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Search for narrow trijet resonances in proton-proton collisions at $\sqrt{s} = 13$ TeV

The CMS Collaboration*

Abstract

The first search for narrow resonances decaying to three well-separated hadronic jets is presented. The search uses proton-proton collision data corresponding to an integrated luminosity of 138 fb^{-1} at $\sqrt{s} = 13$ TeV, collected at the CERN LHC. No significant deviations from the background predictions are observed between 1.75–9.00 TeV. The results provide the first mass limits on a right-handed boson Z_R decaying to three gluons, an excited quark decaying via a vector boson to three quarks, as well as updated limits on a Kaluza–Klein gluon decaying via a radion to three gluons.

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The standard model (SM) of particle physics has several well-established shortcomings, notably lacking explanations for the nonzero neutrino masses, the issue of dark matter, the matter-antimatter asymmetry of the universe, the hierarchy problem, and the proliferation of quarks and leptons. These shortcomings have motivated a comprehensive program of searches for beyond-the-SM physics at high-energy colliders. Searches for new particles coupled to quarks and/or gluons are especially compelling at hadron colliders, such as the CERN LHC, since these particles could potentially be produced with significant cross sections. Searches for new particles decaying to two hadronic jets (dijets) have been performed numerous times at past and current colliders [1–11]. Hadronic resonances with more complicated signatures have also been explored, including searches for the pair production of resonances decaying to two or more jets [12–19] and a search for resonances decaying to a jet and a boosted dijet resonance [20].

In this Letter, we present the first generic search for singly-produced resolved three-jet (trijet) resonances using proton-proton (pp) collision data recorded by the CMS experiment from 2016 to 2018 at a center-of-mass energy of 13 TeV, corresponding to an integrated luminosity of 138 fb^{-1} . Following techniques established by dijet resonance searches, the trijet invariant mass (m_{jjj}) spectrum is scrutinized for narrow excesses compatible with new resonances decaying to three hadronic jets. No assumptions are made on the structure of the decay, except for the requirement that the final state contain three resolved jets with large transverse momentum (p_{T}). The SM background consists of events containing three or more jets produced exclusively through quantum chromodynamics (QCD), and is estimated from data by fitting the m_{jjj} spectrum with smoothly falling empirical functions. New resonances (X) from three signal models that could potentially address the above-mentioned SM shortcomings are considered: a right-handed Z boson, Z_{R} , decaying directly to three gluons (g) [21]; a Kaluza–Klein excitation of a gluon, G_{KK} , decaying via an intermediate radion to three gluons [22, 23]; and an excited quark, q^* , decaying via an intermediate beyond-the-SM vector boson to three quarks (q) [24]. The latter two scenarios, which involve an initial resonance (X) and an intermediate resonance (Y), are referred to as cascade decays and are characterized by the mass ratio $\rho_{\text{m}} = m_{\text{Y}}/m_{\text{X}}$. Values of ρ_{m} between 0.2 and 0.8 that could lead to the resolved trijet signature are considered. At values of ρ_{m} lower than 0.2, the Y is produced with significant transverse momentum, and its decay products are reconstructed as a single jet; this scenario was explored in Ref. [20]. Tabulated results are provided as a HEPData record for this analysis [25].

The CMS apparatus [26] is a multipurpose, nearly hermetic detector, designed to trigger on [27, 28] and identify electrons, muons, photons, and charged and neutral hadrons [29–31]. The “particle-flow” algorithm [32] aims to reconstruct all individual particles in an event, combining information provided by the all-silicon inner tracker, the crystal electromagnetic calorimeters (ECAL), and the brass-scintillator hadron calorimeters (HCAL), all operating inside a 3.8 T superconducting solenoid, with data from the gas-ionization muon detectors embedded in the flux-return yoke outside the solenoid. The reconstructed particles are used to build τ leptons, jets, and missing transverse momentum [33–35].

Events of interest are selected using a two-tiered trigger system. The first level, composed of custom hardware processors, uses information from the calorimeters and muon detectors to select events at a rate of around 100 kHz within a fixed latency of about $4 \mu\text{s}$ [27]. The second level, known as the high-level trigger, consists of a farm of processors running a version of the full event reconstruction software optimized for fast processing, and reduces the event rate to around 1 kHz before data storage [28].

The primary vertex is taken to be the vertex corresponding to the hardest scattering in the

event, evaluated using tracking information alone [36]. The energy of photons is obtained from the ECAL measurement. The energy of electrons is determined from a combination of the electron momentum at the primary interaction vertex as determined by the tracker, the energy of the corresponding ECAL cluster, and the energy sum of all bremsstrahlung photons spatially compatible with originating from the electron track. The energy of muons is obtained from the curvature of the corresponding track. The energy of charged hadrons is determined from a combination of their momentum measured in the tracker and the matching ECAL and HCAL energy deposits, corrected for the response function of the calorimeters to hadronic showers. Finally, the energy of neutral hadrons is obtained from the corresponding corrected ECAL and HCAL energies.

For each event, hadronic jets are clustered from these reconstructed particles using the infrared- and collinear-safe anti- k_T algorithm [37, 38] with a distance parameter of 0.4. Jet momentum is determined as the vectorial sum of all particle momenta in the jet, and is found from simulation to be, on average, within 5 to 10% of the true momentum over the entire p_T spectrum and detector acceptance. Additional pp interactions within the same or nearby bunch crossings (pileup) can contribute additional tracks and calorimetric energy depositions to the jet momentum. To mitigate this effect, charged particles identified to be originating from pileup vertices are discarded, and an offset correction is applied to correct for remaining contributions. Jet energy corrections are derived from simulation to bring the measured response of jets to that of particle level jets on average. In situ measurements of the momentum balance in dijet, photon + jet, Z + jet, and multijet events are used to account for any residual differences in the jet energy scale between data and simulation [34]. The jet energy resolution amounts typically to 15–20% at 30 GeV, 10% at 100 GeV, and 5% at 1 TeV [34]. Additional selection criteria are applied to each jet to remove jets potentially dominated by anomalous contributions from various subdetector components or reconstruction failures [39]. Jets with $p_T < 50$ GeV are further required to pass a multivariate pileup jet rejection algorithm, which has 99% efficiency for jets originating from the primary vertex of interest [40].

To mitigate the loss of resolution from QCD radiation outside the jet boundaries, a “wide-jet” algorithm introduced for previous CMS dijet searches [41] is used. Using the three leading jets with $p_T > 100$ GeV as seeds, the four-vectors of additional jets in the event with $p_T > 30$ GeV are added to the nearest seed jet within an optimized distance of $\Delta R \equiv \sqrt{(\Delta\eta)^2 + (\Delta\phi)^2} < 1.1$, where η is the pseudorapidity and ϕ is the azimuthal angle in radians. The trijet invariant mass, m_{jjj} , is defined as the invariant mass of the three wide jets obtained with this algorithm.

Monte Carlo (MC) simulation is used to model the contributions from the production of new resonances that assume the Breit–Wigner line shape. The $Z_R \rightarrow ggg$ process is simulated at leading order (LO) in QCD with PYTHIA 8.242 [42]. Two width scenarios are simulated: the nominal left-right symmetric scenario from Ref. [21] assuming SM-like couplings with a width of $\Gamma_{Z_R}/m_{Z_R} \sim 3\%$, and a narrow-width scenario with $\Gamma_{Z_R}/m_{Z_R} \sim 0.01\%$ that represents the cases where the width is negligible with respect to the jet energy resolution. The G_{KK} process is modeled at LO with MADGRAPH5_aMC@NLO 2.7.3 [43], interfaced with PYTHIA for the parton shower. The excited quark process is also modeled at LO in QCD with PYTHIA. Simulated samples are also produced for the background QCD process, using MADGRAPH5_aMC@NLO 2.6.5 at LO with up to four partons at the matrix element level. These background samples are used only for the development of the analysis strategy; the background estimation is performed using collision data. The NNPDF3.1 [44] parton distribution function (PDF) set at next-to-next-to-LO accuracy is used for all processes, along with the parton shower tune CP5 [45]. The interaction of particles with the detector is simulated using the GEANT4 toolkit [46]. The effects of pileup are incorporated by overlaying minimum bias events, simulated with PYTHIA,

on the hard scattering interactions, with the multiplicity distribution matching that observed in data.

In the MC simulations, only signal events with a generated resonance mass greater than 85% of the mass point under consideration are used. This requirement helps in avoiding sensitivity to the modeling of far-off-shell low-mass tails that are subject to large uncertainties originating from the choice of PDF, the assumption of resonance width, and the exact description of the relativistic Breit-Wigner line shape. The acceptance of this requirement, which depends on the width of the resonance, is denoted by $\mathcal{A}$.

Events are selected using a set of triggers requiring large hadronic activity, based on either the p_T of a single jet or the event H_T , defined as the scalar sum of jet p_T for all jets with $p_T > 30$ GeV and $|\eta| < 3.0$. The typical thresholds deployed for H_T triggers were 900 (1050) GeV and 360 (500) GeV for the p_T triggers in 2016 (2017–2018). To avoid the phase space where the trigger is inefficient, we apply a selection of $m_{jjj} > 1.50$ TeV (2016) or $m_{jjj} > 1.76$ TeV (2017–2018). The trigger and m_{jjj} thresholds for 2016 data are lower because of the lower instantaneous luminosity. The fully reconstructed events are required to have at least three reconstructed jets with $p_T > 100$ GeV and $|\eta| < 2.5$, prior to the application of the wide-jet algorithm. The aforementioned jet identification criteria, which have an efficiency of more than 98–99%, are applied to each jet used to construct the wide jet. Since the QCD background tends to have jets with higher values of $|\eta|$ than the signal processes, the sensitivity of the analysis is improved by requiring that the maximum difference in η between any two of the three wide jets, $\Delta\eta_{\max}$, be less than 1.6. Similarly, the maximum difference in ΔR between any two jets must be less than 3.0. The efficiencies of the above selection requirements are calculated in simulation to be $\sim 30\%$ for the Z_R process and between 23–35% (26–37%) for the $G_{KK}(q^*)$ process between 2.0 and 9.0 TeV. For resonance masses below ~ 2 TeV, the efficiencies decrease because the signal peaks are truncated by the lower m_{jjj} threshold, reaching $\sim 20\%$ for the Z_R process, $\sim 15\%$ for the G_{KK} process, and $\sim 25\%$ for the q^* process.

