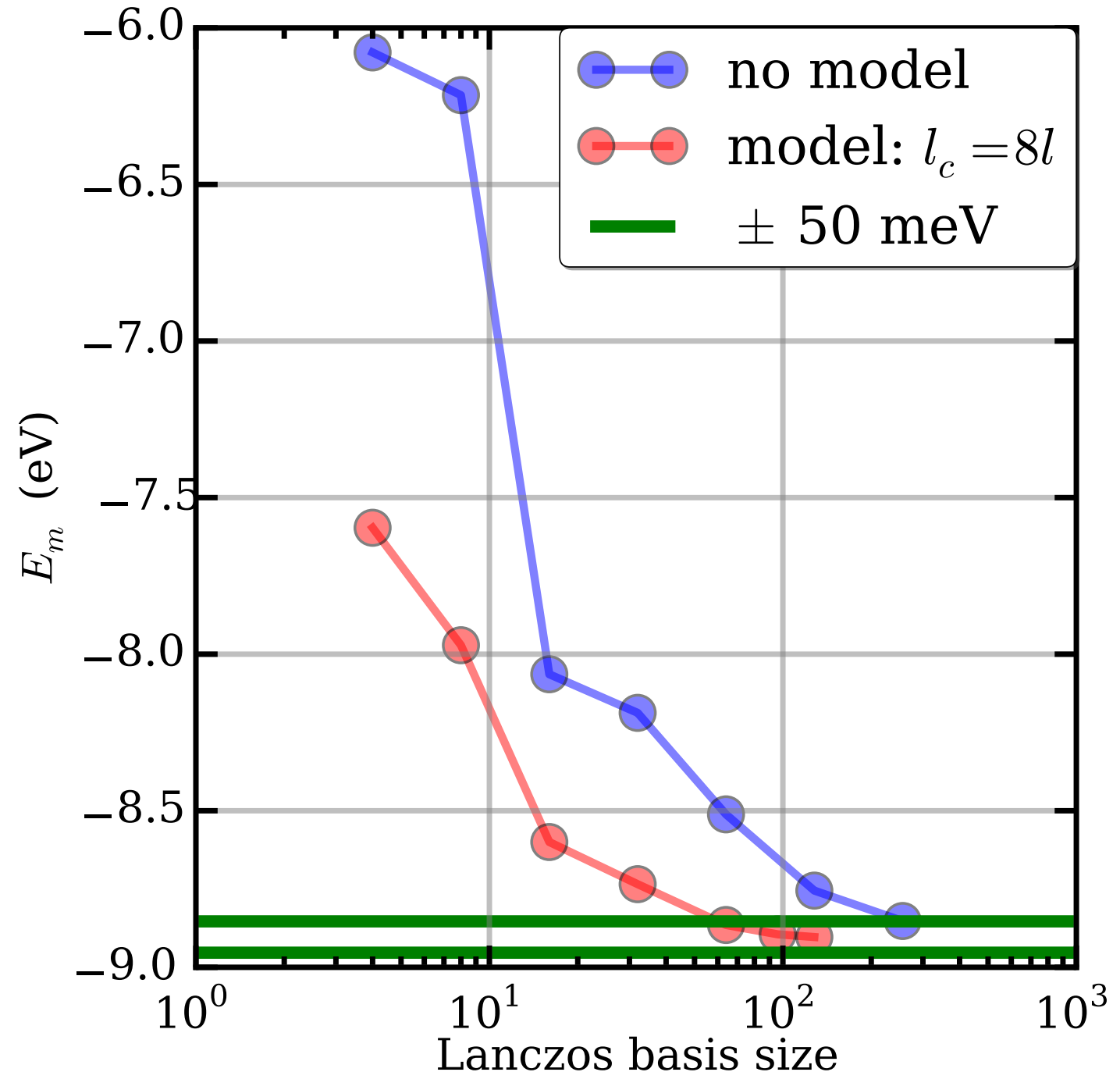
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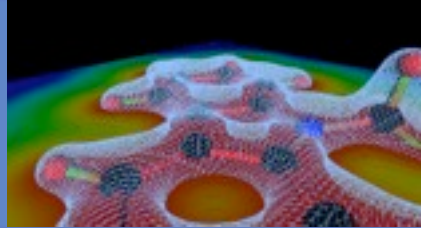
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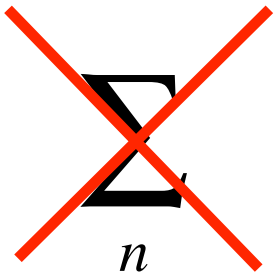
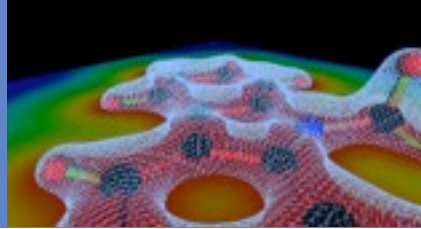
- We have a small $\{|q\rangle\}$
 $\rightarrow \epsilon^{-1} \sim \text{free}$