The background, dominantly arising from the QCD multijet process, is modeled by fitting the m_{jjj} distribution with smoothly falling, empirical functions that have been used in previous searches [19, 20]. Three families of functions are considered:

$$\begin{aligned} f_A(x; N) &= p_0 \frac{(1-x)^{p_1}}{x^{\sum_{i=2}^N p_i \log^{i-2}(x)}}, \\ f_B(x; N) &= p_0 \frac{e^{-p_1 x}}{x^{\sum_{i=2}^N p_i \log^{i-2}(x)}}, \\ f_C(x; N) &= p_0 x^{\sum_{i=1}^N p_i \log^{i-1}(x)}, \end{aligned} \tag{1}$$

where $x = m_{jjj}/\sqrt{s}$, $p_i (i = 0, 1, \dots, N)$ are free parameters, and N is the order of the fit function, which is determined using a Fisher F-test [47]. It is found that the third order yields the best fits to the data for all three families of functions. The signal is extracted using a binned maximum-likelihood fit based on a profile likelihood ratio test statistic [48], with variable bin size approximating the m_{jjj} resolution. The signal component of the fit is taken from binned templates derived from simulation, while the background component corresponds to the integral of the fit function across each bin. The three fit functions are incorporated into a single likelihood using a discrete nuisance parameter [49]. This parameter is profiled in an analogous way to continuous nuisance parameters to consider the extra uncertainty brought by the various possible choices of the background functional form. The background parameters are included as freely floating nuisance parameters. The three data-taking years are treated separately in the fit, in particular with independent background parameterizations for each year,

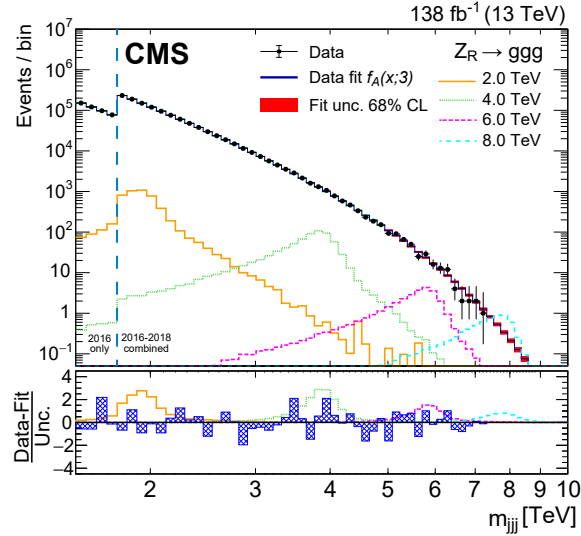


Figure 1: The observed m_{jjj} distribution and the background-only fit to the data using the f_A fit function. Uncertainties in the fit that correspond to the 68% confidence level are depicted with the red band. The expected m_{jjj} distributions for Z_R signal masses of 2.0, 4.0, 6.0, and 8.0 TeV, with nominal width of $\sim 3\%$, are also shown. For illustration purposes, the normalizations correspond to $\sigma\mathcal{B}$ values of 200, 50, 20, and 20 fb, respectively. Only 2016 data are shown for $m_{jjj} < 1.76$ TeV because of the higher trigger thresholds in 2017 and 2018. In the bottom panel, the blue hatched bars show the difference between the observed data and the background prediction divided by the statistical uncertainty, along with expectations for the example Z_R signal points.

and are combined statistically into a single likelihood. The combined result of the background-only fits with the function $f_A(x; 3)$ is shown in Fig. 1 as an example.

To gauge the systematic biases in the signal extraction, bias tests are performed using pseudo-data produced from the background-only fits with injected signal events. The bias is found to be negligible. Systematic uncertainties in the signal are incorporated in the likelihood as constrained nuisance parameters. The jet energy resolution and scale uncertainties translate to $\approx 10\%$ and $\approx 1\%$ uncertainties in the width and the position of the trijet mass shape for the signal, respectively. The integrated luminosities for the 2016, 2017, and 2018 data-taking years have 1.2–2.5% individual uncertainties [50–52], while the overall uncertainty for the 2016–2018 period is 1.6%. Finally, the uncertainty in the predicted signal yield that stems from the pileup reweighting procedure is $\approx 3\%$ for the three years, and is uncorrelated between years. The most significant uncertainties are those associated with background fit parameters in the background prediction, which are evaluated to have an impact on the fitted signal strength by $\approx 10\%$ over the entire search range.

No significant excesses compatible with the signal hypotheses are observed. For the $X \rightarrow ggg$ signal scenario, the largest deviation is observed at $m_X = 4.1$ TeV, corresponding to local significance values of 2.1 and 2.2 standard deviations for the nominal and narrow width hypotheses, respectively. Taking into account the look-elsewhere effect [53], the global significance values are 0.2 and 0.3 standard deviations, respectively. For the $X \rightarrow Y(gg)g$ cascade decay scenario, the largest deviation is observed for $\rho_m = 0.3$ and $m_X = 4.1$ TeV, with a local (global) significance of 2.2 (0.4) standard deviations. For the $X \rightarrow Y(qq)q$ cascade decay scenario, the largest deviation is observed at $\rho_m = 0.7$ and $m_X = 3.9$ TeV, and corresponds to a local (global) significance of 2.1 (0.3) standard deviations.

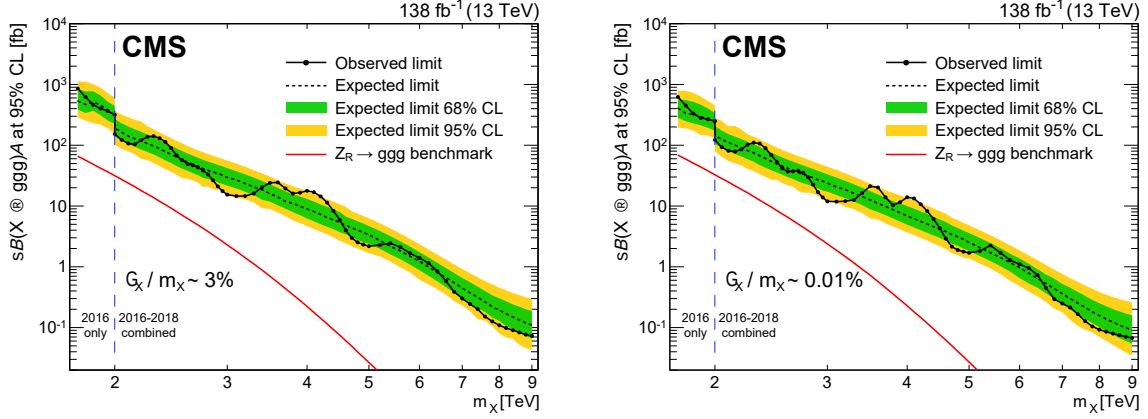


Figure 2: Limits at 95% CL on $\sigma\mathcal{B}(X \rightarrow ggg)\mathcal{A}$ for the nominal (left) and narrow-width (right) scenarios. Only 2016 data are used to derive limits below 2.0 TeV because of the higher trigger thresholds in 2017 and 2018. Theoretical predictions assuming SM-like couplings are depicted with red curves.

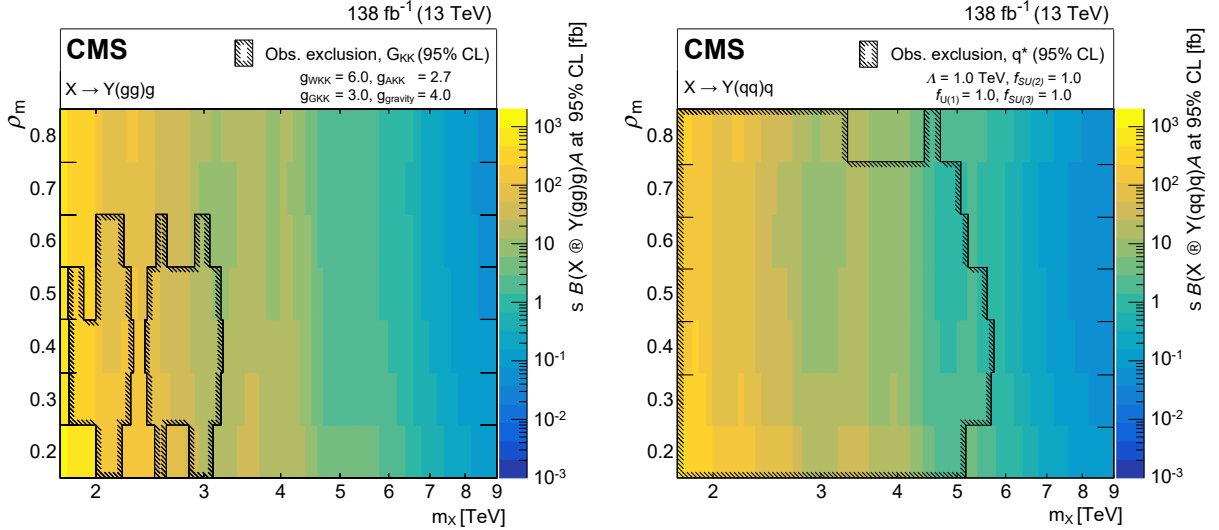


Figure 3: Observed limits at 95% CL as a function of m_X and ρ_m on $\sigma\mathcal{B}(X \rightarrow Y(gg)g)\mathcal{A}$ (left) and $\sigma\mathcal{B}(X \rightarrow Y(qq)q)\mathcal{A}$ (right). Only 2016 data are used to derive limits below 2.0 TeV because of the higher trigger thresholds in 2017 and 2018. The legend shows the model parameters for the chosen benchmark [23, 24], and their corresponding mass exclusion ranges are depicted with areas inside the black hatched contours.

Limits are set for the first time at 95% confidence level (CL) on the product of the signal process production cross section, branching fraction, and acceptance ($\sigma\mathcal{B}\mathcal{A}$) to three resolved jets using the asymptotic CL_s procedure [54–56] in the mass range between 1.75 TeV and 9.00 TeV. Only 2016 data are used to derive limits below 2.0 TeV because of the higher trigger thresholds deployed in 2017 and 2018. Limits are shown in Fig. 2 for the three-body decay of new resonances, ranging from 0.073 (0.067) fb to 0.86 (0.62) pb for the nominal (narrow) width hypotheses. Theoretical predictions of $\sigma\mathcal{B}\mathcal{A}$ assuming SM-like couplings are depicted with red curves. The current data set does not provide sufficient sensitivity to constrain the Z_R model. Limits on the cascade decay of new resonances into ggg (qqq) are presented in Fig. 3, which range from 0.061 (0.047) fb to 1.2 (0.49) pb. Benchmark model coupling parameters [23, 24] are shown in the legend, and their corresponding mass exclusion ranges are depicted with areas inside the black hatched contours. At $\rho_m = 0.2$, this search achieves a similar level of sensitivity to that of the previous search [20] for the ggg decay mode.

In summary, the first generic search for new particles decaying to three hadronic jets has been presented. The search uses proton-proton collision data at $\sqrt{s} = 13$ TeV recorded by the CMS experiment in 2016–2018, corresponding to an integrated luminosity of 138 fb^{-1} . The three-jet invariant mass spectrum is scanned for narrow peaks corresponding to new particles. No significant excesses above the standard model background expectations are observed. Limits are set on the product of the production cross section, branching fraction, and acceptance to three resolved jets. The results are interpreted in the context of a new right-handed boson Z_R decaying to three gluons, a Kaluza–Klein gluon G_{KK} decaying via an intermediate radion to three gluons (ggg), and an excited quark decaying via a vector boson to three quarks (qqq). This is the first search for the three-body decay of high-mass resonances (X) into three resolved jets at the LHC, and also the first search for X that decays into three resolved jets through an intermediate resonance (Y) with a mass ratio m_Y/m_X between 0.3–0.8 for the ggg decay mode and 0.2–0.8 for the qqq decay mode, significantly extending the model parameter space explored by a previous search [20].