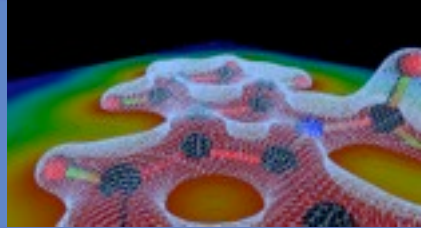
Performance



Performance



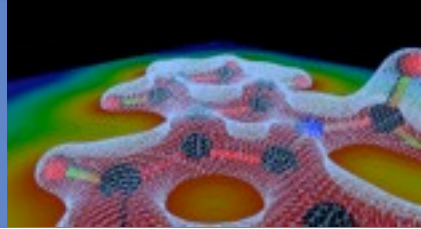
Performance



$$\sum_n$$

$$\epsilon^{-1}$$

(smaller)

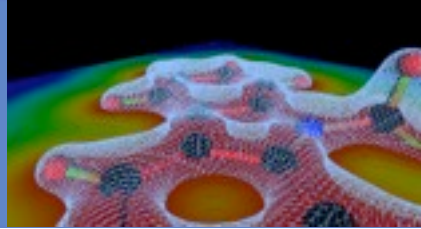


~~$$\sum_n$$~~

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(smaller)

$$A|x\rangle = |b\rangle$$



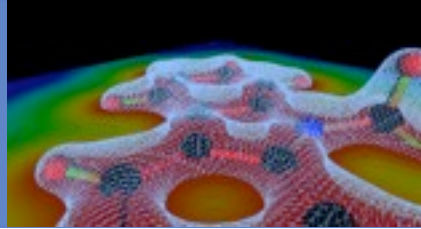
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$$A|x\rangle = |b\rangle$$

$$\{|q\rangle\}$$



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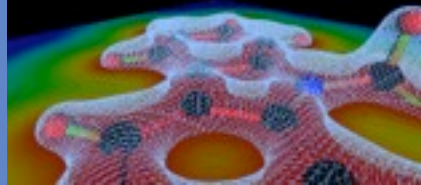
(smaller)

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Faster ?

Performance

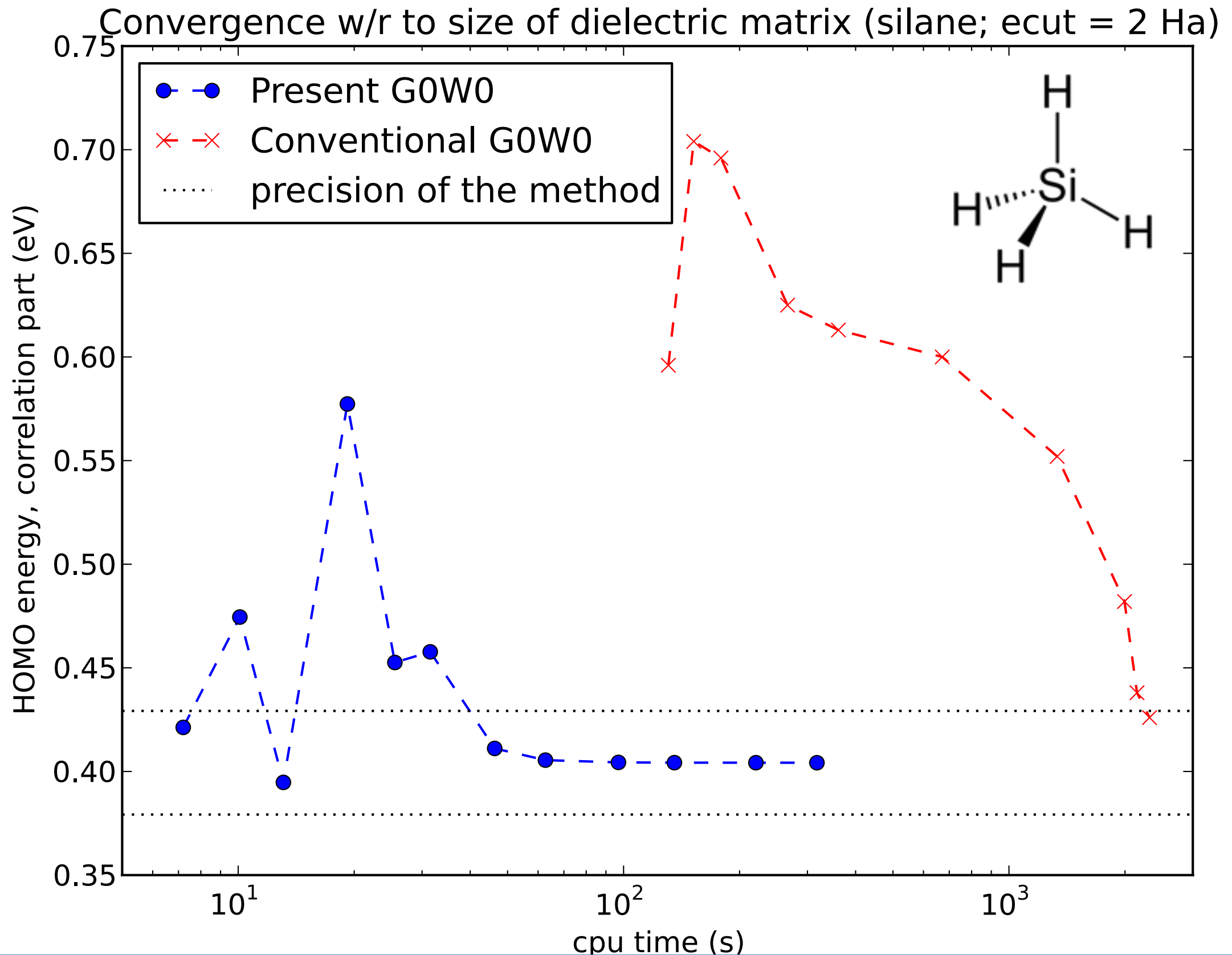


~~Σ_n~~
 ~~ϵ^{-1}~~
 (smaller)