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A The CMS Collaboration

Yerevan Physics Institute, Yerevan, Armenia

A. Hayrapetyan, A. Tumasyan¹ 

Institut für Hochenergiephysik, Vienna, Austria

W. Adam , J.W. Andrejkovic, T. Bergauer , S. Chatterjee , K. Damanakis , M. Dragicevic , A. Escalante Del Valle , P.S. Hussain , M. Jeitler² , N. Krammer , D. Liko , I. Mikulec , J. Schieck² , R. Schöffbeck , D. Schwarz , M. Sonawane , S. Templ , W. Waltenberger , C.-E. Wulz² 

Universiteit Antwerpen, Antwerpen, Belgium

M.R. Darwish³ , T. Janssen , P. Van Mechelen 

Vrije Universiteit Brussel, Brussel, Belgium

E.S. Bols , J. D'Hondt , S. Dansana , A. De Moor , M. Delcourt , H. El Faham , S. Lowette , I. Makarenko , D. Müller , A.R. Sahasransu , S. Tavernier , M. Tytgat⁴ , S. Van Putte , D. Vannerom 

Université Libre de Bruxelles, Bruxelles, Belgium

B. Clerbaux , G. De Lentdecker , L. Favart , D. Hohov , J. Jaramillo , A. Khalilzadeh, K. Lee , M. Mahdavihorrami , A. Malara , S. Paredes , L. Pétré , N. Postiau, L. Thomas , M. Vanden Bemden , C. Vander Velde , P. Vanlaer 

Ghent University, Ghent, Belgium

M. De Coen , D. Dobur , Y. Hong , J. Knolle , L. Lambrecht , G. Mestdach, C. Rendón, A. Samalan, K. Skovpen , N. Van Den Bossche , L. Wezenbeek 

Université Catholique de Louvain, Louvain-la-Neuve, Belgium

A. Benecke , G. Bruno , C. Caputo , C. Delaere , I.S. Donertas , A. Giammanco , K. Jaffel , Sa. Jain , V. Lemaitre, J. Lidrych , P. Mastrapasqua , K. Mondal , T.T. Tran , S. Wertz 

Centro Brasileiro de Pesquisas Fisicas, Rio de Janeiro, Brazil

G.A. Alves , E. Coelho , C. Hensel , T. Menezes De Oliveira, A. Moraes , P. Rebello Teles , M. Soeiro

Universidade do Estado do Rio de Janeiro, Rio de Janeiro, Brazil

W.L. Aldá Júnior , M. Alves Gallo Pereira , M. Barroso Ferreira Filho , H. Brandao Malbouisson , W. Carvalho , J. Chinellato⁵, E.M. Da Costa , G.G. Da Silveira⁶ , D. De Jesus Damiao , S. Fonseca De Souza , J. Martins⁷ , C. Mora Herrera , K. Mota Amarilo , L. Mundim , H. Nogima , A. Santoro , A. Sznajder , M. Thiel , A. Vilela Pereira 

Universidade Estadual Paulista, Universidade Federal do ABC, São Paulo, Brazil

C.A. Bernardes⁶ , L. Calligaris , T.R. Fernandez Perez Tomei , E.M. Gregores , P.G. Mercadante , S.F. Novaes , B. Orzari , Sandra S. Padula 

Institute for Nuclear Research and Nuclear Energy, Bulgarian Academy of Sciences, Sofia, Bulgaria

A. Aleksandrov , G. Antchev , R. Hadjiiska , P. Iaydjiev , M. Misheva , M. Shopova , G. Sultanov 

University of Sofia, Sofia, Bulgaria

A. Dimitrov , T. Ivanov , L. Litov , B. Pavlov , P. Petkov , A. Petrov , E. Shumka 

Instituto De Alta Investigación, Universidad de Tarapacá, Casilla 7 D, Arica, Chile

S. Keshri , S. Thakur 

Beihang University, Beijing, China

T. Cheng , Q. Guo, T. Javaid , M. Mittal , L. Yuan 

Department of Physics, Tsinghua University, Beijing, China

G. Bauer^{8,9}, Z. Hu , J. Liu, K. Yi^{8,10} 

Institute of High Energy Physics, Beijing, China

G.M. Chen¹¹ , H.S. Chen¹¹ , M. Chen¹¹ , F. Iemmi , C.H. Jiang, A. Kapoor¹² ,
H. Liao , Z.-A. Liu¹³ , F. Monti , M.A. Shahzad¹¹, R. Sharma¹⁴ , J.N. Song¹³, J. Tao ,
C. Wang¹¹, J. Wang , Z. Wang¹¹, H. Zhang 

State Key Laboratory of Nuclear Physics and Technology, Peking University, Beijing, China

A. Agapitos , Y. Ban , A. Levin , C. Li , Q. Li , Y. Mao, S.J. Qian , X. Sun ,
D. Wang , H. Yang, L. Zhang , C. Zhou 

Sun Yat-Sen University, Guangzhou, China

Z. You 

University of Science and Technology of China, Hefei, China

N. Lu 

Institute of Modern Physics and Key Laboratory of Nuclear Physics and Ion-beam Application (MOE) - Fudan University, Shanghai, China

X. Gao¹⁵ , D. Leggat, H. Okawa , Y. Zhang 

Zhejiang University, Hangzhou, Zhejiang, China

Z. Lin , C. Lu , M. Xiao 

Universidad de Los Andes, Bogota, Colombia

C. Avila , D.A. Barbosa Trujillo, A. Cabrera , C. Florez , J. Fraga , J.A. Reyes Vega

Universidad de Antioquia, Medellin, Colombia

J. Mejia Guisao , F. Ramirez , M. Rodriguez , J.D. Ruiz Alvarez 

University of Split, Faculty of Electrical Engineering, Mechanical Engineering and Naval Architecture, Split, Croatia

D. Giljanovic , N. Godinovic , D. Lelas , A. Sculac 

University of Split, Faculty of Science, Split, Croatia

M. Kovac , T. Sculac 

Institute Rudjer Boskovic, Zagreb, Croatia

P. Bargassa , V. Brigljevic , B.K. Chitroda , D. Ferencek , S. Mishra ,
A. Starodumov¹⁶ , T. Susa 

University of Cyprus, Nicosia, Cyprus

A. Attikis , K. Christoforou , S. Konstantinou , J. Mousa , C. Nicolaou, F. Ptochos ,
P.A. Razis , H. Rykaczewski, H. Saka , A. Stepennov 

Charles University, Prague, Czech Republic

M. Finger , M. Finger Jr. , A. Kveton 

Escuela Politecnica Nacional, Quito, Ecuador

E. Ayala 

Universidad San Francisco de Quito, Quito, EcuadorE. Carrera Jarrin **Academy of Scientific Research and Technology of the Arab Republic of Egypt, Egyptian Network of High Energy Physics, Cairo, Egypt**S. Elgammal¹⁷, A. Ellithi Kamel¹⁸**Center for High Energy Physics (CHEP-FU), Fayoum University, El-Fayoum, Egypt**M.A. Mahmoud , Y. Mohammed **National Institute of Chemical Physics and Biophysics, Tallinn, Estonia**R.K. Dewanjee¹⁹ , K. Ehataht , M. Kadastik, T. Lange , S. Nandan , C. Nielsen , J. Pata , M. Raidal , L. Tani , C. Veelken **Department of Physics, University of Helsinki, Helsinki, Finland**H. Kirschenmann , K. Osterberg , M. Voutilainen **Helsinki Institute of Physics, Helsinki, Finland**S. Bharthuar , E. Brücken , F. Garcia , J. Havukainen , K.T.S. Kallonen , M.S. Kim , R. Kinnunen, T. Lampén , K. Lassila-Perini , S. Lehti , T. Lindén , M. Lotti, L. Martikainen , M. Myllymäki , M.m. Rantanen , H. Siikonen , E. Tuominen , J. Tuominiemi **Lappeenranta-Lahti University of Technology, Lappeenranta, Finland**P. Luukka , H. Petrow , T. Tuuva[†]**IRFU, CEA, Université Paris-Saclay, Gif-sur-Yvette, France**M. Besancon , F. Couderc , M. Dejardin , D. Denegri, J.L. Faure, F. Ferri , S. Ganjour , P. Gras , G. Hamel de Monchenault , V. Lohezic , J. Malcles , J. Rander, A. Rosowsky , M.Ö. Sahin , A. Savoy-Navarro²⁰ , P. Simkina , M. Titov , M. Tornago **Laboratoire Leprince-Ringuet, CNRS/IN2P3, Ecole Polytechnique, Institut Polytechnique de Paris, Palaiseau, France**C. Baldenegro Barrera , F. Beaudette , A. Buchot Perraguin , P. Busson , A. Cappati , C. Charlot , F. Damas , O. Davignon , A. De Wit , G. Falmagne , B.A. Fontana Santos Alves , S. Ghosh , A. Gilbert , R. Granier de Cassagnac , A. Hakimi , B. Harikrishnan , L. Kalipoliti , G. Liu , J. Motta , M. Nguyen , C. Ochando , L. Portales , R. Salerno , U. Sarkar , J.B. Sauvan , Y. Sirois , A. Tarabini , E. Vernazza , A. Zabi , A. Zghiche **Université de Strasbourg, CNRS, IPHC UMR 7178, Strasbourg, France**J.-L. Agram²¹ , J. Andrea , D. Apparú , D. Bloch , J.-M. Brom , E.C. Chabert , C. Collard , S. Falke , U. Goerlach , C. Grimault, R. Haeberle , A.-C. Le Bihan , M.A. Sessini , P. Van Hove **Institut de Physique des 2 Infinis de Lyon (IP2I), Villeurbanne, France**S. Beauceron , B. Blancon , G. Boudoul , N. Chanon , J. Choi , D. Contardo , P. Depasse , C. Dozen²² , H. El Mamouni, J. Fay , S. Gascon , M. Gouzevitch , C. Greenberg, G. Grenier , B. Ille , I.B. Laktineh, M. Lethuillier , L. Mirabito, S. Perries, A. Purohit , M. Vander Donckt , P. Verdier , J. Xiao **Georgian Technical University, Tbilisi, Georgia**I. Lomidze , T. Toriashvili²³ , Z. Tsamalaidze¹⁶ **RWTH Aachen University, I. Physikalisches Institut, Aachen, Germany**

V. Botta , L. Feld , K. Klein , M. Lipinski , D. Meuser , A. Pauls , N. Röwert , M. Teroerde 

RWTH Aachen University, III. Physikalisches Institut A, Aachen, Germany

S. Diekmann , A. Dodonova , N. Eich , D. Eliseev , F. Engelke , M. Erdmann , P. Fackeldey , B. Fischer , T. Hebbeker , K. Hoepfner , F. Ivone , A. Jung , M.y. Lee , L. Mastrolorenzo, M. Merschmeyer , A. Meyer , S. Mukherjee , D. Noll , A. Novak , F. Nowotny, A. Pozdnyakov , Y. Rath, W. Redjeb , F. Rehm, H. Reithler , V. Sarkisovi , A. Schmidt , A. Sharma , A. Stein , F. Torres Da Silva De Araujo²⁴ , L. Vigilante, S. Wiedenbeck , S. Zaleski

RWTH Aachen University, III. Physikalisches Institut B, Aachen, Germany

C. Dziwok , G. Flügge , W. Haj Ahmad²⁵ , T. Kress , A. Nowack , O. Pooth , A. Stahl , T. Ziemons , A. Zotz 

Deutsches Elektronen-Synchrotron, Hamburg, Germany

H. Aarup Petersen , M. Aldaya Martin , J. Alimena , S. Amoroso, Y. An , S. Baxter , M. Bayatmakou , H. Becerril Gonzalez , O. Behnke , A. Belvedere , S. Bhattacharya , F. Blekman²⁶ , K. Borras²⁷ , D. Brunner , A. Campbell , A. Cardini , C. Cheng, F. Colombina , S. Consuegra Rodríguez , G. Correia Silva , M. De Silva , G. Eckerlin, D. Eckstein , L.I. Estevez Banos , O. Filatov , E. Gallo²⁶ , A. Geiser , A. Giraldi , G. Greau, V. Guglielmi , M. Guthoff , A. Hinzmann , A. Jafari²⁸ , L. Jeppe , N.Z. Jomhari , B. Kaech , M. Kasemann , H. Kaveh , C. Kleinwort , R. Kogler , M. Komm , D. Krücker , W. Lange, D. Leyva Pernia , K. Lipka²⁹ , W. Lohmann³⁰ , R. Mankel , I.-A. Melzer-Pellmann , M. Mendizabal Morentin , J. Metwally, A.B. Meyer , G. Milella , A. Mussgiller , A. Nürnberg , Y. Otari, D. Pérez Adán , E. Ranken , A. Raspereza , B. Ribeiro Lopes , J. Rübenach, A. Saggio , M. Scham^{31,27} , S. Schnake²⁷ , P. Schütze , C. Schwanenberger²⁶ , D. Selivanova , M. Shchedrolosiev , R.E. Sosa Ricardo , L.P. Sreelatha Pramod , D. Stafford, F. Vazzoler , A. Ventura Barroso , R. Walsh , Q. Wang , Y. Wen , K. Wichmann, L. Wiens²⁷ , C. Wissing , S. Wuchterl , Y. Yang , A. Zimmermann Castro Santos

University of Hamburg, Hamburg, Germany

A. Albrecht , S. Albrecht , M. Antonello , S. Bein , L. Benato , M. Bonanomi , P. Connor , M. Eich, K. El Morabit , Y. Fischer , A. Fröhlich, C. Garbers , E. Garutti , A. Grohsjean , M. Hajheidari, J. Haller , H.R. Jabusch , G. Kasieczka , P. Keicher, R. Klanner , W. Korcar , T. Kramer , V. Kutzner , F. Labe , J. Lange , A. Lobanov , C. Matthies , A. Mehta , L. Moureaux , M. Mrowietz, A. Nigamova , Y. Nissan, A. Paasch , K.J. Pena Rodriguez , T. Quadfasel , B. Raciti , M. Rieger , D. Savoia , J. Schindler , P. Schleper , M. Schröder , J. Schwandt , M. Sommerhalder , H. Stadie , G. Steinbrück , A. Tews, M. Wolf

Karlsruher Institut fuer Technologie, Karlsruhe, Germany

S. Brommer , M. Burkart, E. Butz , T. Chwalek , A. Dierlamm , A. Droll, N. Faltermann , M. Giffels , A. Gottmann , F. Hartmann³² , R. Hofsaess , M. Horzela , U. Husemann , M. Klute , R. Koppenhöfer , M. Link, A. Lintuluoto , S. Maier , S. Mitra , M. Mormile , Th. Müller , M. Neukum, M. Oh , G. Quast , K. Rabbertz , B. Regnery , N. Shadskiy , I. Shvetsov , H.J. Simonis , N. Trevisani , R. Ulrich , J. van der Linden , R.F. Von Cube , M. Wassmer , S. Wieland , F. Wittig, R. Wolf , S. Wunsch, X. Zuo

Institute of Nuclear and Particle Physics (INPP), NCSR Demokritos, Aghia Paraskevi,

Greece

G. Anagnostou, P. Assiouras , G. Daskalakis , A. Kyriakis, A. Papadopoulos³², A. Stakia 

National and Kapodistrian University of Athens, Athens, Greece

P. Kontaxakis , G. Melachroinos, A. Panagiotou, I. Papavergou , I. Paraskevas , N. Saoulidou , K. Theofilatos , E. Tziaferi , K. Vellidis , I. Zisopoulos 

National Technical University of Athens, Athens, Greece

G. Bakas , T. Chatzistavrou, G. Karapostoli , K. Kousouris , I. Papakrivopoulos , E. Siamarkou, G. Tsiopolitis, A. Zacharopoulou

University of Ioánnina, Ioánnina, Greece

K. Adamidis, I. Bestintzanos, I. Evangelou , C. Foudas, P. Gianneios , C. Kamtsikis, P. Katsoulis, P. Kokkas , P.G. Kosmoglou Kioseoglou , N. Manthos , I. Papadopoulos , J. Strologas 

HUN-REN Wigner Research Centre for Physics, Budapest, Hungary

M. Bartók³³ , C. Hajdu , D. Horvath^{34,35} , F. Sikler , V. Veszpremi 

MTA-ELTE Lendület CMS Particle and Nuclear Physics Group, Eötvös Loránd University, Budapest, Hungary

M. Csanád , K. Farkas , M.M.A. Gadallah³⁶ , Á. Kadlecik , P. Major , K. Mandal , G. Pásztor , A.J. Rádl³⁷ , G.I. Veres 

Faculty of Informatics, University of Debrecen, Debrecen, Hungary

P. Raics, B. Ujvari³⁸ , G. Zilizi 

Institute of Nuclear Research ATOMKI, Debrecen, Hungary

G. Bencze, S. Czellar, J. Karancsi³³ , J. Molnar, Z. Szillasi

Karoly Robert Campus, MATE Institute of Technology, Gyongyos, Hungary

T. Csorgo³⁷ , F. Nemes³⁷ , T. Novak 

Panjab University, Chandigarh, India

J. Babbar , S. Bansal , S.B. Beri, V. Bhatnagar , G. Chaudhary , S. Chauhan , N. Dhingra³⁹ , A. Kaur , A. Kaur , H. Kaur , M. Kaur , S. Kumar , M. Meena , K. Sandeep , T. Sheokand, J.B. Singh , A. Singla 

University of Delhi, Delhi, India

A. Ahmed , A. Bhardwaj , A. Chhetri , B.C. Choudhary , A. Kumar , M. Naimuddin , K. Ranjan , S. Saumya 

Saha Institute of Nuclear Physics, HBNI, Kolkata, India

S. Acharya⁴⁰ , S. Baradia , S. Barman⁴¹ , S. Bhattacharya , D. Bhowmik, S. Dutta , S. Dutta, B. Gomber⁴⁰ , P. Palit , G. Saha , B. Sahu⁴⁰ , S. Sarkar

Indian Institute of Technology Madras, Madras, India

M.M. Ameen , P.K. Behera , S.C. Behera , S. Chatterjee , P. Jana , P. Kalbhor , J.R. Komaragiri⁴² , D. Kumar⁴² , L. Panwar⁴² , R. Pradhan , P.R. Pujahari , N.R. Saha , A. Sharma , A.K. Sikdar , S. Verma 

Tata Institute of Fundamental Research-A, Mumbai, India

T. Aziz, I. Das , S. Dugad, M. Kumar , G.B. Mohanty , P. Suryadevara

Tata Institute of Fundamental Research-B, Mumbai, India

A. Bala , S. Banerjee , R.M. Chatterjee, M. Guchait , Sh. Jain , S. Karmakar 

S. Kumar , G. Majumder , K. Mazumdar , S. Mukherjee , S. Parolia , A. Thachayath 

National Institute of Science Education and Research, An OCC of Homi Bhabha National Institute, Bhubaneswar, Odisha, India

S. Bahinipati⁴³ , A.K. Das , C. Kar , D. Maity⁴⁴ , P. Mal , T. Mishra , V.K. Muraleedharan Nair Bindhu⁴⁴ , K. Naskar⁴⁴ , A. Nayak⁴⁴ , P. Sadangi , P. Saha , S.K. Swain , S. Varghese⁴⁴ , D. Vats⁴⁴ 

Indian Institute of Science Education and Research (IISER), Pune, India

A. Alpana , S. Dube , B. Kansal , A. Laha , A. Rastogi , S. Sharma 

Isfahan University of Technology, Isfahan, Iran

H. Bakhshiansohi⁴⁵ , E. Khazaie⁴⁶ , M. Zeinali⁴⁷ 

Institute for Research in Fundamental Sciences (IPM), Tehran, Iran

S. Chenarani⁴⁸ , S.M. Etesami , M. Khakzad , M. Mohammadi Najafabadi 

University College Dublin, Dublin, Ireland

M. Grunewald 

INFN Sezione di Bari^a, Università di Bari^b, Politecnico di Bari^c, Bari, Italy

M. Abbrescia^{a,b} , R. Aly^{a,c,49} , A. Colaleo^{a,b} , D. Creanza^{a,c} , B. D'Anzi^{a,b} , N. De Filippis^{a,c} , M. De Palma^{a,b} , A. Di Florio^{a,c} , W. Elmetenawee^{a,b,49} , L. Fiore^a , G. Iaselli^{a,c} , G. Maggi^{a,c} , M. Maggi^a , I. Margjeka^{a,b} , V. Mastrapasqua^{a,b} , S. My^{a,b} , S. Nuzzo^{a,b} , A. Pellecchia^{a,b} , A. Pompili^{a,b} , G. Pugliese^{a,c} , R. Radogna^a , G. Ramirez-Sanchez^{a,c} , D. Ramos^a , A. Ranieri^a , L. Silvestris^a , F.M. Simone^{a,b} , Ü. Sözbilir^a , A. Stamerra^a , R. Venditti^a , P. Verwilligen^a , A. Zaza^{a,b} 

INFN Sezione di Bologna^a, Università di Bologna^b, Bologna, Italy

G. Abbiendi^a , C. Battilana^{a,b} , D. Bonacorsi^{a,b} , L. Borgonovi^a , R. Campanini^{a,b} , P. Capiluppi^{a,b} , A. Castro^{a,b} , F.R. Cavallo^a , M. Cuffiani^{a,b} , G.M. Dallavalle^a , T. Diotallevi^{a,b} , F. Fabbri^a , D. Fasanella^{a,b} , P. Giacomelli^a , L. Giommi^{a,b} , C. Grandi^a , L. Guiducci^{a,b} , S. Lo Meo^{a,50} , L. Lunerti^{a,b} , S. Marcellini^a , G. Masetti^a , F.L. Navarria^{a,b} , A. Perrotta^a , F. Primavera^{a,b} , A.M. Rossi^{a,b} , T. Rovelli^{a,b} , G.P. Siroli^{a,b} 

INFN Sezione di Catania^a, Università di Catania^b, Catania, Italy

S. Costa^{a,b,51} , A. Di Mattia^a , R. Potenza^{a,b} , A. Tricomi^{a,b,51} , C. Tuve^{a,b} 

INFN Sezione di Firenze^a, Università di Firenze^b, Firenze, Italy

G. Barbagli^a , G. Bardelli^{a,b} , B. Camaiani^{a,b} , A. Cassese^a , R. Ceccarelli^a , V. Ciulli^{a,b} , C. Civinini^a , R. D'Alessandro^{a,b} , E. Focardi^{a,b} , T. Kello^a , G. Latino^{a,b} , P. Lenzi^{a,b} , M. Lizzo^a , M. Meschini^a , S. Paoletti^a , A. Papanastassiou^{a,b} , G. Sguazzoni^a , L. Viliani^a 

INFN Laboratori Nazionali di Frascati, Frascati, Italy

L. Benussi , S. Bianco , S. Meola⁵² , D. Piccolo 

INFN Sezione di Genova^a, Università di Genova^b, Genova, Italy

P. Chatagnon^a , F. Ferro^a , E. Robutti^a , S. Tosi^{a,b} 

INFN Sezione di Milano-Bicocca^a, Università di Milano-Bicocca^b, Milano, Italy

A. Benaglia^a , G. Boldrini^{a,b} , F. Brivio^a , F. Cetorelli^a , F. De Guio^{a,b} , M.E. Dinardo^{a,b} , P. Dini^a , S. Gennai^a , R. Gerosa^{a,b} , A. Ghezzi^{a,b} , P. Govoni^{a,b} 

L. Guzzi^a , M.T. Lucchini^{a,b} , M. Malberti^a , S. Malvezzi^a , A. Massironi^a , D. Menasce^a , L. Moroni^a , M. Paganoni^{a,b} , D. Pedrini^a , B.S. Pinolini^a, S. Ragazzi^{a,b} , T. Tabarelli de Fatis^{a,b} , D. Zuolo^a 

INFN Sezione di Napoli^a, Università di Napoli 'Federico II'^b, Napoli, Italy; Università della Basilicata^c, Potenza, Italy; Università G. Marconi^d, Roma, Italy

S. Buontempo^a , A. Cagnotta^{a,b} , F. Carnevali^{a,b}, N. Cavallo^{a,c} , A. De Iorio^{a,b} , F. Fabozzi^{a,c} , A.O.M. Iorio^{a,b} , L. Lista^{a,b,53} , P. Paolucci^{a,32} , B. Rossi^a , C. Sciacca^{a,b} 

INFN Sezione di Padova^a, Università di Padova^b, Padova, Italy; Università di Trento^c, Trento, Italy

R. Ardino^a , P. Azzi^a , N. Bacchetta^{a,54} , P. Bortignon^a , A. Bragagnolo^{a,b} , R. Carlin^{a,b} , P. Checchia^a , T. Dorigo^a , U. Gasparini^{a,b} , G. Grosso^a, M. Gulmini^{a,55} , L. Layer^{a,56}, E. Lusiani^a , M. Margoni^{a,b} , G. Maron^{a,55} , A.T. Meneguzzo^{a,b} , M. Migliorini^{a,b} , J. Pazzini^{a,b} , P. Ronchese^{a,b} , R. Rossin^{a,b} , F. Simonetto^{a,b} , G. Strong^a , M. Tosi^{a,b} , A. Triossi^{a,b} , S. Ventura^a , H. Yarar^{a,b}, M. Zanetti^{a,b} , P. Zotto^{a,b} , A. Zucchetta^{a,b} , G. Zumerle^{a,b} 