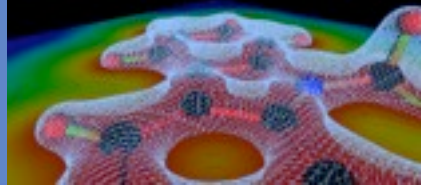
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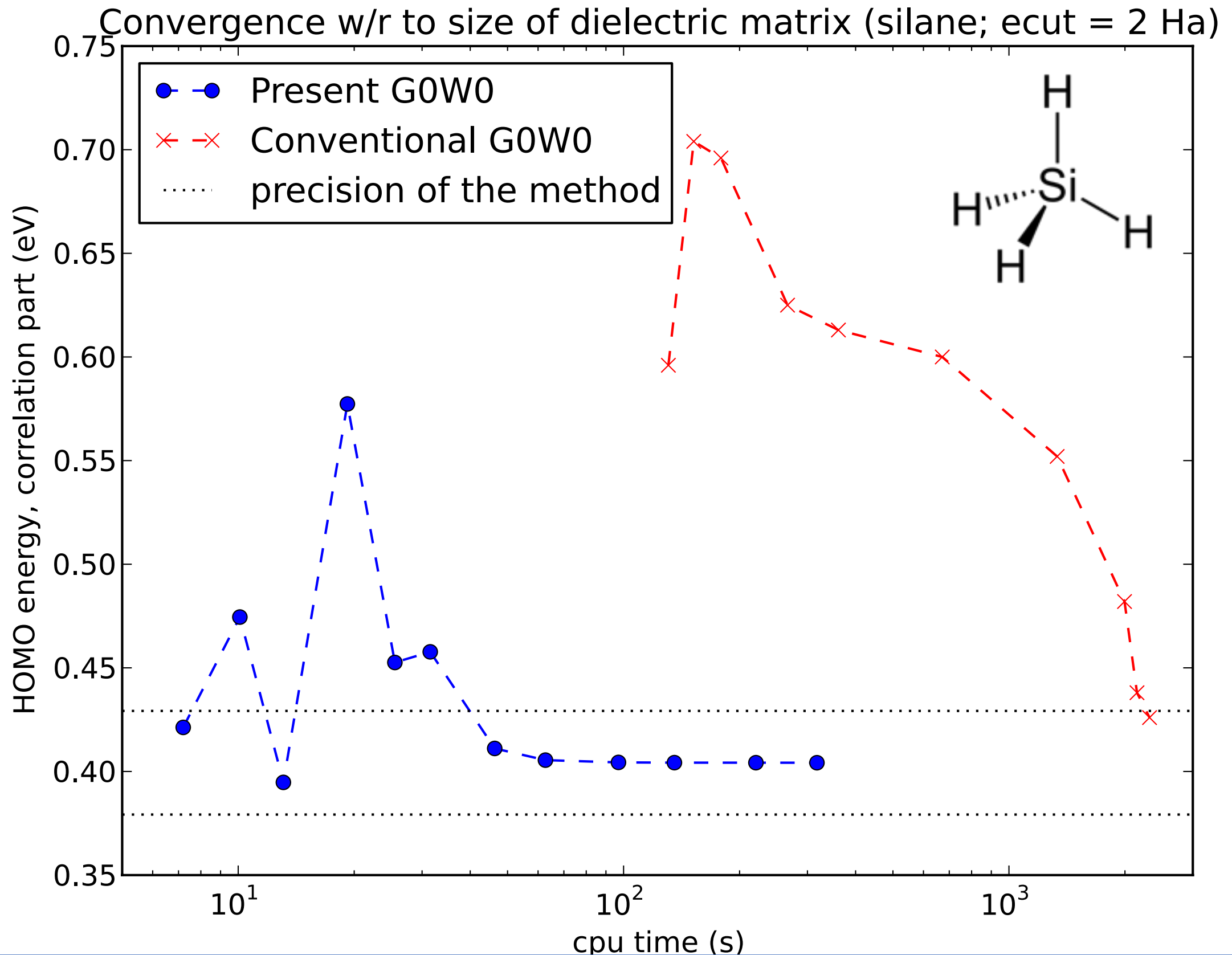
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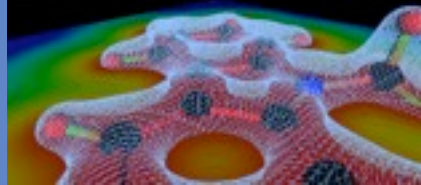
$$\{|q\rangle\}$$

Faster ?

Yes!



Performance



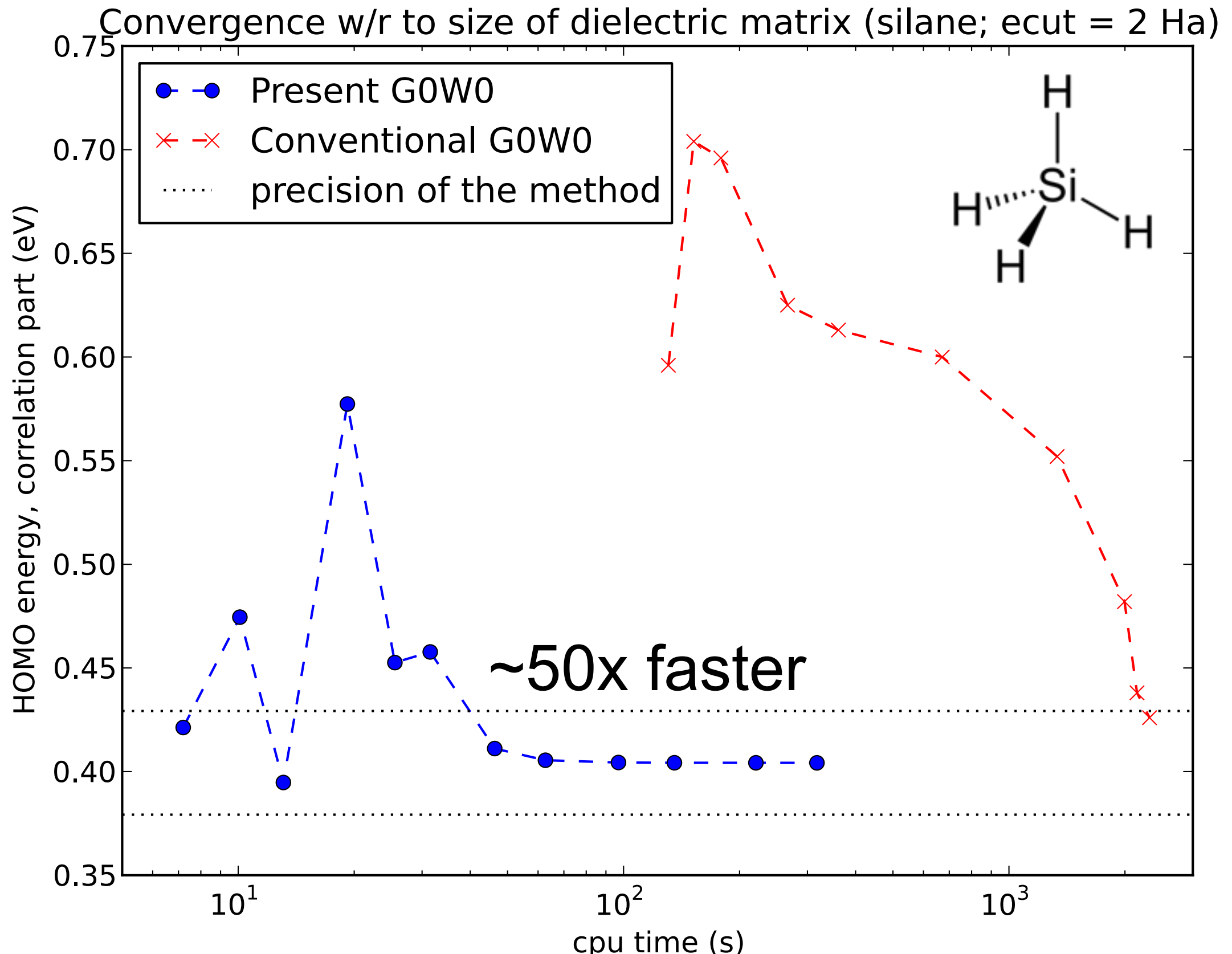
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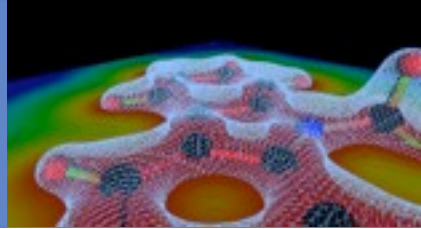
$$\{|q\rangle\}$$

Faster ?