INFN Sezione di Pavia^a, Università di Pavia^b, Pavia, Italy

S. Abu Zeid^{a,57} , C. Aimè^{a,b} , A. Braghieri^a , S. Calzaferri^{a,b} , D. Fiorina^{a,b} , P. Montagna^{a,b} , V. Re^a , C. Riccardi^{a,b} , P. Salvini^a , I. Vai^{a,b} , P. Vitulo^{a,b} 

INFN Sezione di Perugia^a, Università di Perugia^b, Perugia, Italy

S. Ajmal^{a,b} , P. Asenov^{a,58} , G.M. Bilei^a , D. Ciangottini^{a,b} , L. Fanò^{a,b} , M. Magherini^{a,b} , G. Mantovani^{a,b}, V. Mariani^{a,b} , M. Menichelli^a, F. Moscatelli^{a,58} , A. Piccinelli^{a,b} , M. Presilla^{a,b} , A. Rossi^{a,b} , A. Santocchia^{a,b} , D. Spiga^a , T. Tedeschi^{a,b} 

INFN Sezione di Pisa^a, Università di Pisa^b, Scuola Normale Superiore di Pisa^c, Pisa, Italy; Università di Siena^d, Siena, Italy

P. Azzurri^a , G. Bagliesi^a , R. Bhattacharya^a , L. Bianchini^{a,b} , T. Boccali^a , E. Bossini^a , D. Bruschini^{a,c} , R. Castaldi^a , M.A. Ciocci^{a,b} , M. Cipriani^{a,b} , V. D'Amante^{a,d} , R. Dell'Orso^a , S. Donato^a , A. Giassi^a , F. Ligabue^{a,c} , D. Matos Figueiredo^a , A. Messineo^{a,b} , M. Musich^{a,b} , F. Palla^a , A. Rizzi^{a,b} , G. Rolandi^{a,c} , S. Roy Chowdhury^a , T. Sarkar^a , A. Scribano^a , P. Spagnolo^a , R. Tenchini^{a,b} , G. Tonelli^{a,b} , N. Turini^{a,d} , A. Venturi^a , P.G. Verdini^a 

INFN Sezione di Roma^a, Sapienza Università di Roma^b, Roma, Italy

P. Barria^a , M. Campana^{a,b} , F. Cavallari^a , L. Cunqueiro Mendez^{a,b} , D. Del Re^{a,b} , E. Di Marco^a , M. Diemoz^a , F. Errico^{a,b} , E. Longo^{a,b} , P. Meridiani^a , J. Mijuskovic^{a,b} , G. Organtini^{a,b} , F. Pandolfi^a , R. Paramatti^{a,b} , C. Quaranta^{a,b} , S. Rahatlou^{a,b} , C. Rovelli^a , F. Santanastasio^{a,b} , L. Soffi^a 

INFN Sezione di Torino^a, Università di Torino^b, Torino, Italy; Università del Piemonte Orientale^c, Novara, Italy

N. Amapane^{a,b} , R. Arcidiacono^{a,c} , S. Argiro^{a,b} , M. Arneodo^{a,c} , N. Bartosik^a , R. Bellan^{a,b} , A. Bellora^{a,b} , C. Biino^a , N. Cartiglia^a , M. Costa^{a,b} , R. Covarelli^{a,b} , N. Demaria^a , L. Finco^a , M. Grippio^{a,b} , B. Kiani^{a,b} , F. Legger^a , F. Luongo^{a,b} , C. Mariotti^a , S. Maselli^a , A. Mecca^{a,b} , E. Migliore^{a,b} , M. Monteno^a , R. Mulargia^a , M.M. Obertino^{a,b} , G. Ortona^a , L. Pacher^{a,b} , N. Pastrone^a , M. Pelliccioni^a , M. Ruspa^{a,c} , F. Siviero^{a,b} , V. Sola^{a,b} , A. Solano^{a,b} , D. Soldi^{a,b} , A. Staiano^a , C. Tarricone^{a,b} , D. Trocino^a , G. Umoret^{a,b} , E. Vlasov^{a,b} 

INFN Sezione di Trieste^a, Università di Trieste^b, Trieste, Italy

S. Belforte^a , V. Candelise^{a,b} , M. Casarsa^a , F. Cossutti^a , K. De Leo^{a,b} ,
G. Della Ricca^{a,b} 

Kyungpook National University, Daegu, Korea

S. Dogra , J. Hong , C. Huh , B. Kim , D.H. Kim , J. Kim, H. Lee, S.W. Lee ,
C.S. Moon , Y.D. Oh , M.S. Ryu , S. Sekmen , Y.C. Yang 

Chonnam National University, Institute for Universe and Elementary Particles, Kwangju, Korea

G. Bak , P. Gwak , H. Kim , D.H. Moon 

Hanyang University, Seoul, Korea

E. Asilar , D. Kim , T.J. Kim , J.A. Merlin, J. Park 

Korea University, Seoul, Korea

S. Choi , S. Han, B. Hong , K. Lee, K.S. Lee , S. Lee , J. Park, S.K. Park, J. Yoo 

Kyung Hee University, Department of Physics, Seoul, Korea

J. Goh 

Sejong University, Seoul, Korea

H. S. Kim , Y. Kim, S. Lee

Seoul National University, Seoul, Korea

J. Almond, J.H. Bhyun, J. Choi , W. Jun , J. Kim , J.S. Kim, S. Ko , H. Kwon , H. Lee ,
J. Lee , J. Lee , B.H. Oh , S.B. Oh , H. Seo , U.K. Yang, I. Yoon 

University of Seoul, Seoul, Korea

W. Jang , D.Y. Kang, Y. Kang , S. Kim , B. Ko, J.S.H. Lee , Y. Lee , I.C. Park , Y. Roh,
I.J. Watson , S. Yang 

Yonsei University, Department of Physics, Seoul, Korea

S. Ha , H.D. Yoo 

Sungkyunkwan University, Suwon, Korea

M. Choi , M.R. Kim , H. Lee, Y. Lee , I. Yu 

College of Engineering and Technology, American University of the Middle East (AUM), Dasman, Kuwait

T. Beyrouthy, Y. Maghrbi 

Riga Technical University, Riga, Latvia

K. Dreimanis , A. Gaile , G. Pikurs, A. Potrebko , M. Seidel , V. Veckalns⁵⁹ 

University of Latvia (LU), Riga, Latvia

N.R. Strautnieks 

Vilnius University, Vilnius, Lithuania

M. Ambrozas , A. Juodagalvis , A. Rinkevicius , G. Tamulaitis 

National Centre for Particle Physics, Universiti Malaya, Kuala Lumpur, Malaysia

N. Bin Norjoharuddeen , I. Yusuff⁶⁰ , Z. Zolkapli

Universidad de Sonora (UNISON), Hermosillo, Mexico

J.F. Benitez , A. Castaneda Hernandez , H.A. Encinas Acosta, L.G. Gallegos Maríñez,
M. León Coello , J.A. Murillo Quijada , A. Sehrawat , L. Valencia Palomo 

Centro de Investigacion y de Estudios Avanzados del IPN, Mexico City, Mexico

G. Ayala , H. Castilla-Valdez , E. De La Cruz-Burelo , I. Heredia-De La Cruz⁶¹ ,
R. Lopez-Fernandez , C.A. Mondragon Herrera, A. Sánchez Hernández 

Universidad Iberoamericana, Mexico City, Mexico

C. Oropeza Barrera , M. Ramírez García 

Benemerita Universidad Autonoma de Puebla, Puebla, Mexico

I. Bautista , I. Pedraza , H.A. Salazar Ibarguen , C. Uribe Estrada 

University of Montenegro, Podgorica, Montenegro

I. Bubanja, N. Raicevic 

University of Canterbury, Christchurch, New Zealand

P.H. Butler 

National Centre for Physics, Quaid-I-Azam University, Islamabad, Pakistan

A. Ahmad , M.I. Asghar, A. Awais , M.I.M. Awan, H.R. Hoorani , W.A. Khan 

AGH University of Krakow, Faculty of Computer Science, Electronics and Telecommunications, Krakow, Poland

V. Avati, L. Grzanka , M. Malawski 

National Centre for Nuclear Research, Swierk, Poland

H. Bialkowska , M. Bluj , B. Boimska , M. Górski , M. Kazana , M. Szleper ,
P. Zalewski 

Institute of Experimental Physics, Faculty of Physics, University of Warsaw, Warsaw, Poland

K. Bunkowski , K. Doroba , A. Kalinowski , M. Konecki , J. Krolikowski ,
A. Muhammad 

Warsaw University of Technology, Warsaw, Poland

K. Pozniak , W. Zabolotny 

Laboratório de Instrumentação e Física Experimental de Partículas, Lisboa, Portugal

M. Araujo , D. Bastos , C. Beirão Da Cruz E Silva , A. Boletti , M. Bozzo , P. Faccioli ,
M. Gallinaro , J. Hollar , N. Leonardo , T. Niknejad , A. Petrilli , M. Pisano ,
J. Seixas , J. Varela , J.W. Wulff

Faculty of Physics, University of Belgrade, Belgrade, Serbia

P. Adzic , P. Milenovic 

VINCA Institute of Nuclear Sciences, University of Belgrade, Belgrade, Serbia

M. Dordevic , J. Milosevic , V. Rekovic

Centro de Investigaciones Energéticas Medioambientales y Tecnológicas (CIEMAT), Madrid, Spain

M. Aguilar-Benitez, J. Alcaraz Maestre , Cristina F. Bedoya , M. Cepeda , M. Cerrada , N. Colino , B. De La Cruz , A. Delgado Peris , D. Fernández Del Val ,
J.P. Fernández Ramos , J. Flix , M.C. Fouz , O. Gonzalez Lopez , S. Goy Lopez ,
J.M. Hernandez , M.I. Josa , J. León Holgado , D. Moran , C. M. Morcillo Perez ,
Á. Navarro Tobar , C. Perez Dengra , A. Pérez-Calero Yzquierdo , J. Puerta Pelayo ,
I. Redondo , D.D. Redondo Ferrero , L. Romero, S. Sánchez Navas , L. Urda Gómez ,
J. Vazquez Escobar , C. Willmott

Universidad Autónoma de Madrid, Madrid, Spain

J.F. de Trocóniz 

Universidad de Oviedo, Instituto Universitario de Ciencias y Tecnologías Espaciales de Asturias (ICTEA), Oviedo, Spain

B. Alvarez Gonzalez , J. Cuevas , J. Fernandez Menendez , S. Folgueras , I. Gonzalez Caballero , J.R. González Fernández , E. Palencia Cortezon , C. Ramón Álvarez , V. Rodríguez Bouza , A. Soto Rodríguez , A. Trapote , C. Vico Villalba , P. Vischia 

Instituto de Física de Cantabria (IFCA), CSIC-Universidad de Cantabria, Santander, Spain

S. Bhowmik , S. Blanco Fernández , J.A. Brochero Cifuentes , I.J. Cabrillo , A. Calderon , J. Duarte Campderros , M. Fernandez , C. Fernandez Madrazo , G. Gomez , C. Lasosa García , C. Martinez Rivero , P. Martinez Ruiz del Arbol , F. Matorras , P. Matorras Cuevas , E. Navarrete Ramos , J. Piedra Gomez , L. Scodellaro , I. Vila , J.M. Vizán García 