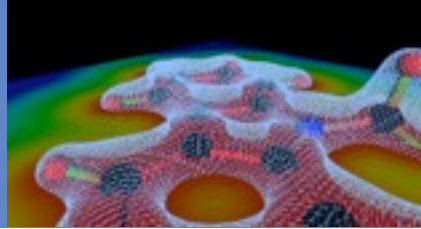
Yes!



Preliminary results

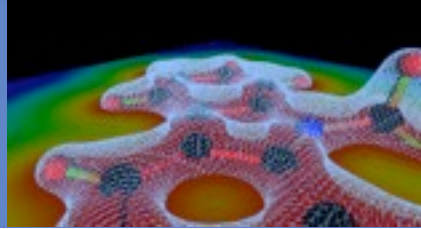


Preliminary results



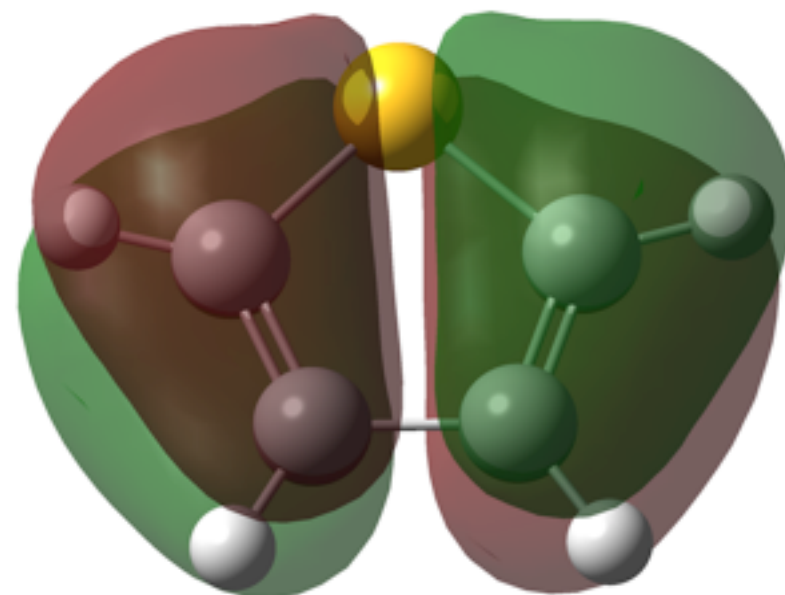
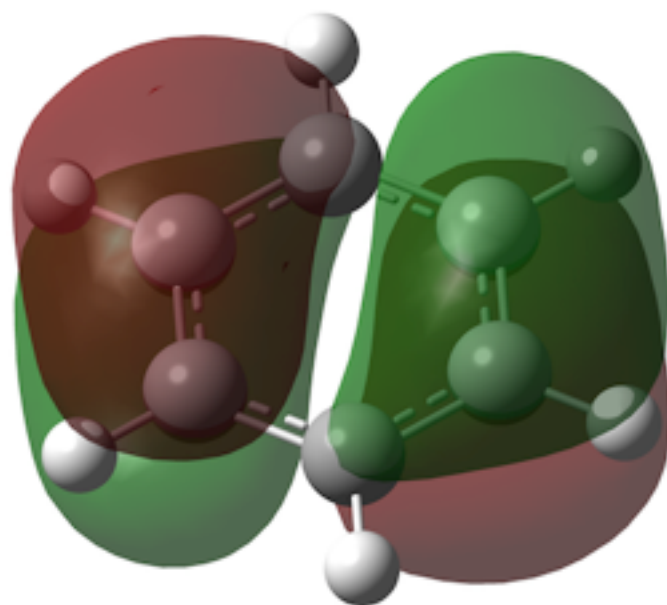
- Working in reals systems?