University of Colombo, Colombo, Sri Lanka

M.K. Jayananda , B. Kailasapathy⁶² , D.U.J. Sonnadara , D.D.C. Wickramarathna 

University of Ruhuna, Department of Physics, Matara, Sri Lanka

W.G.D. Dharmaratna⁶³ , K. Liyanage , N. Perera , N. Wickramage 

CERN, European Organization for Nuclear Research, Geneva, Switzerland

D. Abbaneo , C. Amendola , E. Auffray , G. Auzinger , J. Baechler, D. Barney , A. Bermúdez Martínez , M. Bianco , B. Bilin , A.A. Bin Anuar , A. Bocci , E. Brondolin , C. Caillol , T. Camporesi , G. Cerminara , N. Chernyavskaya , D. d'Enterria , A. Dabrowski , A. David , A. De Roeck , M.M. Defranchis , M. Deile , M. Dobson , F. Fallavollita⁶⁴, L. Forthomme , G. Franzoni , W. Funk , S. Giani, D. Gigi, K. Gill , F. Glege , L. Gouskos , M. Haranko , J. Hegeman , B. Huber, V. Innocente , T. James , P. Janot , J. Kieseler , S. Laurila , P. Lecoq , E. Leutgeb , C. Lourenço , B. Maier , L. Malgeri , M. Mannelli , A.C. Marini , M. Matthewman, F. Meijers , S. Mersi , E. Meschi , V. Milosevic , F. Moortgat , M. Mulders , S. Orfanelli, F. Pantaleo , G. Petrucciani , A. Pfeiffer , M. Pierini , D. Piparo , H. Qu , D. Rabady , G. Reales Gutiérrez, M. Rovere , H. Sakulin , S. Scarfi , C. Schwick, M. Selvaggi , A. Sharma , K. Shchelina , P. Silva , P. Sphicas⁶⁵ , A.G. Stahl Leiton , A. Steen , S. Summers , D. Treille , P. Tropea , A. Tsirou, D. Walter , J. Wanczyk⁶⁶ , K.A. Wozniak⁶⁷ , P. Zehetner , P. Zjendl , W.D. Zeuner

Paul Scherrer Institut, Villigen, Switzerland

T. Bevilacqua⁶⁸ , L. Caminada⁶⁸ , A. Ebrahimi , W. Erdmann , R. Horisberger , Q. Ingram , H.C. Kaestli , D. Kotlinski , C. Lange , M. Missiroli⁶⁸ , L. Noehte⁶⁸ , T. Rohe 

ETH Zurich - Institute for Particle Physics and Astrophysics (IPA), Zurich, Switzerland

T.K. Aarrestad , K. Androsova⁶⁶ , M. Backhaus , A. Calandri , C. Cazzaniga , K. Datta , A. De Cosa , G. Dissertori , M. Dittmar, M. Donegà , F. Eble , M. Galli , K. Gedia , F. Glessgen , C. Grab , D. Hits , W. Lustermann , A.-M. Lyon , R.A. Manzoni , M. Marchegiani , L. Marchese , C. Martin Perez , A. Mascellani⁶⁶ , F. Nessi-Tedaldi , F. Pauss , V. Perovic , S. Pigazzini , M.G. Ratti , M. Reichmann , C. Reissel , T. Reitenspiess , B. Ristic , F. Riti , D. Ruini, D.A. Sanz Becerra , R. Seidita , J. Steggemann⁶⁶ , D. Valsecchi , R. Wallny 

Universität Zürich, Zurich, Switzerland

C. Amsler⁶⁹ , P. Bäertschi , C. Botta , D. Brzhechko, M.F. Canelli , K. Cormier 

R. Del Burgo, J.K. Heikkilä , M. Huwiler , W. Jin , A. Jofrehei , B. Kilminster , S. Leontsinis , S.P. Liechti , A. Macchiolo , P. Meiring , V.M. Mikuni , U. Molinatti , I. Neutelings , A. Reimers , P. Robmann, S. Sanchez Cruz , K. Schweiger , M. Senger , Y. Takahashi , R. Tramontano 

National Central University, Chung-Li, Taiwan

C. Adloff⁷⁰, C.M. Kuo, W. Lin, P.K. Rout , P.C. Tiwari⁴² , S.S. Yu 

National Taiwan University (NTU), Taipei, Taiwan

L. Ceard, Y. Chao , K.F. Chen , P.s. Chen, Z.g. Chen, W.-S. Hou , T.h. Hsu, Y.w. Kao, R. Khurana, G. Kole , Y.y. Li , R.-S. Lu , E. Paganis , A. Psallidas, X.f. Su, J. Thomas-Wilsker , L.s. Tsai, H.y. Wu, E. Yazgan 

High Energy Physics Research Unit, Department of Physics, Faculty of Science, Chulalongkorn University, Bangkok, Thailand

C. Asawatrangkuldee , N. Srimanobhas , V. Wachirapusanand 

Çukurova University, Physics Department, Science and Art Faculty, Adana, Turkey

D. Agyel , F. Boran , Z.S. Demiroglu , F. Dolek , I. Dumanoglu⁷¹ , E. Eskut , Y. Guler⁷² , E. Gurpinar Guler⁷² , C. Isik , O. Kara, A. Kayis Topaksu , U. Kiminsu , G. Onengut , K. Ozdemir⁷³ , A. Polatoz , B. Tali⁷⁴ , U.G. Tok , S. Turkcapar , E. Uslan , I.S. Zorbakir 

Middle East Technical University, Physics Department, Ankara, Turkey

M. Yalvac⁷⁵ 

Bogazici University, Istanbul, Turkey

B. Akgun , I.O. Atakisi , E. Gülmez , M. Kaya⁷⁶ , O. Kaya⁷⁷ , S. Tekten⁷⁸ 

Istanbul Technical University, Istanbul, Turkey

A. Cakir , K. Cankocak^{71,79} , Y. Komurcu , S. Sen⁸⁰ 

Istanbul University, Istanbul, Turkey

O. Aydilek , S. Cerci⁷⁴ , V. Epshteyn , B. Haciasahinoglu , I. Hos⁸¹ , B. Isildak⁸² , B. Kaynak , S. Ozkorucuklu , O. Potok , H. Sert , C. Simsek , D. Sunar Cerci⁷⁴ , C. Zorbilmez 

Institute for Scintillation Materials of National Academy of Science of Ukraine, Kharkiv, Ukraine

A. Boyaryntsev , B. Grynyov 

National Science Centre, Kharkiv Institute of Physics and Technology, Kharkiv, Ukraine

L. Levchuk 

University of Bristol, Bristol, United Kingdom

D. Anthony , J.J. Brooke , A. Bundock , F. Bury , E. Clement , D. Cussans , H. Flacher , M. Glowacki, J. Goldstein , H.F. Heath , L. Kreczko , B. Krikler , S. Paramesvaran , S. Seif El Nasr-Storey, V.J. Smith , N. Stylianou⁸³ , K. Walkingshaw Pass, R. White 

Rutherford Appleton Laboratory, Didcot, United Kingdom

A.H. Ball, K.W. Bell , A. Belyaev⁸⁴ , C. Brew , R.M. Brown , D.J.A. Cockerill , C. Cooke , K.V. Ellis, K. Harder , S. Harper , M.-L. Holmberg⁸⁵ , J. Linacre , K. Manolopoulos, D.M. Newbold , E. Olaiya, D. Petyt , T. Reis , G. Salvi , T. Schuh, C.H. Shepherd-Themistocleous , I.R. Tomalin , T. Williams 

Imperial College, London, United Kingdom

R. Bainbridge , P. Bloch , C.E. Brown , O. Buchmuller, V. Cacchio, C.A. Carrillo Montoya , G.S. Chahal⁸⁶ , D. Colling , J.S. Dancu, P. Dauncey , G. Davies , J. Davies, M. Della Negra , S. Fayer, G. Fedi , G. Hall , M.H. Hassanshahi , A. Howard, G. Iles , M. Knight , J. Langford , L. Lyons , A.-M. Magnan , S. Malik, A. Martelli , M. Mieskolainen , J. Nash⁸⁷ , M. Pesaresi, B.C. Radburn-Smith , A. Richards, A. Rose , C. Seez , R. Shukla , A. Tapper , K. Uchida , G.P. Uttley , L.H. Vage, T. Virdee³² , M. Vojinovic , N. Wardle , D. Winterbottom 

Brunel University, Uxbridge, United Kingdom

K. Coldham, J.E. Cole , A. Khan, P. Kyberd , I.D. Reid 

Baylor University, Waco, Texas, USA

S. Abdullin , A. Brinkerhoff , B. Caraway , J. Dittmann , K. Hatakeyama , J. Hiltbrand , A.R. Kanuganti , B. McMaster , M. Saunders , S. Sawant , C. Sutantawibul , M. Toms⁸⁸ , J. Wilson 

Catholic University of America, Washington, DC, USA

R. Bartek , A. Dominguez , C. Huerta Escamilla, A.E. Simsek , R. Uniyal , A.M. Vargas Hernandez 

The University of Alabama, Tuscaloosa, Alabama, USA

R. Chudasama , S.I. Cooper , S.V. Gleyzer , C.U. Perez , P. Rumerio⁸⁹ , E. Usai , C. West , R. Yi 

Boston University, Boston, Massachusetts, USA

A. Akpinar , A. Albert , D. Arcaro , C. Cosby , Z. Demiragli , C. Erice , E. Fontanesi , D. Gastler , S. Jeon , J. Rohlf , K. Salyer , D. Sperka , D. Spitzbart , I. Suarez , A. Tsatsos , S. Yuan 

Brown University, Providence, Rhode Island, USA

G. Benelli , X. Coubez²⁷, D. Cutts , M. Hadley , U. Heintz , J.M. Hogan⁹⁰ , T. Kwon , G. Landsberg , K.T. Lau , D. Li , J. Luo , S. Mondal , M. Narain[†] , N. Pervan , S. Sagir⁹¹ , F. Simpson , M. Stamenkovic , W.Y. Wong, X. Yan , W. Zhang

University of California, Davis, Davis, California, USA

S. Abbott , J. Bonilla , C. Brainerd , R. Breedon , M. Calderon De La Barca Sanchez , M. Chertok , M. Citron , J. Conway , P.T. Cox , R. Erbacher , F. Jensen , O. Kukral , G. Mocellin , M. Mulhearn , D. Pellett , W. Wei , Y. Yao , F. Zhang 

University of California, Los Angeles, California, USA

M. Bachtis , R. Cousins , A. Datta , J. Hauser , M. Ignatenko , M.A. Iqbal , T. Lam , E. Manca , D. Saltzberg , V. Valuev 

University of California, Riverside, Riverside, California, USA

R. Clare , M. Gordon, G. Hanson , W. Si , S. Wimpenny[†] 

University of California, San Diego, La Jolla, California, USA

J.G. Branson , S. Cittolin , S. Cooperstein , D. Diaz , J. Duarte , L. Giannini , J. Guiang , R. Kansal , V. Krutelyov , R. Lee , J. Letts , M. Masciovecchio , F. Mokhtar , M. Pieri , M. Quinnan , B.V. Sathia Narayanan , V. Sharma , M. Tadel , E. Vourliotis , F. Würthwein , Y. Xiang , A. Yagil 

University of California, Santa Barbara - Department of Physics, Santa Barbara, California, USA

A. Barzdukas , L. Brennan , C. Campagnari , G. Collura , A. Dorsett , J. Incandela , M. Kilpatrick , J. Kim , A.J. Li , P. Masterson , H. Mei , M. Oshiro , J. Richman , U. Sarica , R. Schmitz , F. Setti , J. Sheplock , D. Stuart , S. Wang 

California Institute of Technology, Pasadena, California, USA

A. Bornheim , O. Cerri , A. Latorre , J.M. Lawhorn , J. Mao , H.B. Newman , T. Q. Nguyen , M. Spiropulu , J.R. Vlimant , C. Wang , S. Xie , R.Y. Zhu 

Carnegie Mellon University, Pittsburgh, Pennsylvania, USA

J. Alison , S. An , M.B. Andrews , P. Bryant , V. Dutta , T. Ferguson , A. Harilal , C. Liu , T. Mudholkar , S. Murthy , M. Paulini , A. Roberts , A. Sanchez , W. Terrill 

University of Colorado Boulder, Boulder, Colorado, USA

J.P. Cumalat , W.T. Ford , A. Hassani , G. Karathanasis , E. MacDonald , N. Manganello , F. Marini , A. Perloff , C. Savard , N. Schonbeck , K. Stenson , K.A. Ulmer , S.R. Wagner , N. Zipper 