Preliminary results

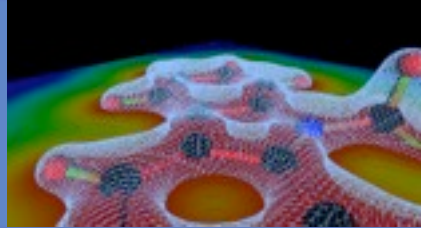


- Working in reals systems?

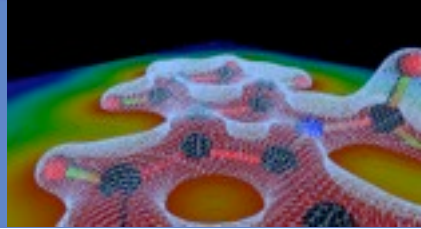
	HOMO (eV)		
	LDA	G ₀ W ₀	Exp.
Benzene	-6.51	-9.22	-9.30
Thiophene	-6.05	-8.94	-8.85



Conclusion

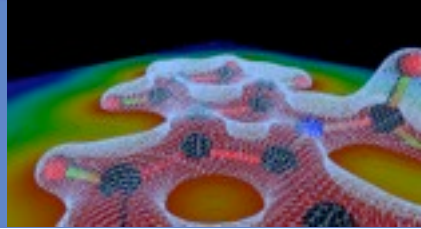


Conclusion



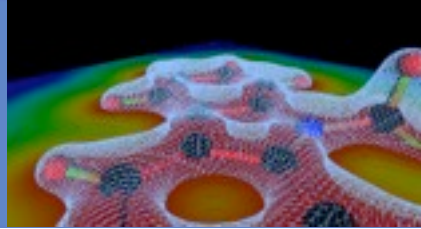
- Bottleneck assessed :

Conclusion

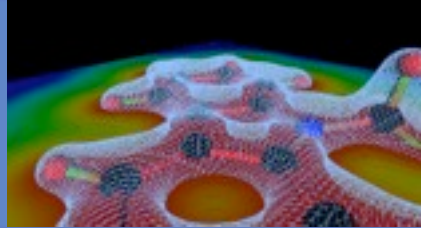


- Bottleneck assessed :
 - no knowledge of conduction states required

Conclusion

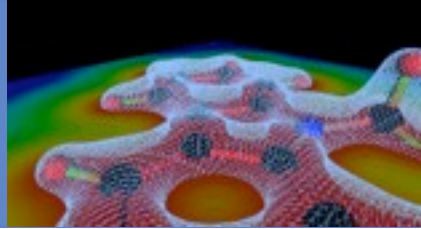


- Bottleneck assessed :
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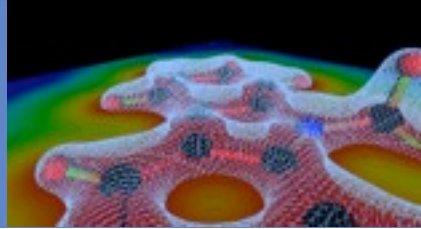
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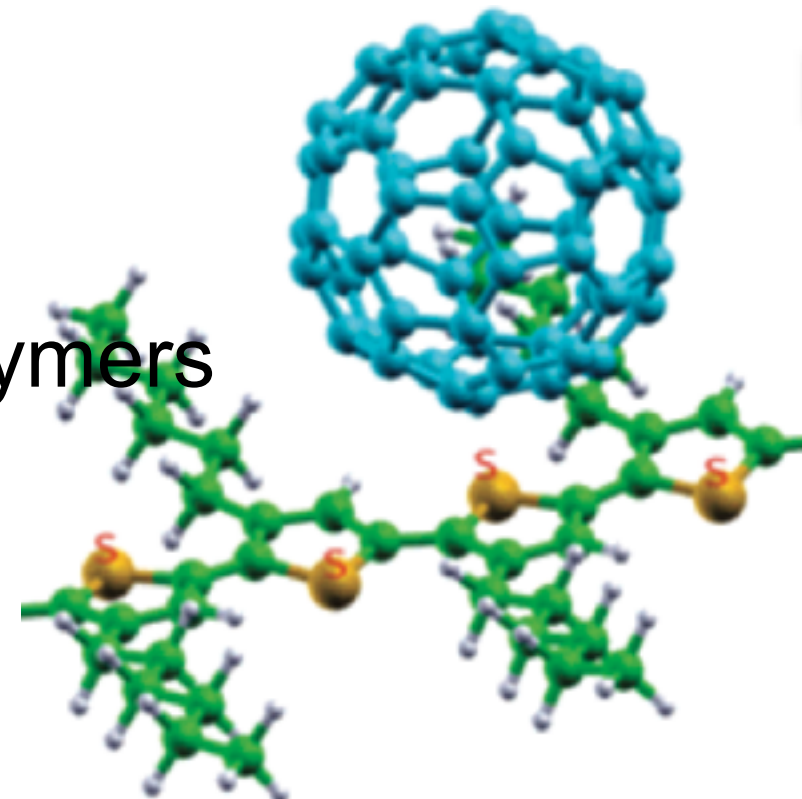