Cornell University, Ithaca, New York, USA

J. Alexander , S. Bright-Thonney , X. Chen , D.J. Cranshaw , J. Fan , X. Fan , D. Gadkari , S. Hogan , J. Monroy , J.R. Patterson , J. Reichert , M. Reid , A. Ryd , J. Thom , P. Wittich , R. Zou 

Fermi National Accelerator Laboratory, Batavia, Illinois, USA

M. Albrow , M. Alyari , O. Amram , G. Apollinari , A. Apresyan , L.A.T. Bauerdick , D. Berry , J. Berryhill , P.C. Bhat , K. Burkett , J.N. Butler , A. Canepa , G.B. Cerati , H.W.K. Cheung , F. Chlebana , G. Cummings , J. Dickinson , I. Dutta , V.D. Elvira , Y. Feng , J. Freeman , A. Gandrakota , Z. Gecse , L. Gray , D. Green , A. Grummer , S. Grünendahl , D. Guerrero , O. Gutsche , R.M. Harris , R. Heller , T.C. Herwig , J. Hirschauer , L. Horyn , B. Jayatilaka , S. Jindariani , M. Johnson , U. Joshi , T. Klijnsma , B. Klima , K.H.M. Kwok , S. Lammel , D. Lincoln , R. Lipton , T. Liu , C. Madrid , K. Maeshima , C. Mantilla , D. Mason , P. McBride , P. Merkel , S. Mrenna , S. Nahn , J. Ngadiuba , D. Noonan , V. Papadimitriou , N. Pastika , K. Pedro , C. Pena⁹² , F. Ravera , A. Reinsvold Hall⁹³ , L. Ristori , E. Sexton-Kennedy , N. Smith , A. Soha , L. Spiegel , S. Stoynev , J. Strait , L. Taylor , S. Tkaczyk , N.V. Tran , L. Uplegger , E.W. Vaandering , I. Zoi

University of Florida, Gainesville, Florida, USA

C. Aruta , P. Avery , D. Bourilkov , L. Cadamuro , P. Chang , V. Cherepanov , R.D. Field , E. Koenig , M. Kolosova , J. Konigsberg , A. Korytov , K.H. Lo , K. Matchev , N. Menendez , G. Mitselmakher , K. Mohrman , A. Muthirakalayil Madhu , N. Rawal , D. Rosenzweig , S. Rosenzweig , K. Shi , J. Wang 

Florida State University, Tallahassee, Florida, USA

T. Adams , A. Al Kadhimi , A. Askew , N. Bower , R. Habibullah , V. Hagopian , R. Hashmi , R.S. Kim , S. Kim , T. Kolberg , G. Martinez , H. Prosper , P.R. Prova , O. Viazlo , M. Wulansatiti , R. Yohay , J. Zhang 

Florida Institute of Technology, Melbourne, Florida, USA

B. Alsufyani , M.M. Baarmand , S. Butalla , T. Elkafrawy⁵⁷ , M. Hohlmann , R. Kumar Verma , M. Rahmani 

University of Illinois Chicago, Chicago, USA, Chicago, USA

M.R. Adams , C. Bennett , R. Cavanaugh , S. Dittmer , R. Escobar Franco , O. Evdoki-

mov , C.E. Gerber , D.J. Hofman , J.h. Lee , D. S. Lemos , A.H. Merrit , C. Mills , S. Nanda , G. Oh , B. Ozek , D. Pilipovic , T. Roy , S. Rudrabhatla , M.B. Tonjes , N. Varelas , X. Wang , Z. Ye , J. Yoo 

The University of Iowa, Iowa City, Iowa, USA

M. Alhusseini , D. Blend , K. Dilsiz⁹⁴ , L. Emediato , G. Karaman , O.K. Köseyan , J.-P. Merlo , A. Mestvirishvili⁹⁵ , J. Nachtman , O. Neogi , H. Ogul⁹⁶ , Y. Onel , A. Penzo , C. Snyder , E. Tiras⁹⁷ 

Johns Hopkins University, Baltimore, Maryland, USA

B. Blumenfeld , L. Corcodilos , J. Davis , A.V. Gritsan , L. Kang , S. Kyriacou , P. Maksimovic , M. Roguljic , J. Roskes , S. Sekhar , M. Swartz , T.Á. Vámi 

The University of Kansas, Lawrence, Kansas, USA

A. Abreu , L.F. Alcerro Alcerro , J. Anguiano , P. Baringer , A. Bean , Z. Flowers , D. Grove , J. King , G. Krintiras , M. Lazarovits , C. Le Mahieu , C. Lindsey , J. Marquez , N. Minafra , M. Murray , M. Nickel , M. Pitt , S. Popescu⁹⁸ , C. Rogan , C. Royon , R. Salvatico , S. Sanders , C. Smith , Q. Wang , G. Wilson 

Kansas State University, Manhattan, Kansas, USA

B. Allmond , A. Ivanov , K. Kaadze , A. Kalogeropoulos , D. Kim , Y. Maravin , K. Nam , J. Natoli , D. Roy , G. Sorrentino 

Lawrence Livermore National Laboratory, Livermore, California, USA

F. Rebassoo , D. Wright 

University of Maryland, College Park, Maryland, USA

A. Baden , A. Belloni , A. Bethani , Y.M. Chen , S.C. Eno , N.J. Hadley , S. Jabeen , R.G. Kellogg , T. Koeth , Y. Lai , S. Lascio , A.C. Mignerey , S. Nabili , C. Palmer , C. Papageorgakis , M.M. Paranjpe , L. Wang 

Massachusetts Institute of Technology, Cambridge, Massachusetts, USA

J. Bendavid , W. Busza , I.A. Cali , Y. Chen , M. D'Alfonso , J. Eysermans , C. Freer , G. Gomez-Ceballos , M. Goncharov , P. Harris , D. Hoang , D. Kovalskyi , J. Krupa , L. Lavezzo , Y.-J. Lee , K. Long , C. Mironov , C. Paus , D. Rankin , C. Roland , G. Roland , S. Rothman , Z. Shi , G.S.F. Stephans , J. Wang , Z. Wang , B. Wyslouch , T. J. Yang 

University of Minnesota, Minneapolis, Minnesota, USA

B. Crossman , B.M. Joshi , C. Kapsiak , M. Krohn , D. Mahon , J. Mans , B. Marzocchi , S. Pandey , M. Revering , R. Rusack , R. Saradhy , N. Schroeder , N. Strobbe , M.A. Wadud 

University of Mississippi, Oxford, Mississippi, USA

L.M. Cremaldi 

University of Nebraska-Lincoln, Lincoln, Nebraska, USA

K. Bloom , M. Bryson , D.R. Claes , C. Fangmeier , F. Golf , G. Haza , J. Hossain , C. Joo , I. Kravchenko , I. Reed , J.E. Siado , W. Tabb , A. Vagnerini , A. Wightman , F. Yan , D. Yu , A.G. Zecchinelli 

State University of New York at Buffalo, Buffalo, New York, USA

G. Agarwal , H. Bandyopadhyay , L. Hay , I. Iashvili , A. Kharchilava , M. Morris , D. Nguyen , S. Rappoccio , H. Rejeb Sfar , A. Williams 

Northeastern University, Boston, Massachusetts, USA

E. Barberis , Y. Haddad , Y. Han , A. Krishna , J. Li , M. Lu , G. Madigan ,
R. McCarthy , D.M. Morse , V. Nguyen , T. Orimoto , A. Parker , L. Skinnari ,
A. Tishelman-Charny , B. Wang , D. Wood 

Northwestern University, Evanston, Illinois, USA

S. Bhattacharya , J. Bueghly , Z. Chen , K.A. Hahn , Y. Liu , Y. Miao , D.G. Monk ,
M.H. Schmitt , A. Taliencio , M. Velasco

University of Notre Dame, Notre Dame, Indiana, USA

R. Band , R. Bucci , S. Castells , M. Cremonesi , A. Das , R. Goldouzian , M. Hildreth ,
K.W. Ho , K. Hurtado Anampa , C. Jessop , K. Lannon , J. Lawrence , N. Loukas ,
L. Lutton , J. Mariano , N. Marinelli , I. Mcalister , T. McCauley , C. Mcgrady , C. Moore ,
Y. Musienko¹⁶ , H. Nelson , M. Osherson , R. Ruchti , A. Townsend , M. Wayne ,
H. Yockey , M. Zarucki , L. Zygala 

The Ohio State University, Columbus, Ohio, USA

A. Basnet , B. Bylsma , M. Carrigan , L.S. Durkin , C. Hill , M. Joyce , A. Lesauvage ,
M. Nunez Ornelas , K. Wei , B.L. Winer , B. R. Yates 

Princeton University, Princeton, New Jersey, USA

F.M. Addesa , H. Bouchamaoui , P. Das , G. Dezoort , P. Elmer , A. Frankenthal ,
B. Greenberg , N. Haubrich , S. Higginbotham , G. Kopp , S. Kwan , D. Lange ,
A. Loeliger , D. Marlow , I. Ojalvo , J. Olsen , A. Shevelev , D. Stickland , C. Tully 

University of Puerto Rico, Mayaguez, Puerto Rico, USA

S. Malik 

Purdue University, West Lafayette, Indiana, USA

A.S. Bakshi , V.E. Barnes , S. Chandra , R. Chawla , S. Das , A. Gu , L. Gutay ,
M. Jones , A.W. Jung , D. Kondratyev , A.M. Koshy , M. Liu , G. Negro ,
N. Neumeister , G. Paspalaki , S. Piperov , V. Scheurer , J.F. Schulte , M. Stojanovic ,
J. Thieman , A. K. Viridi , F. Wang , W. Xie 

Purdue University Northwest, Hammond, Indiana, USA

J. Dolen , N. Parashar , A. Pathak 

Rice University, Houston, Texas, USA

D. Acosta , A. Baty , T. Carnahan , K.M. Ecklund , P.J. Fernández Manteca , S. Freed ,
P. Gardner , F.J.M. Geurts , A. Kumar , W. Li , O. Miguel Colin , B.P. Padley ,
R. Redjimi , J. Rotter , E. Yigitbasi , Y. Zhang 

University of Rochester, Rochester, New York, USA

A. Bodek , P. de Barbaro , R. Demina , J.L. Dulemba , C. Fallon , A. Garcia-Bellido , O. Hindrichs , A. Khukhunaishvili , P. Parygin⁸⁸ , E. Popova⁸⁸ , R. Taus ,
G.P. Van Onsem 

The Rockefeller University, New York, New York, USA

K. Goulianos 

Rutgers, The State University of New Jersey, Piscataway, New Jersey, USA

B. Chiarito , J.P. Chou , Y. Gershtein , E. Halkiadakis , A. Hart , M. Heindl ,
D. Jaroslawski , O. Karacheban³⁰ , I. Laflotte , A. Lath , R. Montalvo , K. Nash ,
H. Routray , S. Salur , S. Schnetzer , S. Somalwar , R. Stone , S.A. Thayil , S. Thomas ,
J. Vora , H. Wang 

University of Tennessee, Knoxville, Tennessee, USA

H. Acharya, D. Ally , A.G. Delannoy , S. Fiorendi , T. Holmes , N. Karunarathna , L. Lee , E. Nibigira , S. Spanier 

Texas A&M University, College Station, Texas, USA

D. Aebi , M. Ahmad , O. Bouhali⁹⁹ , M. Dalchenko , R. Eusebi , J. Gilmore , T. Huang , T. Kamon¹⁰⁰ , H. Kim , S. Luo , S. Malhotra, R. Mueller , D. Overton , D. Rathjens , A. Safonov 

Texas Tech University, Lubbock, Texas, USA

N. Akchurin , J. Damgov , V. Hegde , A. Hussain , Y. Kazhykarim, K. Lamichhane , S.W. Lee , A. Mankel , T. Mengke, S. Muthumuni , T. Peltola , I. Volobouev , A. Whitbeck 

Vanderbilt University, Nashville, Tennessee, USA

E. Appelt , S. Greene, A. Gurrola , W. Johns , R. Kunnawalkam Elayavalli , A. Melo , F. Romeo , P. Sheldon , S. Tuo , J. Velkovska , J. Viinikainen 

University of Virginia, Charlottesville, Virginia, USA

B. Cardwell , B. Cox , J. Hakala , R. Hirosky , A. Ledovskoy , A. Li , C. Neu , C.E. Perez Lara 

Wayne State University, Detroit, Michigan, USA

P.E. Karchin 

University of Wisconsin - Madison, Madison, Wisconsin, USA

A. Aravind, S. Banerjee , K. Black , T. Bose , S. Dasu , I. De Bruyn , P. Everaerts , C. Galloni, H. He , M. Herndon , A. Herve , C.K. Koraka , A. Lanaro, R. Loveless , J. Madhusudanan Sreekala , A. Mallampalli , A. Mohammadi , S. Mondal, G. Parida , D. Pinna, A. Savin, V. Shang , V. Sharma , W.H. Smith , D. Teague, H.F. Tsoi , W. Vetens , A. Warden 

Authors affiliated with an institute or an international laboratory covered by a cooperation agreement with CERN

S. Afanasiev , V. Andreev , Yu. Andreev , T. Aushev , M. Azarkin , A. Babaev , A. Belyaev , V. Blinov¹⁰¹ , E. Boos , V. Borshch , D. Budkouski , V. Bunichev , M. Chadeeva¹⁰¹ , V. Chekhovsky, M. Danilov¹⁰¹ , A. Dermenev , T. Dimova¹⁰¹ , D. Druzhkin¹⁰² , M. Dubinin⁹² , L. Dudko , G. Gavrillov , V. Gavrillov , S. Gninenko , V. Golovtsov , N. Golubev , I. Golutvin , I. Gorbunov , Y. Ivanov , V. Kachanov , L. Kardapoltsev¹⁰¹ , V. Karjavine , A. Karneyeu , V. Kim¹⁰¹ , M. Kirakosyan, D. Kirpichnikov , M. Kirsanov , V. Klyukhin , O. Kodolova¹⁰³ , D. Konstantinov , V. Korenkov , A. Kozyrev¹⁰¹ , N. Krasnikov , A. Lanev , P. Levchenko¹⁰⁴ , N. Lychkovskaya , V. Makarenko , A. Malakhov , V. Matveev¹⁰¹ , V. Murzin , A. Nikitenko^{105,103} , S. Obraztsov , V. Oreshkin , V. Palichik , V. Perelygin , M. Perfilov, S. Petrushanko , S. Polikarpov¹⁰¹ , V. Popov, O. Radchenko¹⁰¹ , M. Savina , V. Savrin , V. Shalaev , S. Shmatov , S. Shulha , Y. Skovpen¹⁰¹ , S. Slabospitskii , V. Smirnov , A. Snigirev , D. Sosnov , V. Sulimov , E. Tcherniaev , A. Terkulov , O. Teryaev , I. Tlisova , A. Toropin , L. Uvarov , A. Uzunian , A. Vorobyev[†], G. Vorotnikov , N. Voytishin , B.S. Yuldashev¹⁰⁶, A. Zarubin , I. Zhizhin , A. Zhokin 

†: Deceased

¹Also at Yerevan State University, Yerevan, Armenia

²Also at TU Wien, Vienna, Austria

-
- ³Also at Institute of Basic and Applied Sciences, Faculty of Engineering, Arab Academy for Science, Technology and Maritime Transport, Alexandria, Egypt
- ⁴Also at Ghent University, Ghent, Belgium
- ⁵Also at Universidade Estadual de Campinas, Campinas, Brazil
- ⁶Also at Federal University of Rio Grande do Sul, Porto Alegre, Brazil
- ⁷Also at UFMS, Nova Andradina, Brazil
- ⁸Also at Nanjing Normal University, Nanjing, China
- ⁹Now at Henan Normal University, Xinxiang, China
- ¹⁰Now at The University of Iowa, Iowa City, Iowa, USA
- ¹¹Also at University of Chinese Academy of Sciences, Beijing, China
- ¹²Also at China Center of Advanced Science and Technology, Beijing, China
- ¹³Also at University of Chinese Academy of Sciences, Beijing, China
- ¹⁴Also at China Spallation Neutron Source, Guangdong, China
- ¹⁵Also at Université Libre de Bruxelles, Bruxelles, Belgium
- ¹⁶Also at an institute or an international laboratory covered by a cooperation agreement with CERN
- ¹⁷Now at British University in Egypt, Cairo, Egypt
- ¹⁸Now at Cairo University, Cairo, Egypt
- ¹⁹Also at Birla Institute of Technology, Mesra, Mesra, India
- ²⁰Also at Purdue University, West Lafayette, Indiana, USA
- ²¹Also at Université de Haute Alsace, Mulhouse, France
- ²²Also at Department of Physics, Tsinghua University, Beijing, China
- ²³Also at Tbilisi State University, Tbilisi, Georgia
- ²⁴Also at The University of the State of Amazonas, Manaus, Brazil
- ²⁵Also at Erzincan Binali Yildirim University, Erzincan, Turkey
- ²⁶Also at University of Hamburg, Hamburg, Germany
- ²⁷Also at RWTH Aachen University, III. Physikalisches Institut A, Aachen, Germany
- ²⁸Also at Isfahan University of Technology, Isfahan, Iran
- ²⁹Also at Bergische University Wuppertal (BUW), Wuppertal, Germany
- ³⁰Also at Brandenburg University of Technology, Cottbus, Germany
- ³¹Also at Forschungszentrum Jülich, Juelich, Germany
- ³²Also at CERN, European Organization for Nuclear Research, Geneva, Switzerland
- ³³Also at Institute of Physics, University of Debrecen, Debrecen, Hungary
- ³⁴Also at Institute of Nuclear Research ATOMKI, Debrecen, Hungary
- ³⁵Now at Universitatea Babes-Bolyai - Facultatea de Fizica, Cluj-Napoca, Romania
- ³⁶Also at Physics Department, Faculty of Science, Assiut University, Assiut, Egypt
- ³⁷Also at HUN-REN Wigner Research Centre for Physics, Budapest, Hungary
- ³⁸Also at Faculty of Informatics, University of Debrecen, Debrecen, Hungary
- ³⁹Also at Punjab Agricultural University, Ludhiana, India
- ⁴⁰Also at University of Hyderabad, Hyderabad, India
- ⁴¹Also at University of Visva-Bharati, Santiniketan, India
- ⁴²Also at Indian Institute of Science (IISc), Bangalore, India
- ⁴³Also at IIT Bhubaneswar, Bhubaneswar, India
- ⁴⁴Also at Institute of Physics, Bhubaneswar, India
- ⁴⁵Also at Deutsches Elektronen-Synchrotron, Hamburg, Germany
- ⁴⁶Also at Department of Physics, Isfahan University of Technology, Isfahan, Iran
- ⁴⁷Also at Sharif University of Technology, Tehran, Iran
- ⁴⁸Also at Department of Physics, University of Science and Technology of Mazandaran, Behshahr, Iran

⁴⁹Also at Helwan University, Cairo, Egypt

⁵⁰Also at Italian National Agency for New Technologies, Energy and Sustainable Economic Development, Bologna, Italy

⁵¹Also at Centro Siciliano di Fisica Nucleare e di Struttura Della Materia, Catania, Italy

⁵²Also at Università degli Studi Guglielmo Marconi, Roma, Italy

⁵³Also at Scuola Superiore Meridionale, Università di Napoli 'Federico II', Napoli, Italy

⁵⁴Also at Fermi National Accelerator Laboratory, Batavia, Illinois, USA

⁵⁵Also at Laboratori Nazionali di Legnaro dell'INFN, Legnaro, Italy

⁵⁶Also at Università di Napoli 'Federico II', Napoli, Italy

⁵⁷Also at Ain Shams University, Cairo, Egypt

⁵⁸Also at Consiglio Nazionale delle Ricerche - Istituto Officina dei Materiali, Perugia, Italy

⁵⁹Also at Riga Technical University, Riga, Latvia

⁶⁰Also at Department of Applied Physics, Faculty of Science and Technology, Universiti Kebangsaan Malaysia, Bangi, Malaysia

⁶¹Also at Consejo Nacional de Ciencia y Tecnología, Mexico City, Mexico

⁶²Also at Trincomalee Campus, Eastern University, Sri Lanka, Nilaveli, Sri Lanka

⁶³Also at Saegis Campus, Nugegoda, Sri Lanka

⁶⁴Also at INFN Sezione di Pavia, Università di Pavia, Pavia, Italy

⁶⁵Also at National and Kapodistrian University of Athens, Athens, Greece

⁶⁶Also at Ecole Polytechnique Fédérale Lausanne, Lausanne, Switzerland

⁶⁷Also at University of Vienna Faculty of Computer Science, Vienna, Austria

⁶⁸Also at Universität Zürich, Zurich, Switzerland

⁶⁹Also at Stefan Meyer Institute for Subatomic Physics, Vienna, Austria

⁷⁰Also at Laboratoire d'Annecy-le-Vieux de Physique des Particules, IN2P3-CNRS, Annecy-le-Vieux, France

⁷¹Also at Near East University, Research Center of Experimental Health Science, Mersin, Turkey

⁷²Also at Konya Technical University, Konya, Turkey

⁷³Also at Izmir Bakircay University, Izmir, Turkey

⁷⁴Also at Adiyaman University, Adiyaman, Turkey

⁷⁵Also at Bozok Universitetesi Rektörlüğü, Yozgat, Turkey

⁷⁶Also at Marmara University, Istanbul, Turkey

⁷⁷Also at Milli Savunma University, Istanbul, Turkey

⁷⁸Also at Kafkas University, Kars, Turkey

⁷⁹Now at Istanbul Okan University, Istanbul, Turkey

⁸⁰Also at Hacettepe University, Ankara, Turkey

⁸¹Also at Istanbul University - Cerrahpasa, Faculty of Engineering, Istanbul, Turkey

⁸²Also at Yildiz Technical University, Istanbul, Turkey

⁸³Also at Vrije Universiteit Brussel, Brussel, Belgium

⁸⁴Also at School of Physics and Astronomy, University of Southampton, Southampton, United Kingdom

⁸⁵Also at University of Bristol, Bristol, United Kingdom

⁸⁶Also at IPPP Durham University, Durham, United Kingdom

⁸⁷Also at Monash University, Faculty of Science, Clayton, Australia

⁸⁸Now at an institute or an international laboratory covered by a cooperation agreement with CERN

⁸⁹Also at Università di Torino, Torino, Italy

⁹⁰Also at Bethel University, St. Paul, Minnesota, USA

⁹¹Also at Karamanoğlu Mehmetbey University, Karaman, Turkey

⁹²Also at California Institute of Technology, Pasadena, California, USA

⁹³Also at United States Naval Academy, Annapolis, Maryland, USA

⁹⁴Also at Bingol University, Bingol, Turkey

⁹⁵Also at Georgian Technical University, Tbilisi, Georgia

⁹⁶Also at Sinop University, Sinop, Turkey

⁹⁷Also at Erciyes University, Kayseri, Turkey

⁹⁸Also at Horia Hulubei National Institute of Physics and Nuclear Engineering (IFIN-HH), Bucharest, Romania

⁹⁹Also at Texas A&M University at Qatar, Doha, Qatar

¹⁰⁰Also at Kyungpook National University, Daegu, Korea

¹⁰¹Also at another institute or international laboratory covered by a cooperation agreement with CERN

¹⁰²Also at Universiteit Antwerpen, Antwerpen, Belgium

¹⁰³Also at Yerevan Physics Institute, Yerevan, Armenia

¹⁰⁴Also at Northeastern University, Boston, Massachusetts, USA

¹⁰⁵Also at Imperial College, London, United Kingdom

¹⁰⁶Also at Institute of Nuclear Physics of the Uzbekistan Academy of Sciences, Tashkent, Uzbekistan