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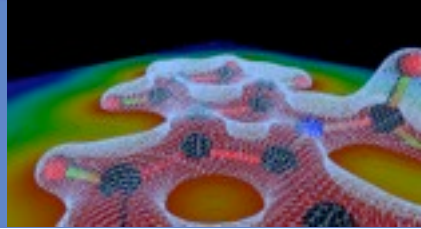
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- Future work :
 - refine DFT calculations for candidate polymers
 - interface states : what do they look like?

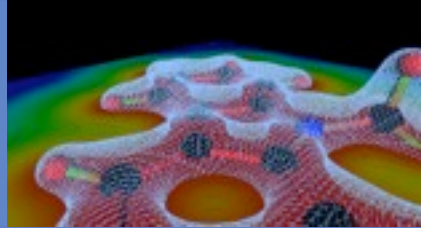


Acknowledgments



Thank you!

Acknowledgments



Michel Côté's group

Thank you!



Bruno Rousseau



Nicolas Bérubé



Gabriel Antonius

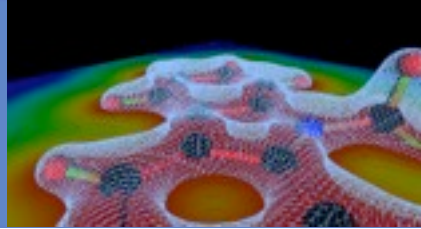


Simon Blackburn

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Vincent Gosselin

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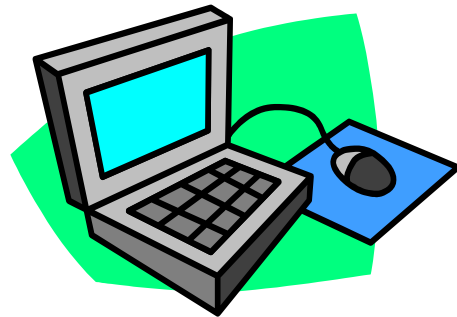
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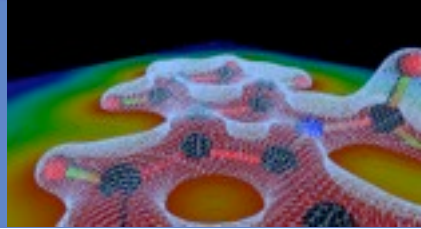
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Calcul Québec



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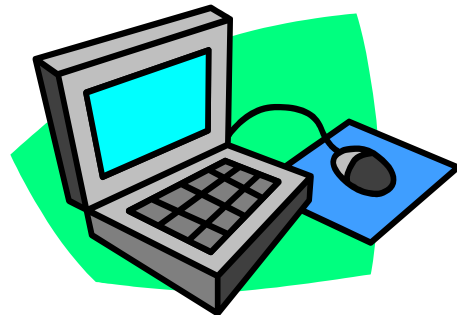
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